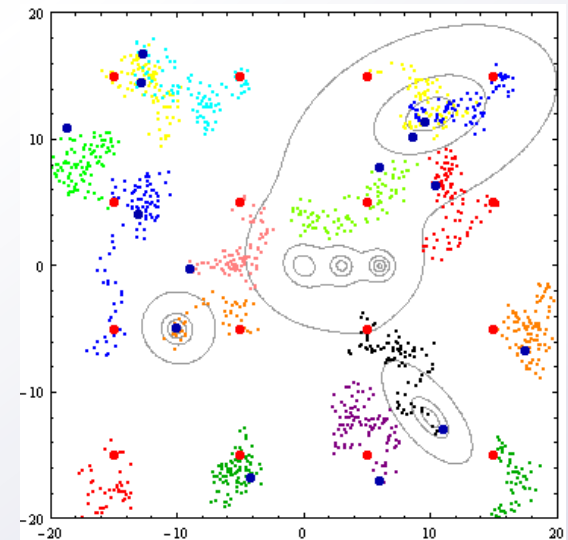
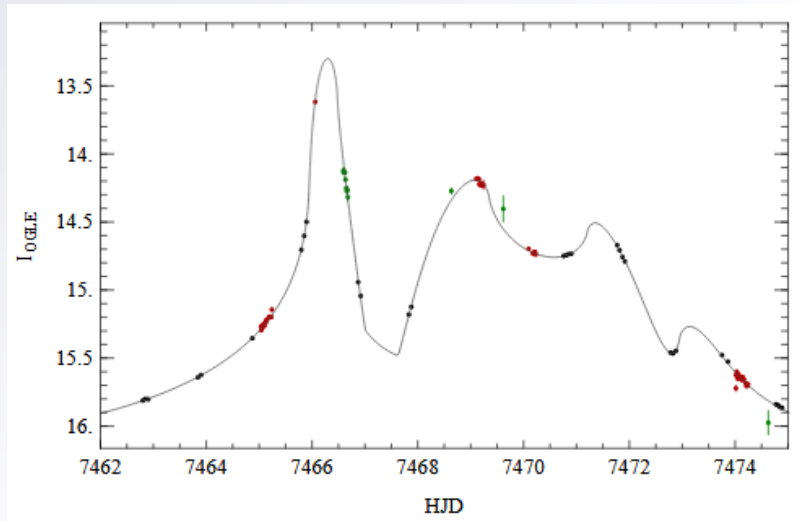


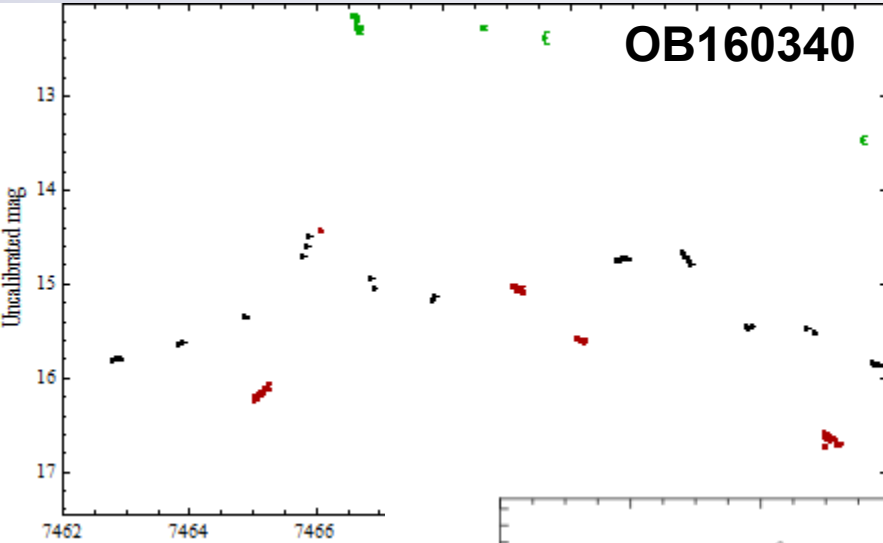
Fitting Light Curves in Microlensing



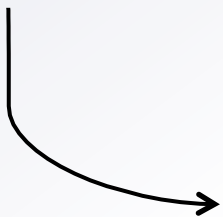
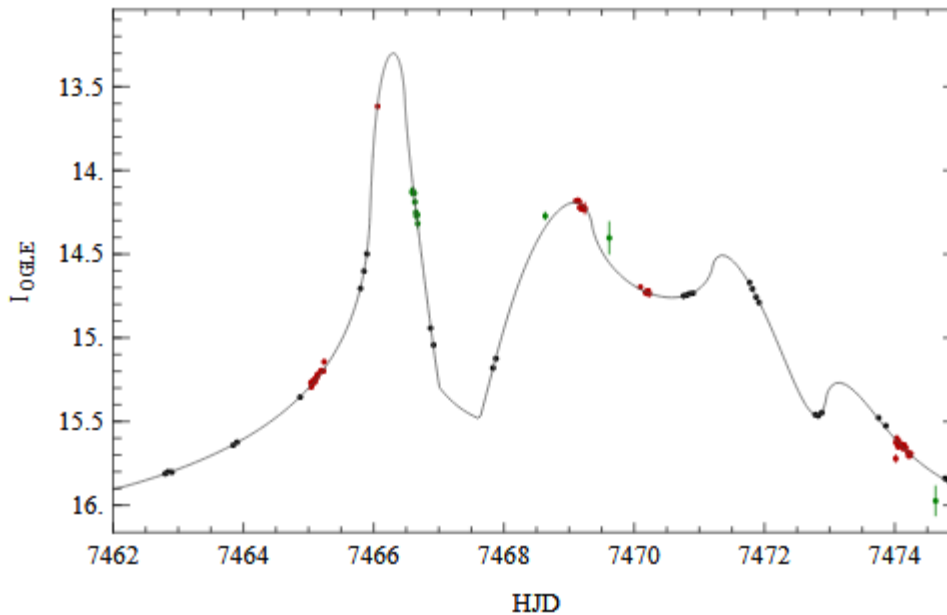
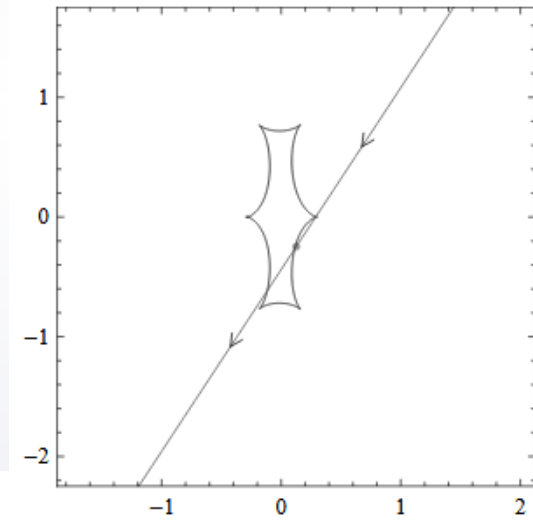
Valerio Bozza
University of Salerno

Our goal

- Given the data, we look for a model to explain them



$$\begin{aligned}t_0 &= 7467.45 \\t_E &= 7.62 \\u_0 &= 0.241 \\\alpha &= 0.988 \\\rho_* &= 0.033 \\s &= 0.906 \\q &= 0.935\end{aligned}$$



Summary

- Finding the best model
 - Downhill methods
 - Markov Chain
- Uncertainty assessment
- Degeneracies
- Bayesian analysis
- Initial conditions

1. Finding the best model

Fitting microlensing events

- If we are able to calculate the magnification for a **given model** at any times, we can easily evaluate the **corresponding χ^2** .
- Binary microlensing light curves are characterized by a minimum of **7 parameters**.
- In addition, for each dataset we have two calibration parameters: **source and background flux**.

$$y_i = F_* f(t_i, \mathbf{p}) + F_B$$

- These parameters come linearly and can be found analytically by a **least-squares fit** for any given model.

$$F_* = \frac{\sum \frac{1}{\sigma_i^2} \sum \frac{f_i y_i}{\sigma_i^2} - \sum \frac{f_i}{\sigma_i^2} \sum \frac{y_i}{\sigma_i^2}}{\sum \frac{1}{\sigma_i^2} \sum \frac{f_i^2}{\sigma_i^2} - \left(\sum \frac{f_i}{\sigma_i^2} \right)^2}; \quad F_B = \frac{\sum \frac{y_i}{\sigma_i^2} \sum \frac{f_i^2}{\sigma_i^2} - \sum \frac{f_i}{\sigma_i^2} \sum \frac{f_i y_i}{\sigma_i^2}}{\sum \frac{1}{\sigma_i^2} \sum \frac{f_i^2}{\sigma_i^2} - \left(\sum \frac{f_i}{\sigma_i^2} \right)^2}$$

- **Now we need a minimization algorithm!**

Steepest descent

- If χ^2 depends on m parameters $\mathbf{p} = \{p_1, \dots, p_m\}$, its gradient is

$$\nabla \chi^2(\mathbf{p}) = -2 \sum_i \mathbf{J}_i \left[\frac{y_i - f(t_i, \mathbf{p})}{\sigma_i^2} \right]$$

$$\text{where } \mathbf{J}_i = \begin{pmatrix} \frac{\partial f_i}{\partial p_1} & \dots & \frac{\partial f_i}{\partial p_m} \end{pmatrix}$$

- The steepest descent is then implemented by choosing

$$\mathbf{p}_{n+1} = \mathbf{p}_n - \alpha \nabla \chi^2$$

- α is determined by a search along the direction of the gradient.

Gauss-Newton method

- Let us set $\mathbf{p}_{n+1} = \mathbf{p}_n + \Delta$.
- If Δ is such that \mathbf{p}_{n+1} is a minimum, then

$$0 = \nabla \chi^2(\mathbf{p}_n + \Delta) = -2 \sum_i \mathbf{J}_i \left[\frac{y_i - f(t_i, \mathbf{p}_n + \Delta)}{\sigma_i^2} \right] \cong$$

$$\cong -2 \sum_i \mathbf{J}_i \left[\frac{y_i - f(t_i, \mathbf{p}_n) - \mathbf{J}_i \cdot \Delta}{\sigma_i^2} \right]$$

- The approximate solution for Δ is obtained by a linear set of equations

$$\sum_i \mathbf{J}_i [\mathbf{J}_i \cdot \Delta] = \sum_i \mathbf{J}_i \left[\frac{y_i - f(t_i, \mathbf{p}_n)}{\sigma_i^2} \right]$$

- Convergence is not guaranteed if we are too far from minimum

Levenberg method

- Interpolates between the two methods, switching from Gauss-Newton to steepest descent when the first fails.
- We modify the normal equations by introducing a **parameter λ**

$$\sum_i \mathbf{J}_i [\mathbf{J}_i \cdot \Delta] + \lambda \Delta = \sum_i \mathbf{J}_i [y_i - f(t_i, \mathbf{p}_i)]$$

- If λ is small, the normal equations work as in Gauss-Newton.
- If λ is large, the new term dominates and Δ is rotated toward the steepest descent direction.

Levenberg-Marquardt algorithm

- Steepest descent may be inefficient if there are directions in which χ^2 is **very flat**.
- The final version of the modified normal equations is

$$\sum_i \left[\mathbf{J}_i (\mathbf{J}_i \cdot \Delta) + \lambda |\mathbf{J}_i|^2 \Delta \right] = \sum_i \mathbf{J}_i [y_i - f(t_i, \mathbf{p}_i)]$$

- In Levenberg-Marquardt algorithm, we start from a value of λ close to 1.
- We calculate Δ ; if $\chi^2(\mathbf{p}_n + \Delta) < \chi^2(\mathbf{p}_n)$, we accept the new point $\mathbf{p}_{n+1} = \mathbf{p}_n + \Delta$ and decrease λ .
- If not, we reject the new point and increase λ .

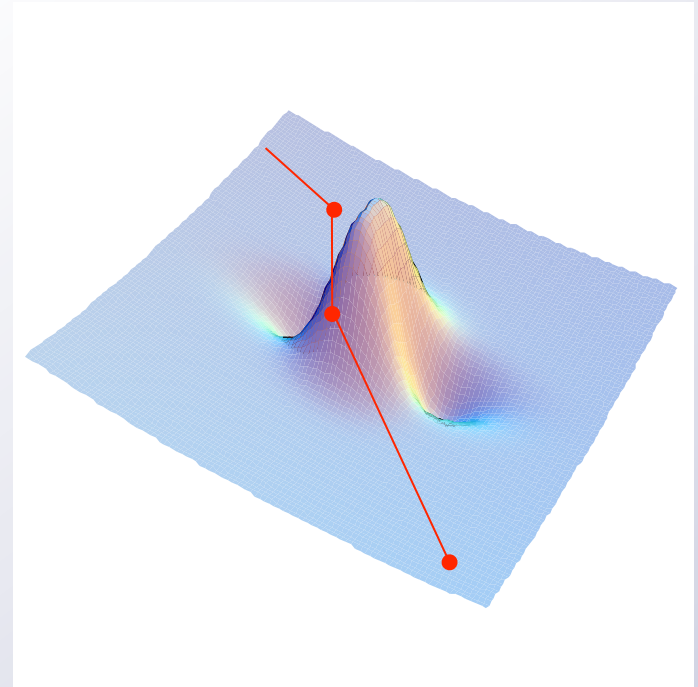
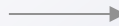
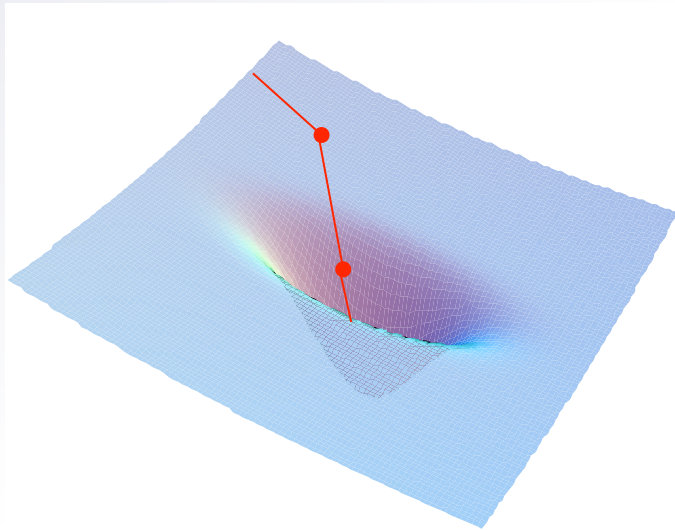
Implementation of Levenberg-Marquardt

- We need to calculate the gradient vector $\mathbf{J}_i = \left(\frac{\partial f_i}{\partial p_1}, \dots, \frac{\partial f_i}{\partial p_m} \right)$
- The derivatives require the calculation of magnification at two points spaced by dp_i . This is the slowest step.
- The resolution of normal equations can be done by standard Gauss method, Cholesky decomposition...
- Levenberg-Marquardt algorithm (nearly) always finds a local minimum.
- It is also very very fast.
- It might get stuck at a local minimum.
- How do we find the best minimum?

1.1 Downhill methods

Jumping out of minima

- One possibility to enlarge our search is to add a penalty on the χ^2 function.
- Once we find the first minimum, we try to fill it with a bumper and run the fit again.
- If the bumper is small, the fit will still remain in the same dip.
- If the bumper is large enough, the fit will jump out of the hole and discover a different minimum.



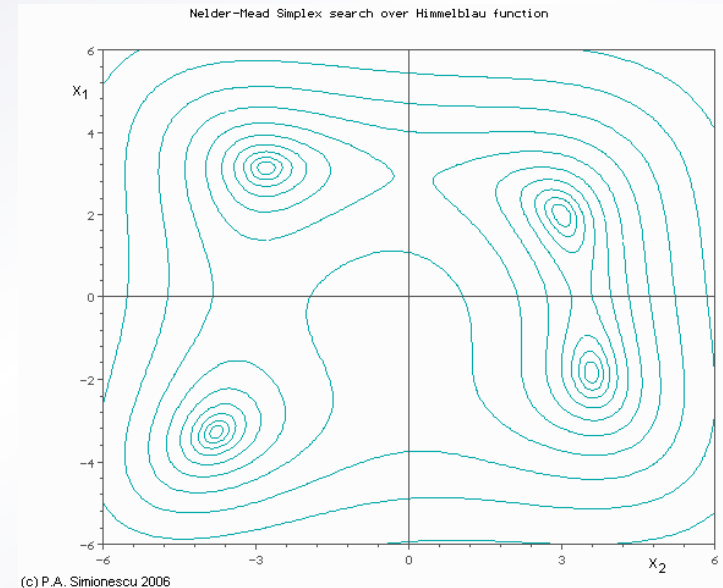
1.1 Downhill methods

Downhill simplex (Nelder-Mead)

- In m dimensions, consider a simplex made of $m+1$ points $\{\mathbf{x}_1, \dots, \mathbf{x}_{m+1}\}$
- Let \mathbf{x}_0 be the barycenter of the best m points.
- The worst point is replaced by its reflection with respect to \mathbf{x}_0 :

$$\mathbf{x}_{new} = \mathbf{x}_0 + \gamma(\mathbf{x}_0 - \mathbf{x}_{m+1})$$

- There are rules for expansion or contraction by tuning γ .
- No need to calculate gradients.

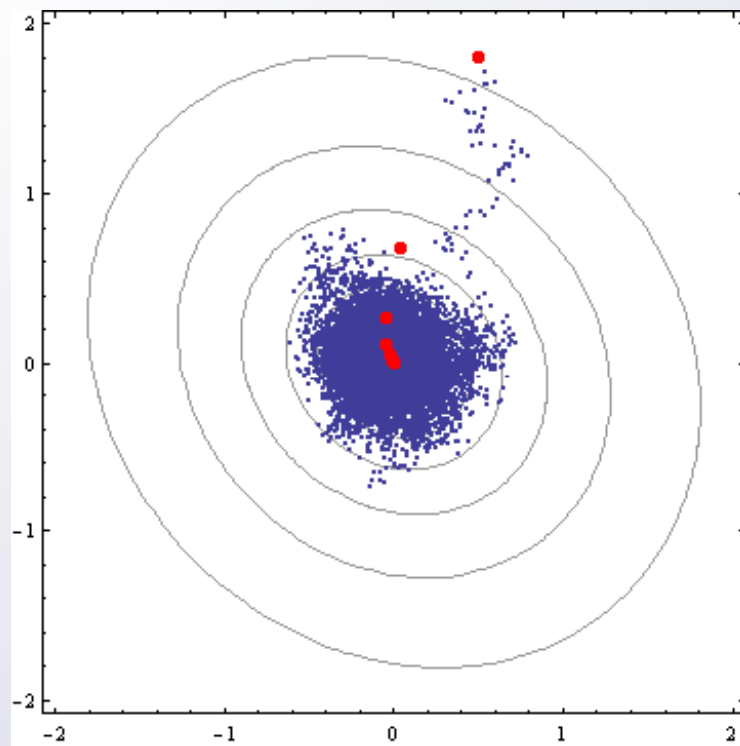


Differential evolution

- Start from a population of $NP \geq 4$ points (“agents”) $\{\mathbf{x}_1, \dots, \mathbf{x}_{NP}\}$
- For each agent \mathbf{x} , pick three more random agents a, b, c .
- Generate a new point \mathbf{y} whose components are
$$y_i = a_i + w(b_i - c_i) \quad \text{with some probability CR}$$
$$y_i = x_i \quad \text{otherwise.}$$
- One random component is always changed.
- If $\chi^2(\mathbf{y}) < \chi^2(\mathbf{x})$ then the new agent replaces the old one.

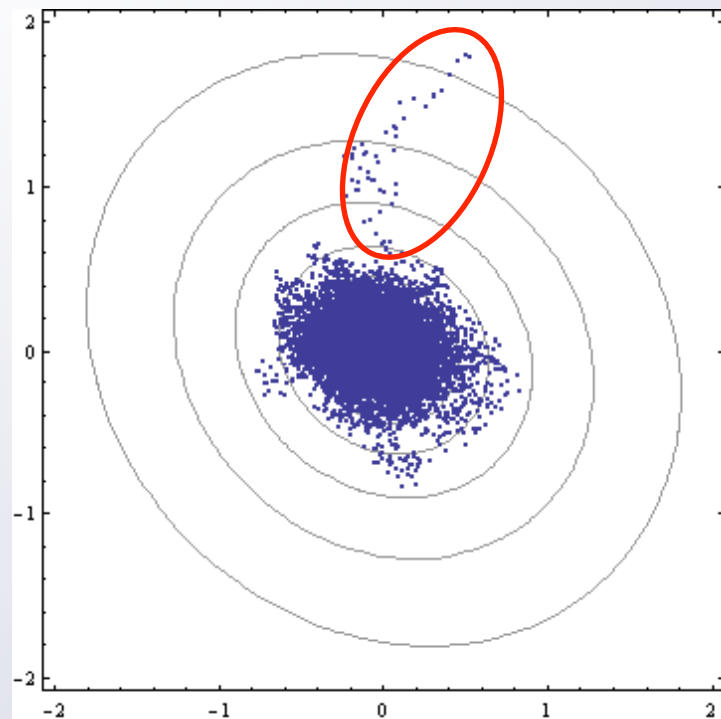
Markov Chain Monte Carlo

- For a recent review see:
“Markov Chain Monte Carlo Methods for Bayesian Data Analysis in Astronomy”, S. Sharma, arXiv:1706.01629.
- MCMC is **NOT** a minimization algorithm!
- MCMC samples a probability distribution:
the best model is just a by-product.
- In this example, after 10000 points, a **Markov chain** finds the best model at accuracy 3×10^{-3} .
- The same accuracy is reached by a **steepest descent** algorithm in 8 steps.



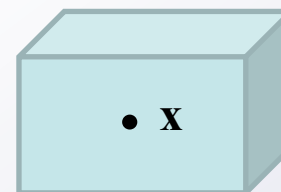
Markov Chain Monte Carlo

- Given the point \mathbf{x}_n in the chain, we randomly draw a candidate new point \mathbf{y} from a **proposal** probability distribution $q(\mathbf{y}|\mathbf{x})$.
- If $p(\mathbf{y}) > p(\mathbf{x})$, we accept the proposal and set $\mathbf{x}_{n+1} = \mathbf{y}$.
- If $p(\mathbf{y}) < p(\mathbf{x})$, we accept the proposal with probability $p(\mathbf{y})/p(\mathbf{x})$, otherwise we set $\mathbf{x}_{n+1} = \mathbf{x}_n$. (*Metropolis algorithm*)
- In the limit of large numbers, the chain will become a **representative sampling** of the probability distribution p .
- The “burn-in” must be discarded.
- In our optimization problems, we set $p = \mathcal{L} = \exp(-\chi^2/2)$.



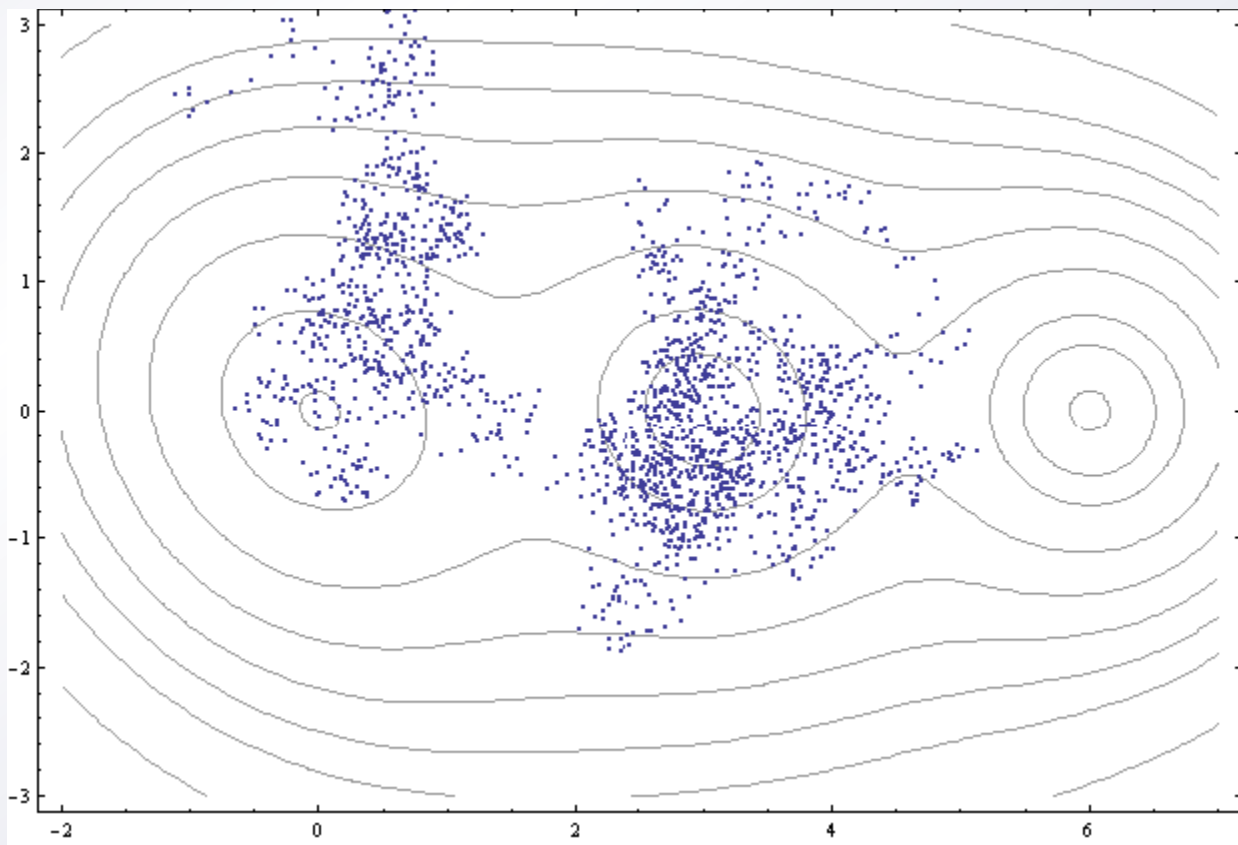
Efficient Markov chains

- The proposal probability distribution $q(\mathbf{y}|\mathbf{x})$ is crucial to sample the space in the shortest time.
- It is forbidden to change it during the Markov chain.
- We can use a uniform distribution centered on \mathbf{x} within some ranges, a multivariate gaussian or similar.
- The size in each direction can be adapted using the local gradient at the initial conditions.
- A too large $q(\mathbf{y}|\mathbf{x})$ will generate very unlikely proposals
- A too small $q(\mathbf{y}|\mathbf{x})$ will only sample locally and never reach convergence.
- The **acceptance rate** should be in the range $[0.2, 0.6]$, with a preference for smaller values at large dimensions.
(0.23 is optimal for infinite dimensions)



Convergence

- Markov chains have the ability of jumping out of local minima.
- A Markov chain has converged if, divided into several chunks, each chunk represents a sampling of the same distribution.
- Convergence tests include autocorrelation measures or correlations among several independent chains.

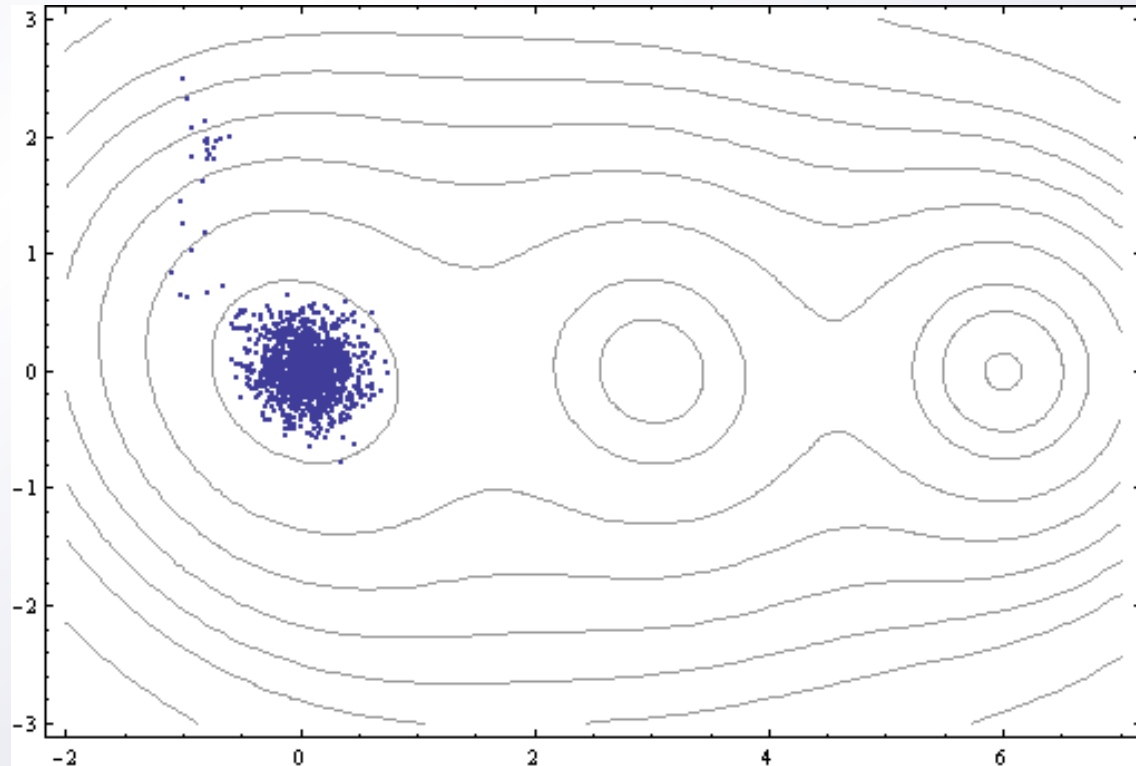


Simulated annealing

- Let us introduce the “**temperature**” T , modifying the probability:

$$p = \exp\left(-\frac{\chi^2}{2T}\right)$$

- At high temperature, all probability ratios tend to 1 and the Markov chain is free to move everywhere.
- The idea (Kirkpatrick et al. 1983) is to start at high temperature to explore the whole parameter space and gradually lower the temperature to pinpoint the best model.



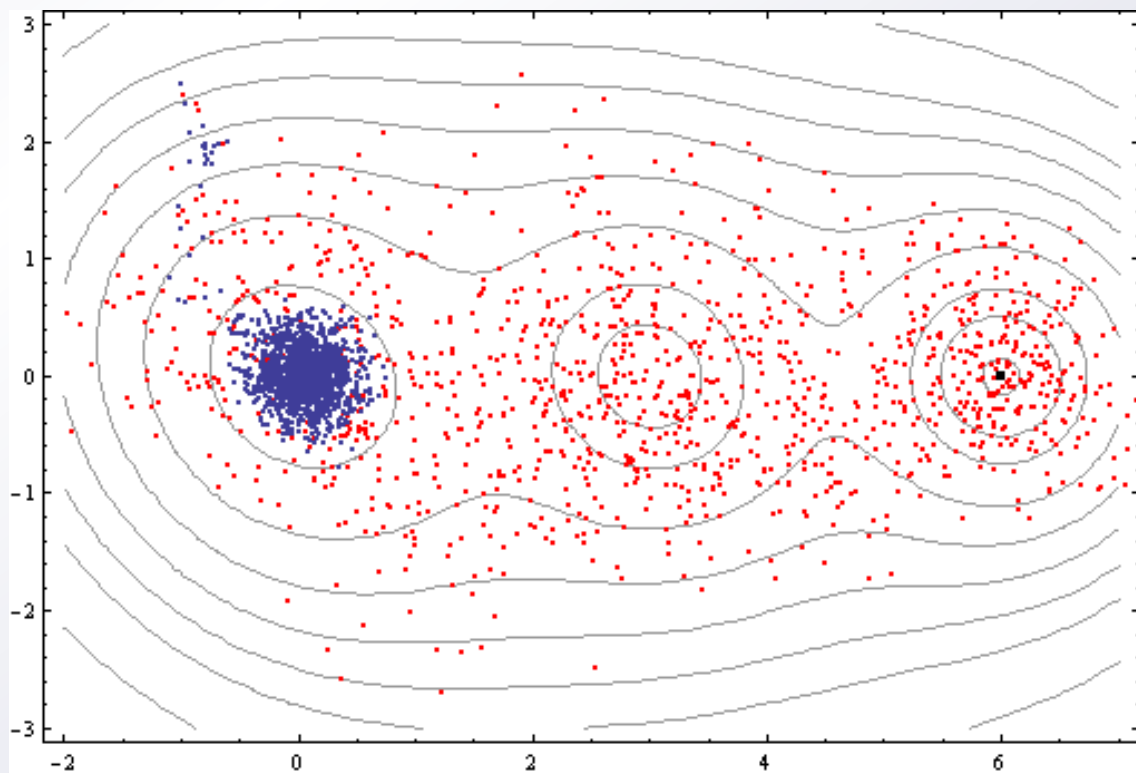
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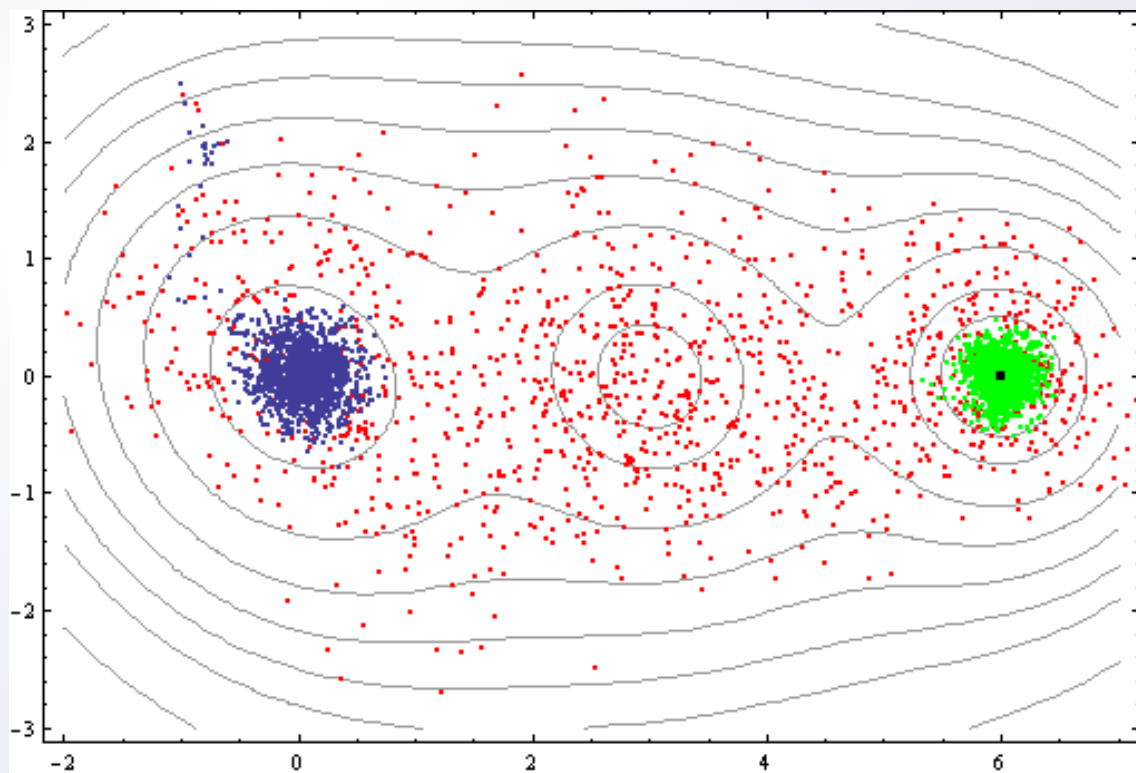
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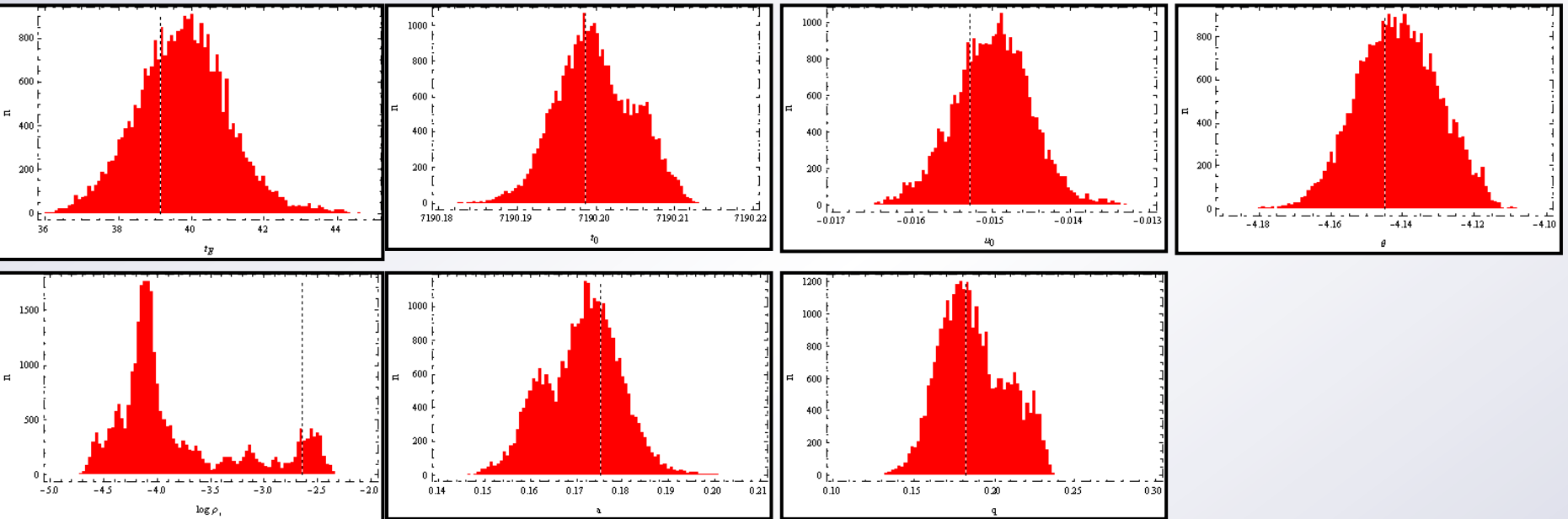
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2. Uncertainty assessment

Confidence intervals

- Once we have sampled our likelihood, we can build **histograms** on any parameters.

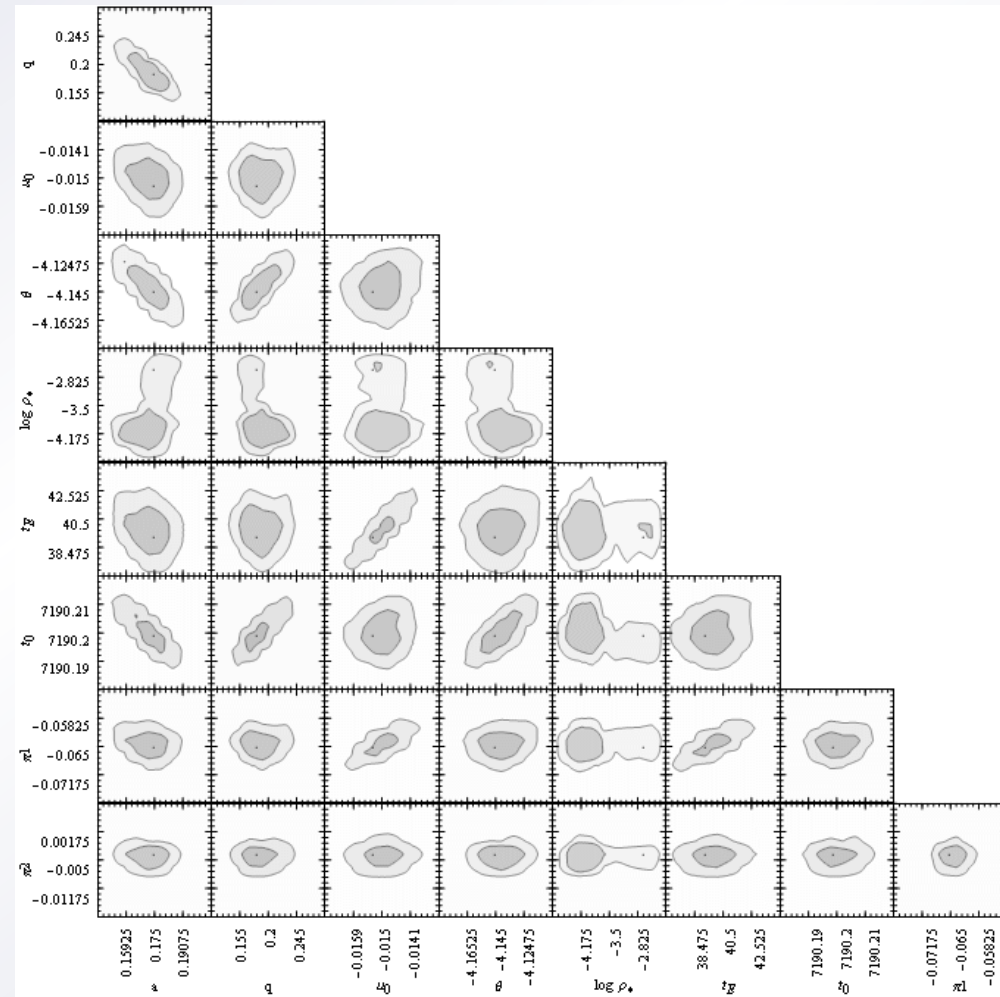


- Confidence intervals** can be obtained:
 - Sort bins according to their height
 - Retain higher bins until you reach the desired CL (e.g. 90%)
 - The CL range is then given by the positions of the two farthest bins on left and right.

2. Uncertainty assessment

Correlation plots

- We can produce density plots on planes defined by any pair of parameters.
- We can define confidence **contours** in the same way.
- This is useful to visualize and detect **degeneracies**.



2. Uncertainty assessment

Fisher and covariance matrices

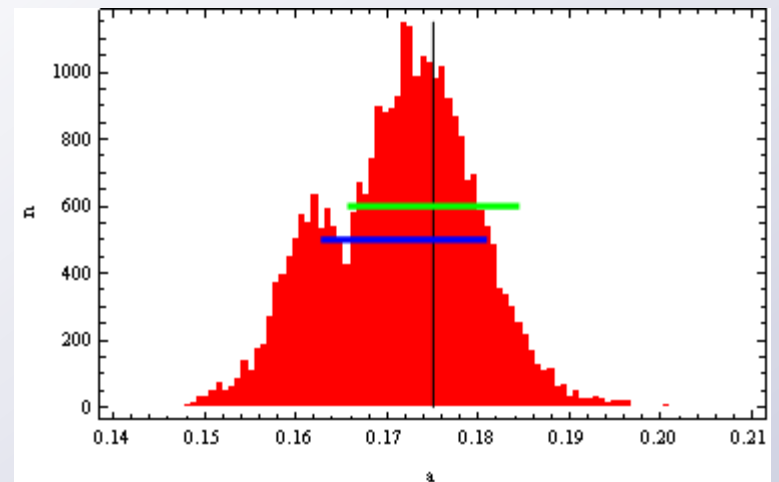
- A common misconception is that MCMC is the only way to obtain the uncertainties in our parameter estimates.
- If you get the best model from other algorithms (e.g. LM), the shape of the minimum is obtained by the **Fisher** matrix

$$F_{mn} = \sum_i \frac{1}{\sigma_i^2} \frac{\partial f(t_i; \mathbf{p})}{\partial p_m} \frac{\partial f(t_i; \mathbf{p})}{\partial p_n}$$

- The **covariance** matrix is just the inverse of the Fisher matrix

$$\text{COV}_{mn} = \left(F^{-1} \right)_{mn}$$

- The variance of each parameter is read along the diagonal of this matrix.



3. Degeneracies

Degeneracies in microlensing

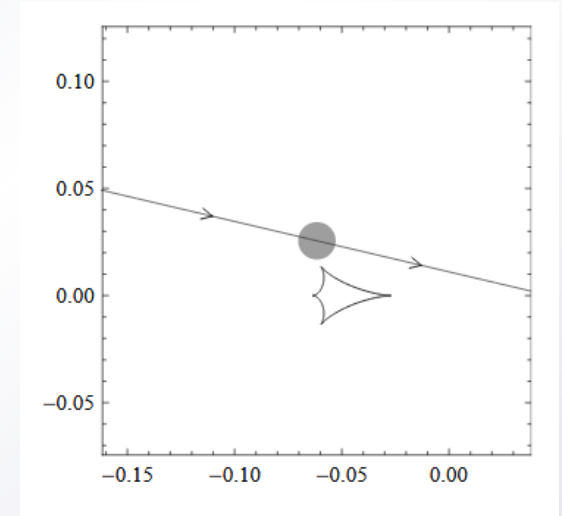
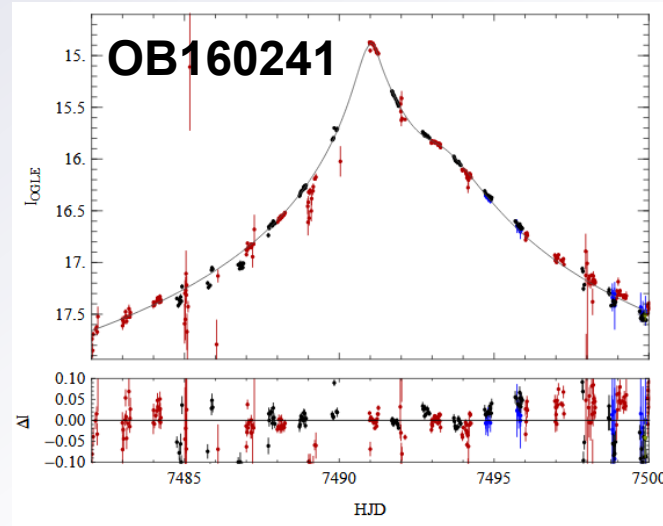
- A degeneracy exists when the same data can be explained by many **different models** with the same likelihood.
- We can have **continuous** degeneracies (e.g. q/s)
- ... or **discrete** degeneracies (e.g. wide/close)

- Degeneracies can be “**strong**” i.e. inherent to gravitational lensing physics itself,
- ... or “**accidental**” if they arise only because of observational shortcomings (gaps, poor sampling, noise, systematics).

3. Degeneracies

Discrete degeneracies

- **Close/Wide** degeneracy in **planets**

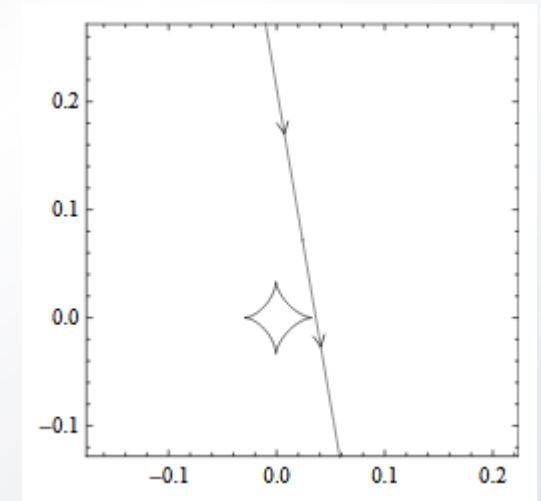
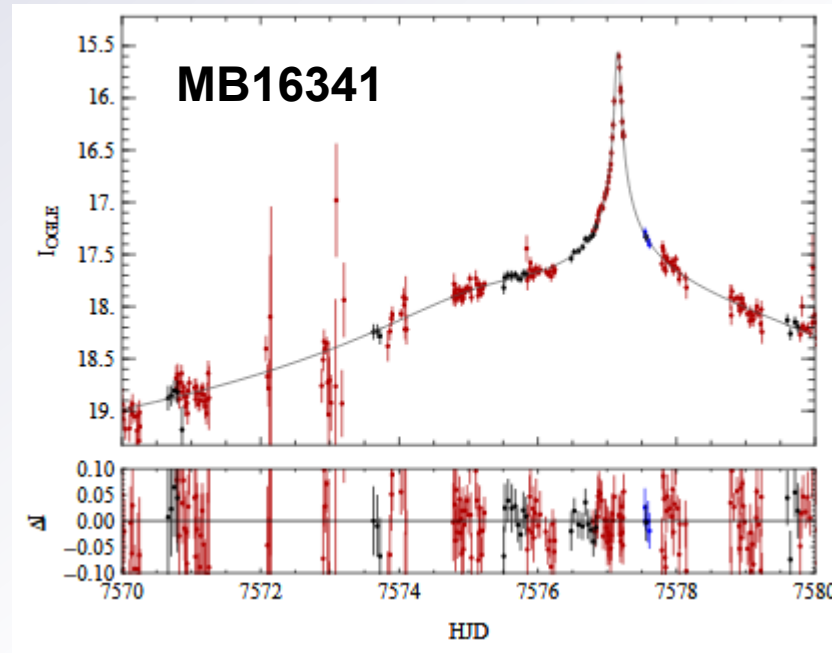


- The central caustic is **invariant** under the transformation $S \leftrightarrow \frac{1}{S}$
- All planetary perturbations due to the central caustic suffer from this degeneracy.

3. Degeneracies

Discrete degeneracies

- **Close/Wide degeneracy in binaries**

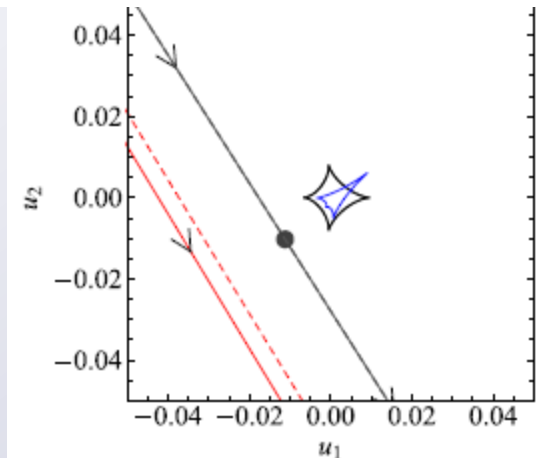
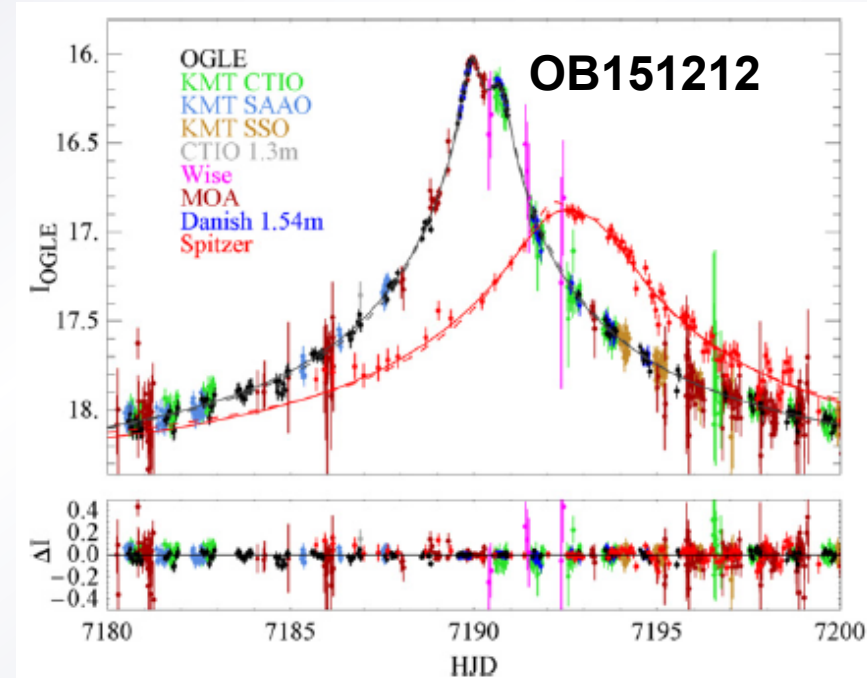


- The **Chang & Refsdal** caustic in the wide regime and the **quadrupole** caustic in the close regime are very similar.
- In addition, the four cusps of a Chang & Refsdal are equivalent (4 possible sub-cases).

3. Degeneracies

Discrete degeneracies

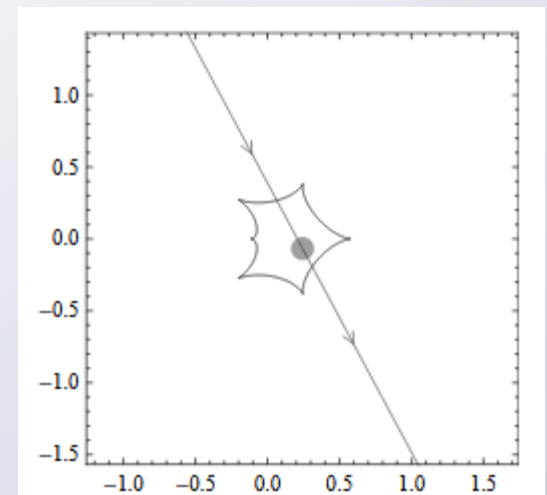
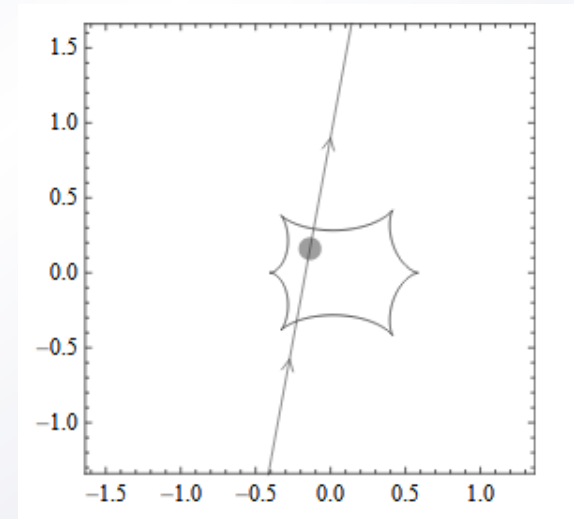
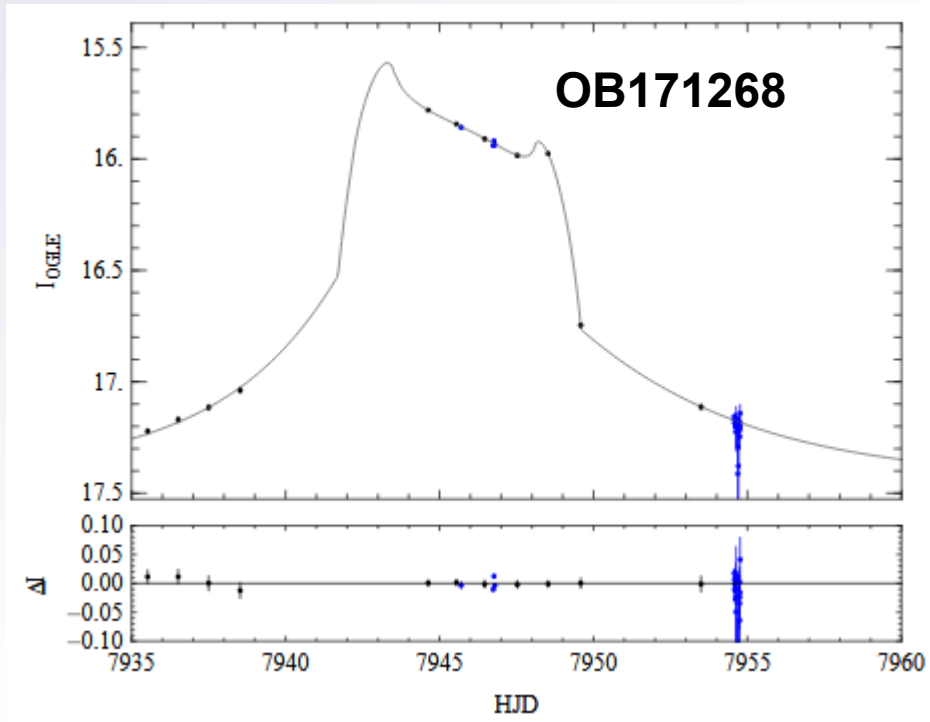
- **Han & Gaudi (2008)** degeneracy
- Light curves with a double peak can be explained by a close approach to a Chang & Refsdal astroidal caustic
- ... or by the approach to the back of a central caustic in the planetary regime.
- In either case we have the close/wide sub-cases and all possible cusp approaches for the binary.



3. Degeneracies

Discrete degeneracies

- **Intermediate** binary degeneracies
The intermediate binary caustic is very extended. Trajectories crossing different folds may lead to very similar light curves

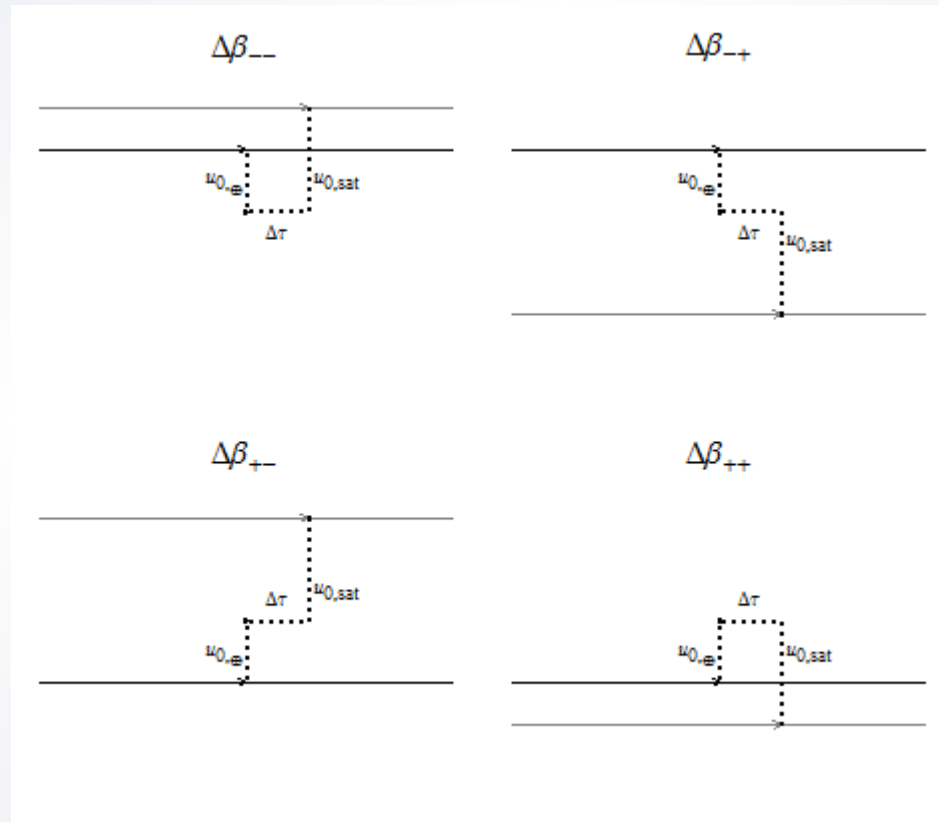


3. Degeneracies

Discrete degeneracies

- **Satellite** degeneracy

Similarly to what happens for PSPL events, if we have observations from space, we have four options for the signs of $u_{0,\text{Earth}}$ and $u_{0,\text{satellite}}$



3. Degeneracies

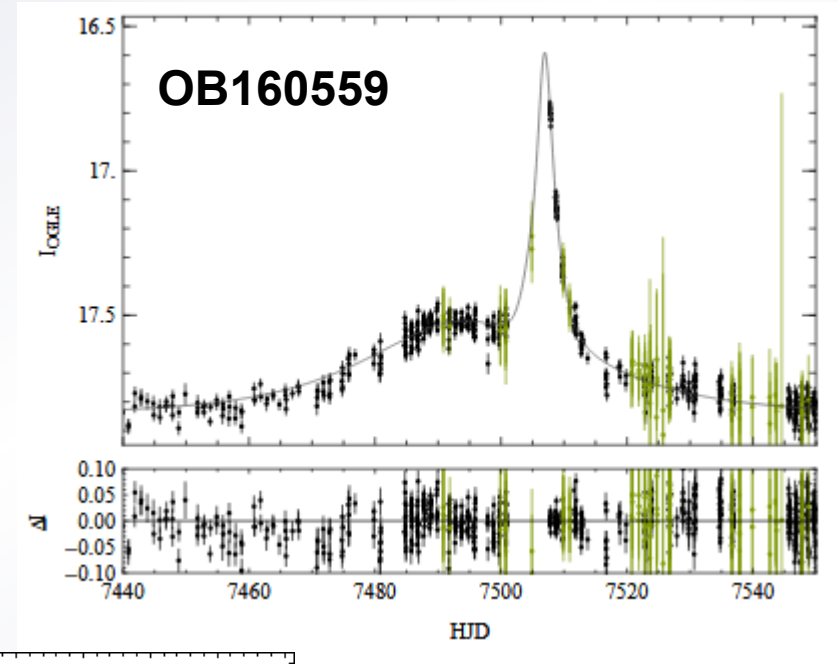
Continuous degeneracies

- **s/q** degeneracy

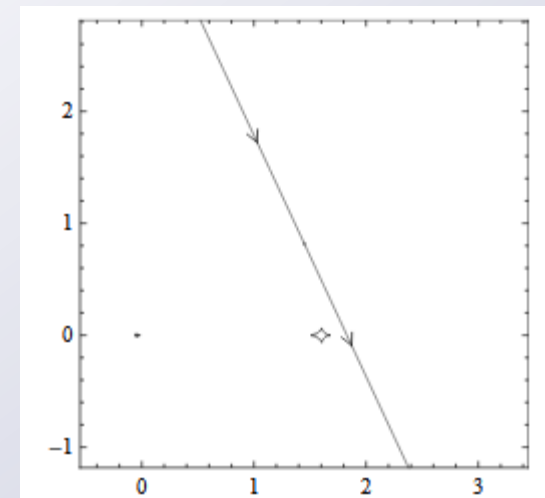
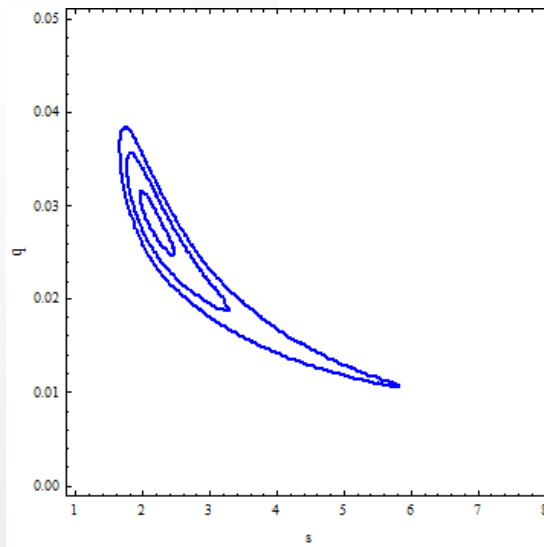
The size of an astroid caustic depends on the combinations

$$\frac{1}{s^2} \sqrt{\frac{q}{(1+q)^3}} \quad \text{Wide regime}$$

$$s^2 \frac{q}{(1+q)^2} \quad \text{Close regime}$$



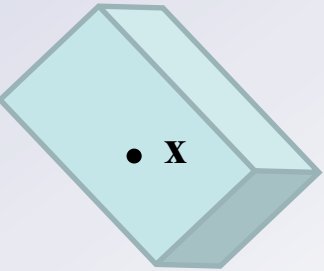
- The mass ratio and separation are highly correlated and poorly known.



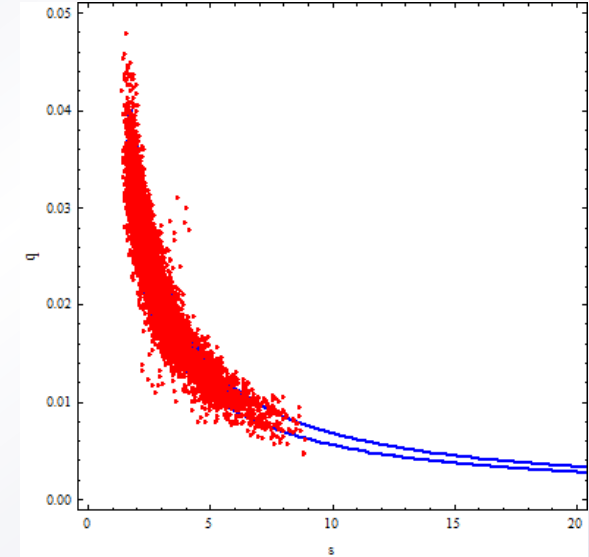
3. Degeneracies

Alternative parameterizations

- The **exploration** of continuous degeneracies is particularly **painful**.



- We can rotate the box of the proposal distribution (easily achieved if we diagonalize the local Fisher matrix before starting the chain)



- If the degeneracy is non-linear, choose **new parameters**

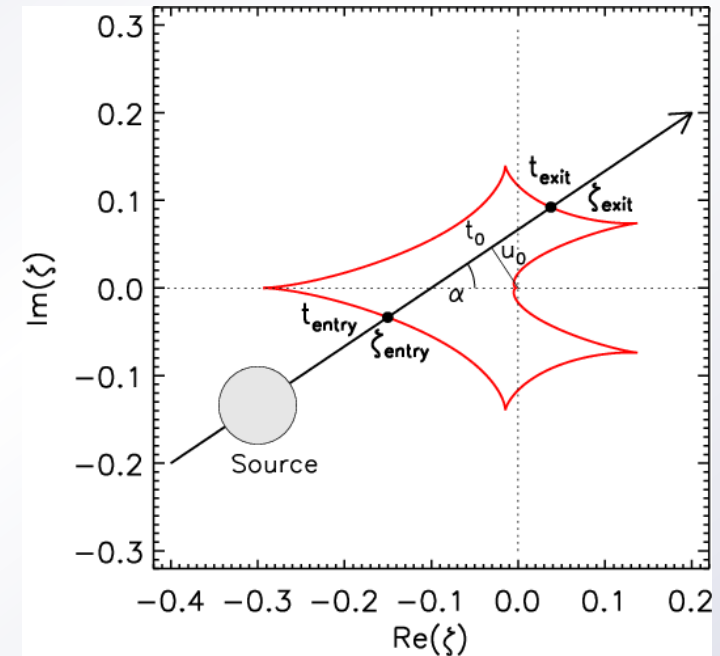
$$p_1 = \frac{1}{s^2} \sqrt{\frac{q}{(1+q)^3}}; \quad p_2 = s^2 \sqrt{\frac{q}{(1+q)^3}}$$

- Define the origin at the center of the caustic we wish to study.
- Fit the log of some parameters (s, q, ρ, t_E).
- Use other combinations that are clearly established by the data (e.g. source crossing time t_* , time of caustic crossing, ...)

3. Degeneracies

Cassan parameters

- Cassan (2008) proposed to use the **curvilinear abscissa** along the caustic.
- u_0 , α , t_0 , t_E are replaced by t_{entry} , s_{entry} , t_{exit} , s_{exit} .



3. Degeneracies

Walking through degeneracies

- In general, distinct features in the lightcurve occurring at definite times (caustic crossing) **couple** the **Einstein time** to the (s,q) values.
- This mitigates the degeneracy between t_E and u_0 plaguing the PSPL.
- For the same reason, different models may predict very different t_E and thus very different **blending** ratios.
- Typically, planetary models mimicking binary models come at negative blending.
- Other hints may come from unlikely source radii or unlikely Einstein times.
- How do we quantify unlikeliness?

4. Bayesian analysis

Bayes' theorem

- For all parameters we can define an expected range of possible values.
- A **uniform prior** can be easily implemented by requiring that the proposal point is within the prior.
- However, we may wish to use the information coming from previous studies to decide which model is more likely (stellar luminosity and mass functions, spatial distributions and velocities)
- This information typically comes in the form of (**prior distributions**).
- Bayes theorem states that the **posterior probability** is the product of the likelihood from the data with the prior expectations:

$$p(\mathbf{p} | y_i) = \frac{p(y_i | \mathbf{p})p(\mathbf{p})}{p(y_i)}$$

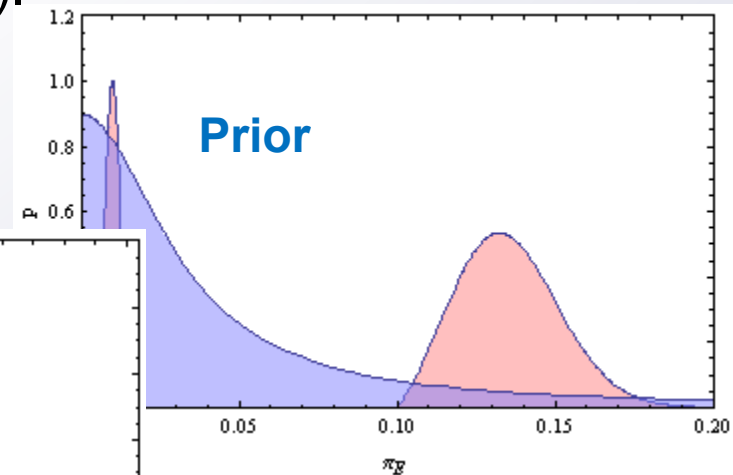
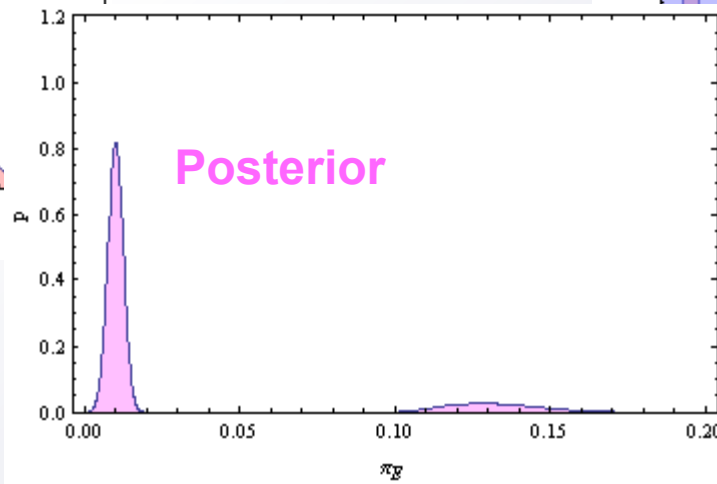
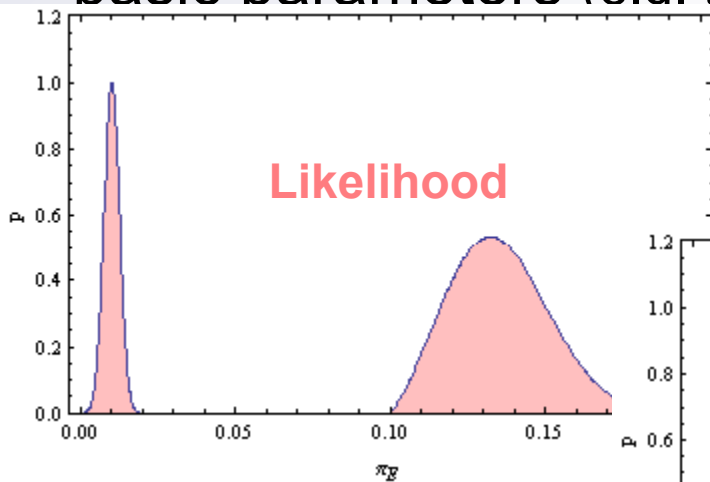
4. Bayesian analysis

Bayes in MCMC

$$p(\mathbf{p} | y_i) = \frac{p(y_i | \mathbf{p})p(\mathbf{p})}{p(y_i)}$$

- In our MCMC we just have to sample the product $\exp\left(-\frac{\chi^2}{2}\right)p(\mathbf{p})$

- The **normalization** $p(y_i)$ cancels if we are only interested in relative posterior probabilities (ratios).
- Note that the priors may be distributions on combinations of the basic parameters (e.g. the mass of the lens).



4. Bayesian analysis

Priors in microlensing

- Microlensing events are normally occurring on source stars in the bulge lensed by stars in the disk/bulge.
- These objects follow some spatial distributions, mass, luminosity and velocity functions.
- In order to use Bayesian approach in microlensing we need a **Galactic model**.
- “Stochastic distributions of lens and source properties for observed galactic microlensing events”, Dominik (2006).
- “A synthetic view on structure and evolution of the Milky Way”, Robin et al. (2003) (Besançon model)
- “Stellar Contribution to the Galactic Bulge Microlensing Optical Depth”, Han & Gould (2003)
- Another combination of models is in Bennett et al. (2008)
- **Blending** light gives a further constraint (see Beaulieu’s talk).

5. Initial conditions

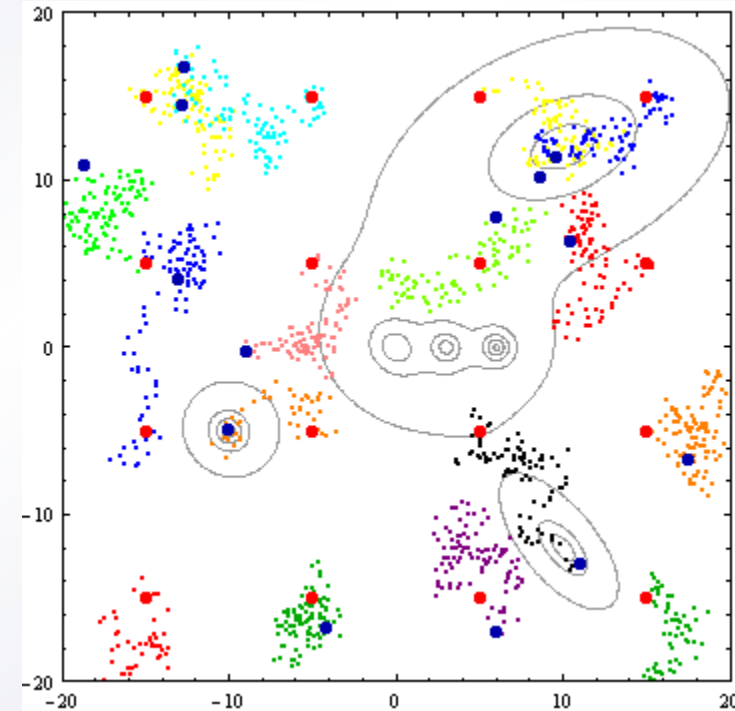
Initial conditions

- Microlensing **parameter space** is huge and full with local χ^2 minima.
- If we start from an arbitrary initial condition we would seldom end in the global minimum.
- We need to **explore** all the relevant parameter space and make sure we find the true best model(s).
- Two ways:
 - Grid search
 - Template library

5. Initial conditions

Grid search

- We may define a **grid** in the parameter space and start fits from all points.
- Many fits will just never converge
- Many fits will end up in the same minima.
- Many minima will be missed.

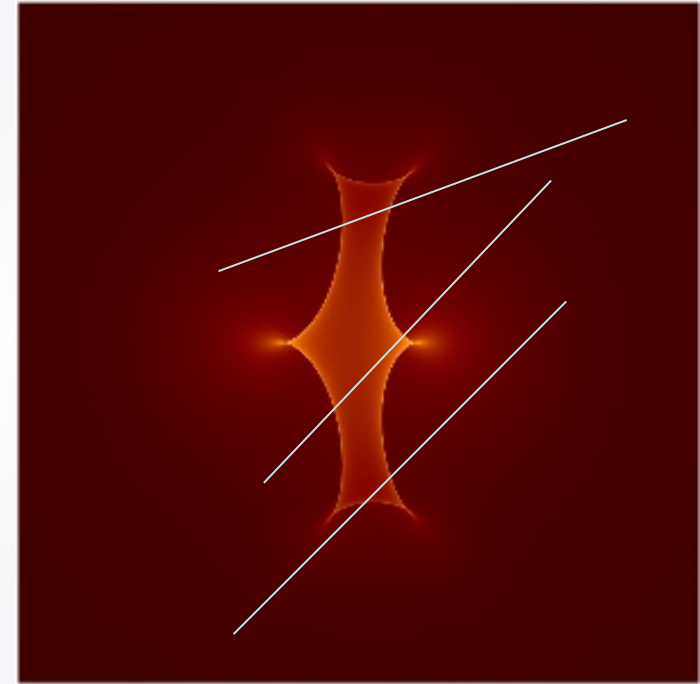
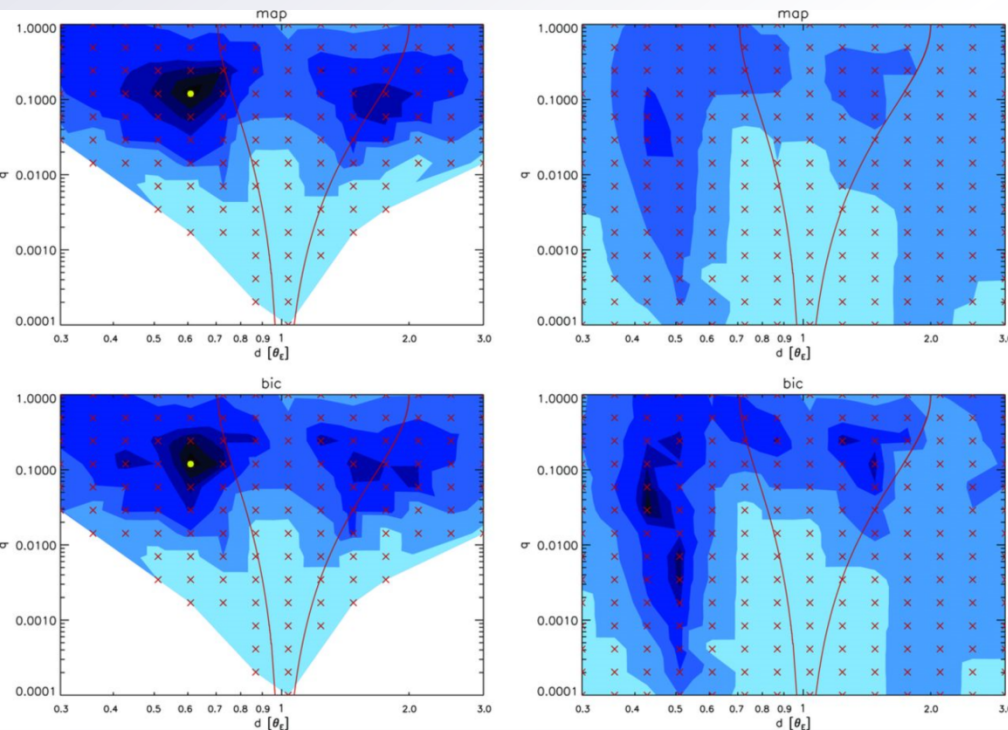


- A **too coarse** grid may miss possible candidate models.
- A **too dense** grid has many redundant or useless fits.

5. Initial conditions

Two-steps grid search

- Inverse-ray-shooting codes may keep **(s,q)** fixed in a first search, so as to use the same magnification map.
- Once these preliminary models are found, we can run a **full fit** including **(s,q)**.

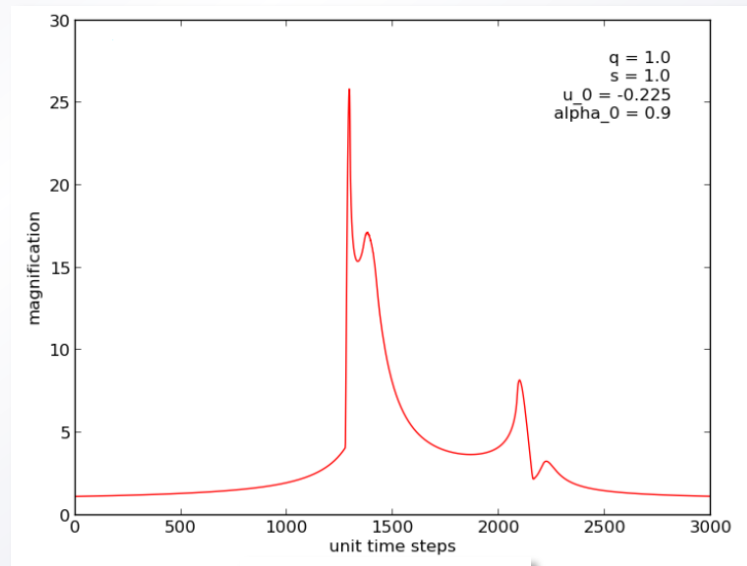


- Codes for a full Bayesian approach along these lines are available (ML/MAP/BIC) (Kains et al. 2012)

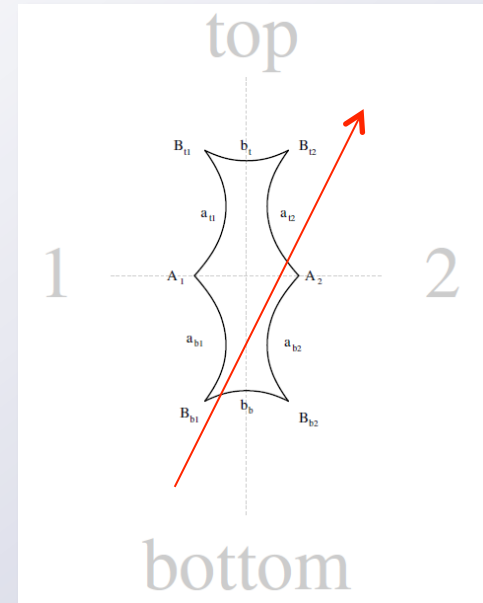
5. Initial conditions

Template libraries

- It would be much more efficient to start the fit from an initial condition that resembles our data.
- We need to build a **library of light curves** covering all possible morphologies (Di Stefano & Mao 1996, Night et al. 2005).
- The most systematic attempt has found **73** different morphologies out of **232** regions in the parameter space (Liebig et al. 2015)
- It was limited to equal-mass binaries!



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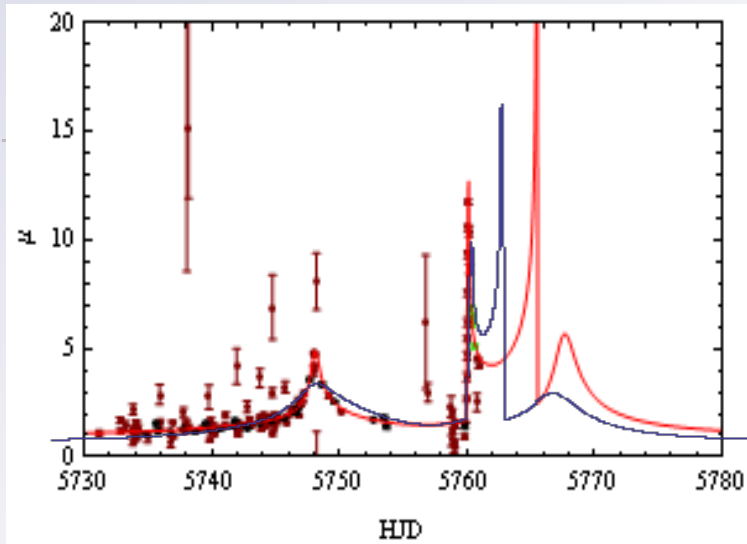


$[b_b a_{b1} a_{t2}] A_2$

5. Initial conditions

Matching a template to data

- **Peaks in the dataset** must be identified and ranked according to their prominence.



- The two most prominent peaks must be **matched** to the most prominent peaks in the template.
- We get $(s, q, u_0, \alpha, \rho)$ from the template.
- (t_0, t_E) are obtained by the peak matching.

- If there is only one peak, the **anomaly time** can be taken as the position of the second peak.

5. Initial conditions

RTModel

- RTModel (<http://www.fisica.unisa.it/GravitationAstrophysics/RTModel.htm>) is an **automatic platform** for real-time modeling.
- It takes data and anomaly alerts from ARTEMiS (<http://www.artemis-uk.org/>).
- It uses matching from a library of 244 templates.
- For each initial condition, the Levenberg-Marquardt fit is repeated five times using the bumpers method.
- The calculation of the magnification is done by VBBinaryLensing.
- All models found are ranked by their χ^2 .
- Duplicates are removed if they fall within the same covariance ellipsoid.
- Models are posted on a public webpage automatically.
- RTModel runs on a 8-core workstation taking 2 hours per event.

Outlook

- Higher order effects (parallax, orbital motion) may dramatically increase the number of light curve morphologies.
- Grid searches in too many dimensions are unfeasible.
- Template libraries require a long construction.
- Similar issues hold for triple and multiple lenses.
- In view of WFIRST, we need to improve our automatic modeling capabilities.
pyLIMA, MulensModel