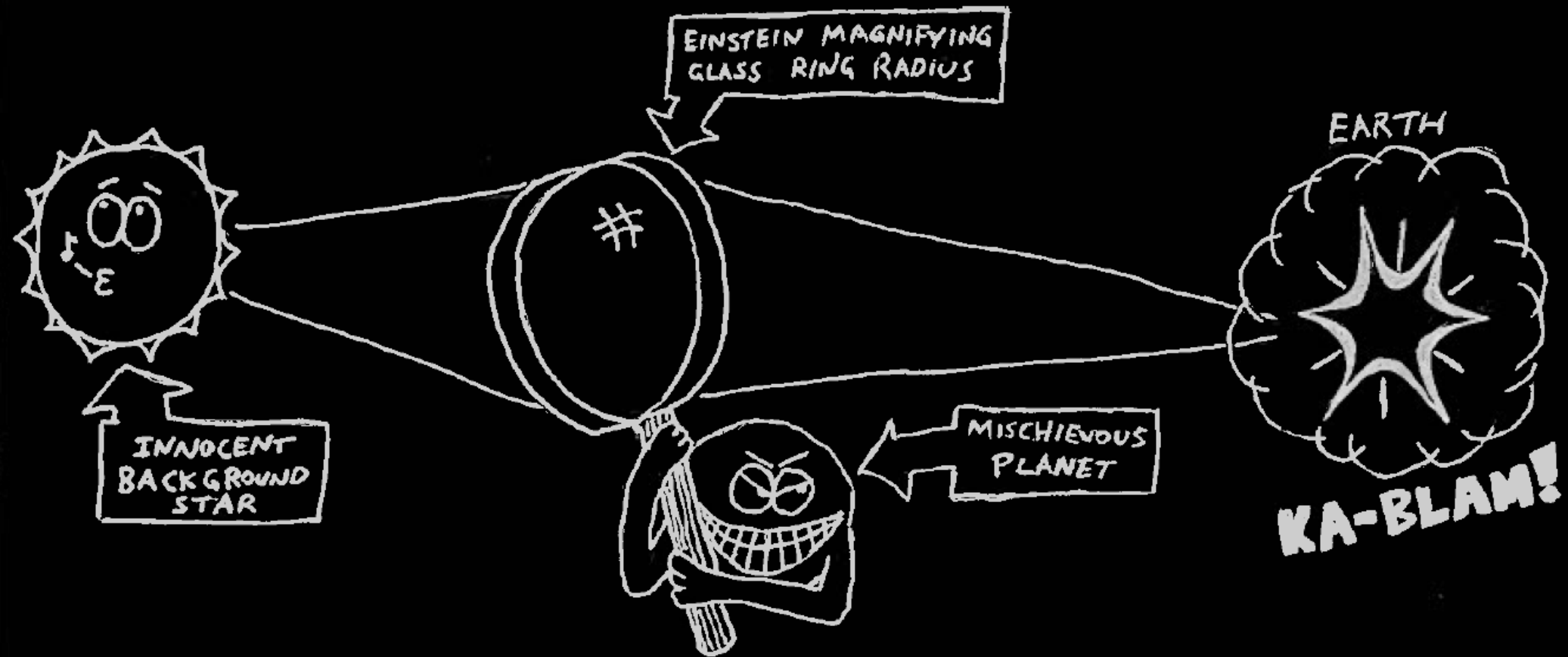


# MICROLENSING MAGNIFICATION CALCULATIONS WITH POINT-SOURCE AND FINITE-SOURCE EFFECTS

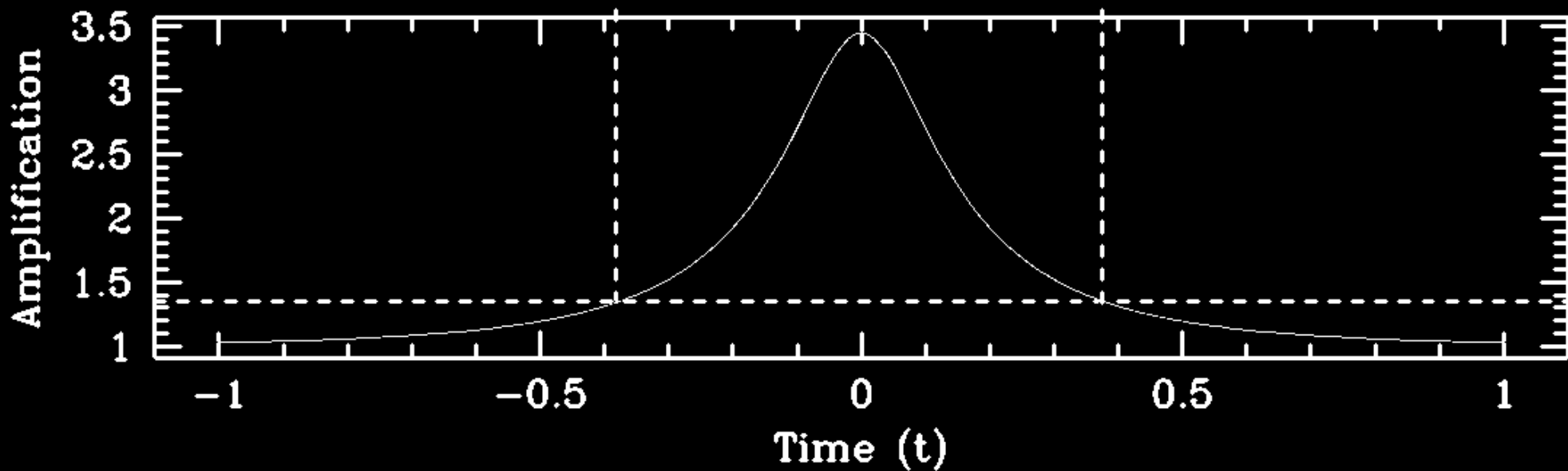
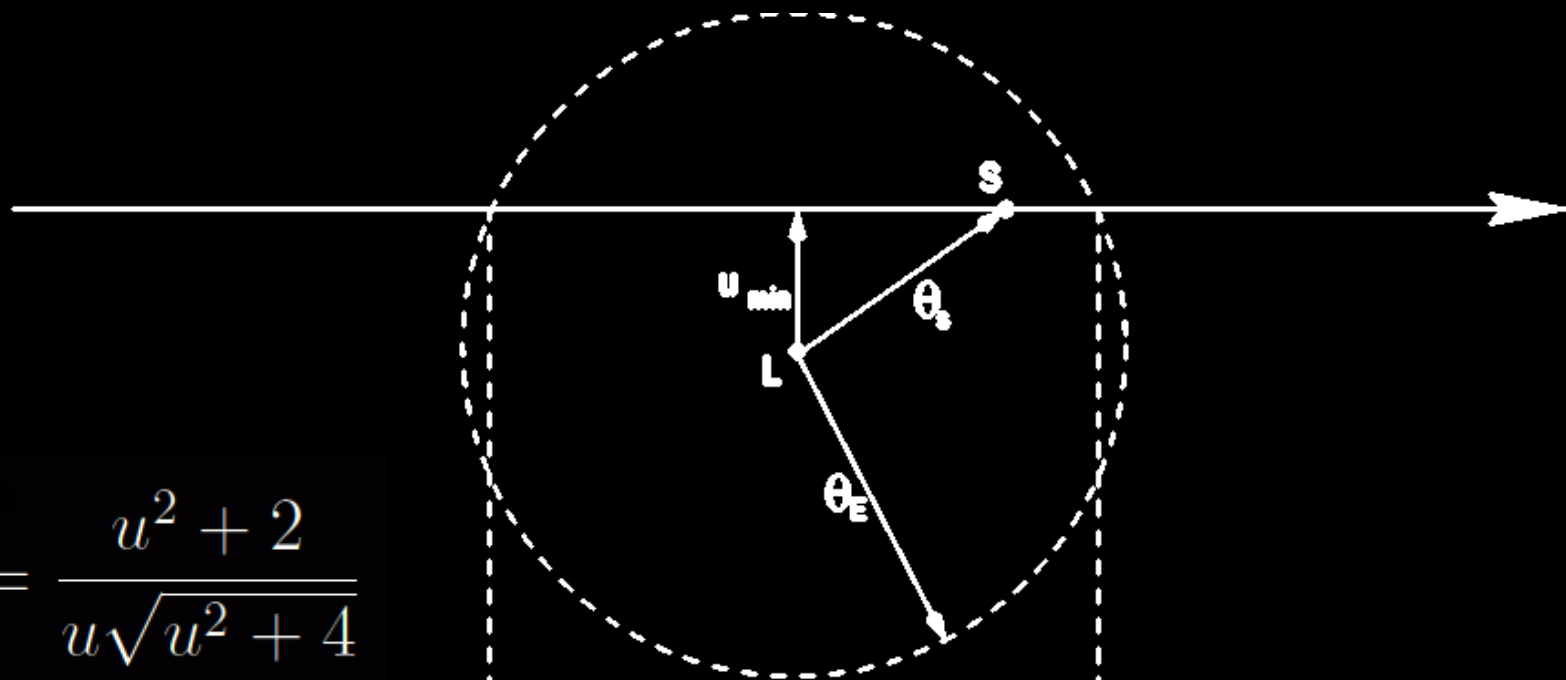
Stephen Kane

NASA Exoplanet Science Institute



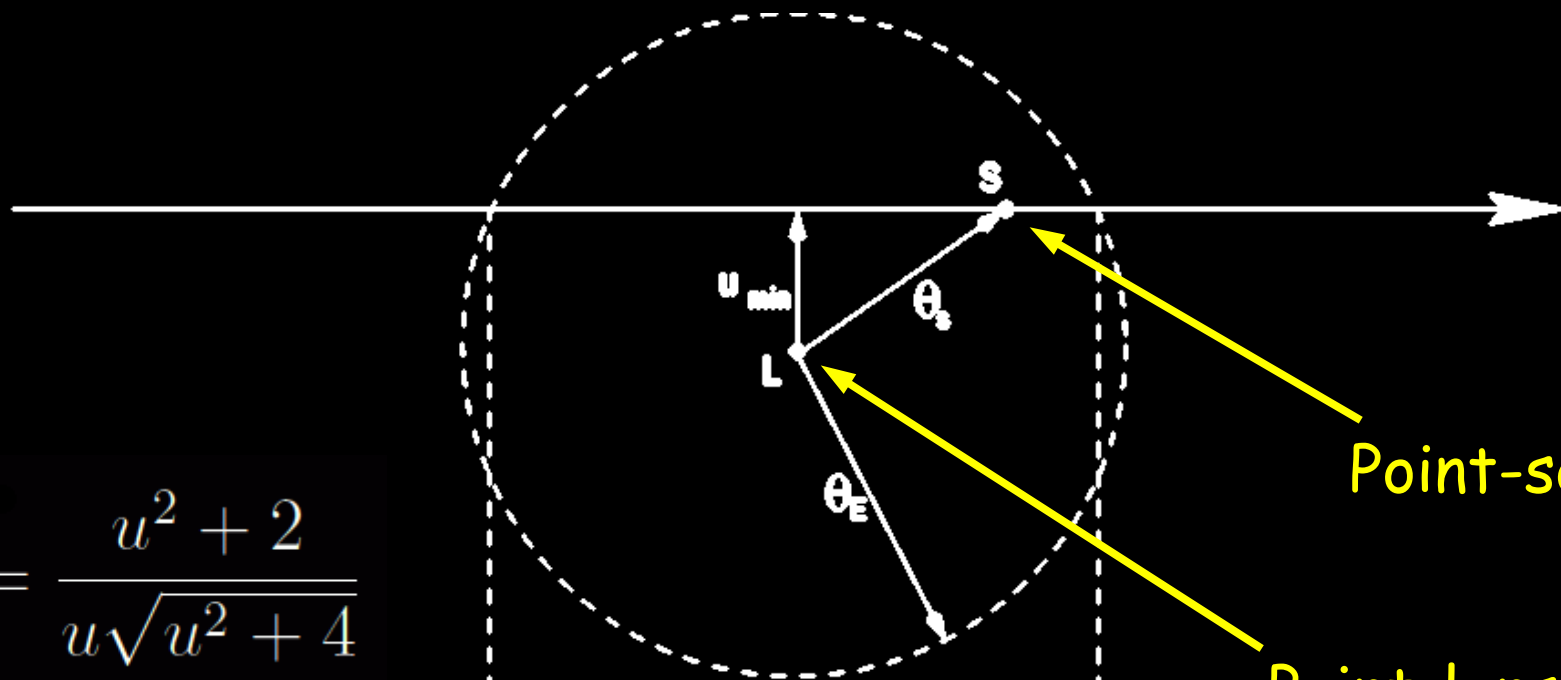
MICROLENSING FOR DUMMIES

$$A = \frac{u^2 + 2}{u\sqrt{u^2 + 4}}$$



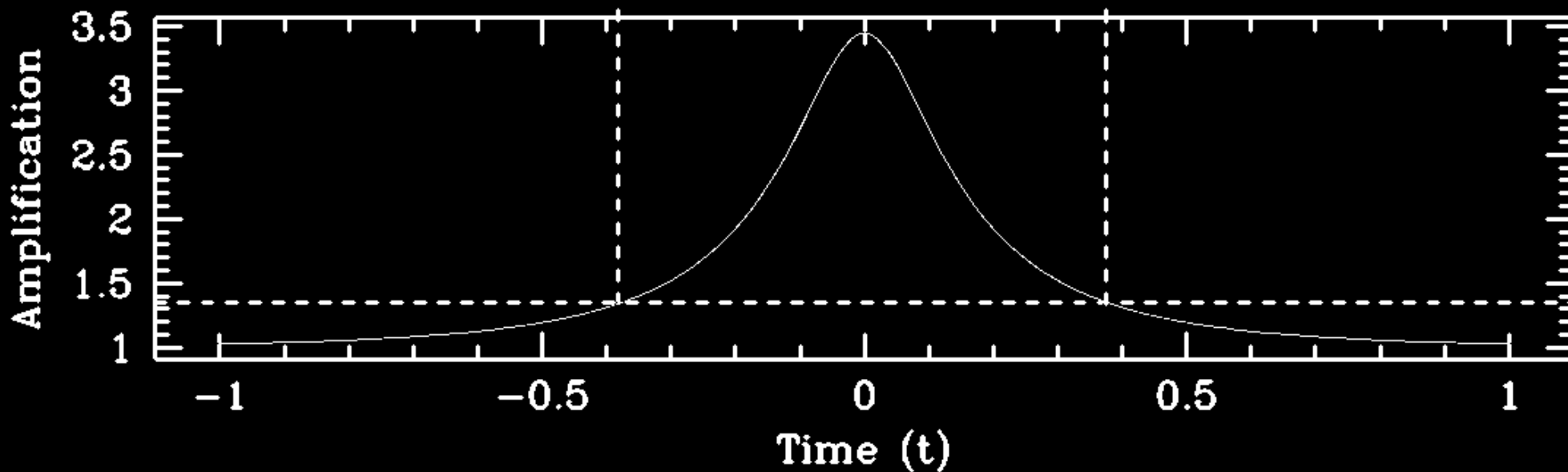
$$A = \frac{u^2 + 2}{u\sqrt{u^2 + 4}}$$

Infinite where  $u = 0$



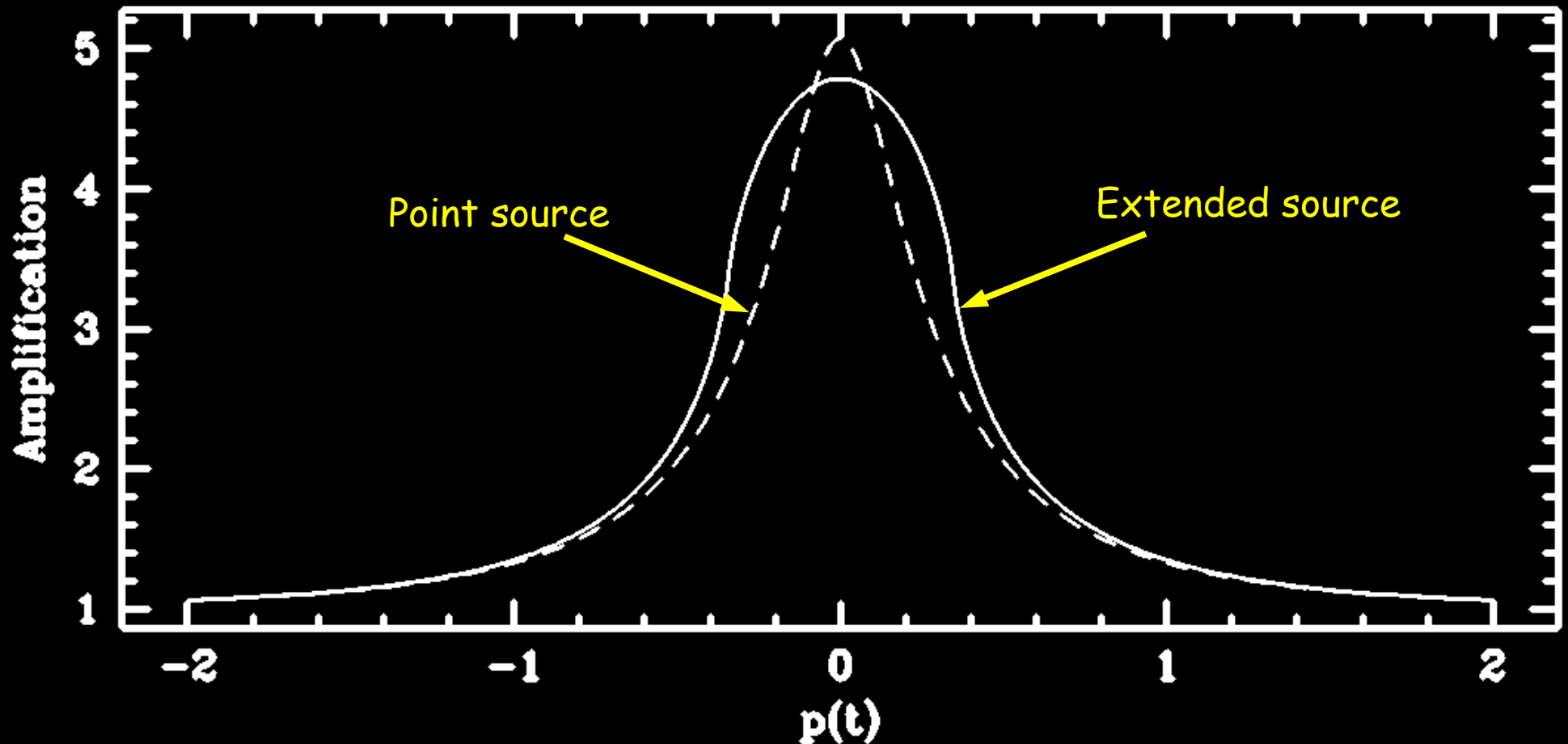
Point-source

Point-lens



# Can we assume a point-source?

- Point-source is often a reasonable assumption
- Assumption breaks down when the source size becomes comparable to the minimum lens-source separation



# Uniform source brightness

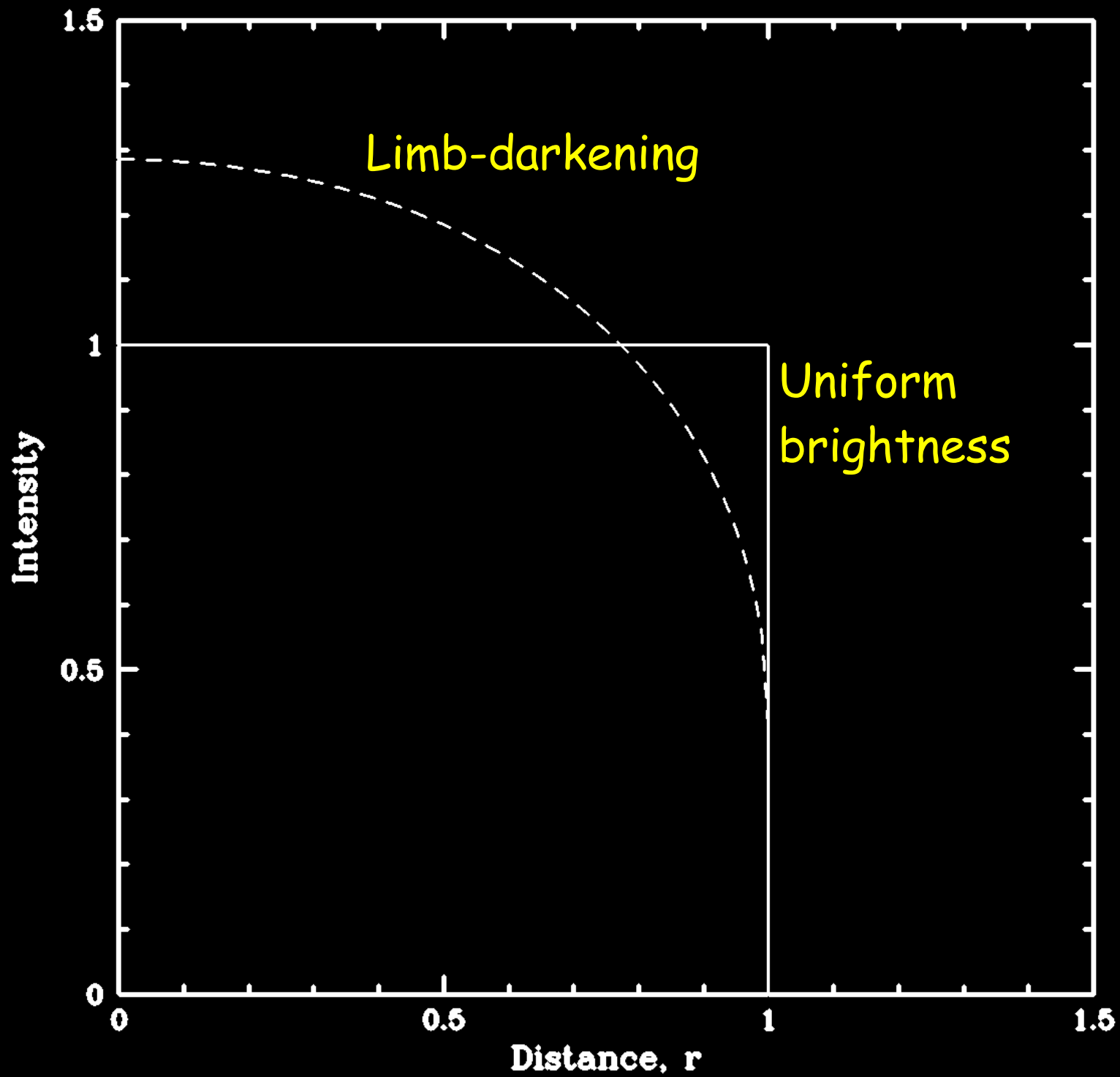
For a source with uniform brightness, the magnification is the ratio of the area of the images to the area of the source.

$$A = \frac{\int d^2y A_p(\mathbf{y})}{\int d^2y}$$

$$A = \begin{cases} \frac{2}{\pi R_s^2} \int_{u-R_s}^{u+R_s} \frac{r^2+2}{\sqrt{r^2+4}} \arccos \frac{u^2+r^2-R_s^2}{2ur} dr & \text{for } u > R_s \\ \frac{2}{\pi R_s^2} \int_{R_s-u}^{R_s+u} \frac{r^2+2}{\sqrt{r^2+4}} \arccos \frac{u^2+r^2-R_s^2}{2ur} + \frac{R_s-u}{R_s^2} \sqrt{(R_s^2-u)^2+4} dr & \text{for } u < R_s \\ \frac{2}{\pi} \left[ \left(1 + \frac{1}{R_s^2}\right) \arcsin \frac{1}{\sqrt{1+\frac{1}{R_s^2}}} + \frac{1}{R_s} \right] & \text{for } u = R_s \end{cases}$$

$$A_{max} = \sqrt{1 + \frac{4}{R_s^2}}$$

Magnification NOT infinite where  $u = 0$



# Limb-darkening

For a source with limb-darkening, different parts of the source will be magnified differently.

$$A = \frac{\int d^2y A_p(\mathbf{y})}{\int d^2y}$$

$$A = \frac{\int d^2y I(\mathbf{y}) A_p(\mathbf{y})}{\int d^2y I(\mathbf{y})}$$

$$I(r) = I(0)[1 - \kappa_1 Y - \kappa_2 Y^2]$$

$$Y = 1 - \sqrt{1 - r^2}$$

# Limb-darkening

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$$I(r) = I(0)[1 - \kappa_1 Y - \kappa_2 Y^2]$$

$$Y = 1 - \sqrt{1 - r^2}$$

**This is not trivially solved for most binary lens events!**

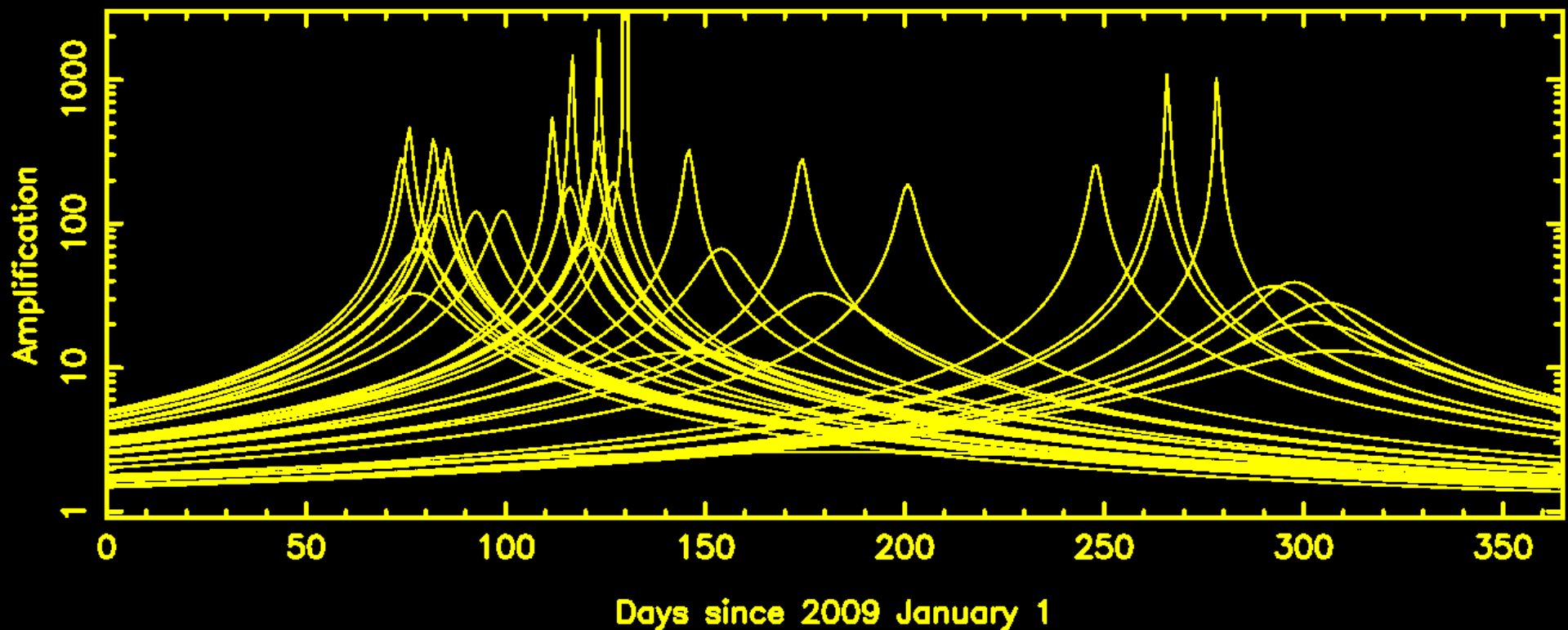


# Can we assume a point-source?

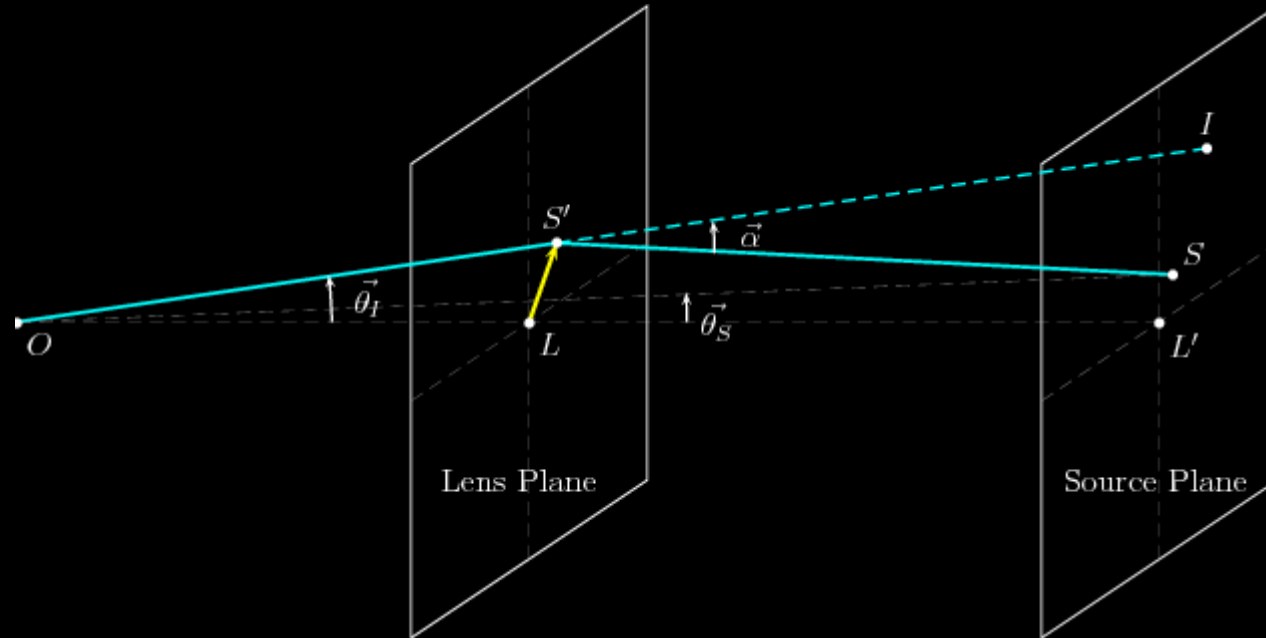
- Well it depends ...
- Assumption breaks down when the source size becomes comparable to the minimum lens-source separation (or planetary Einstein ring)

# Can we assume a point-source?

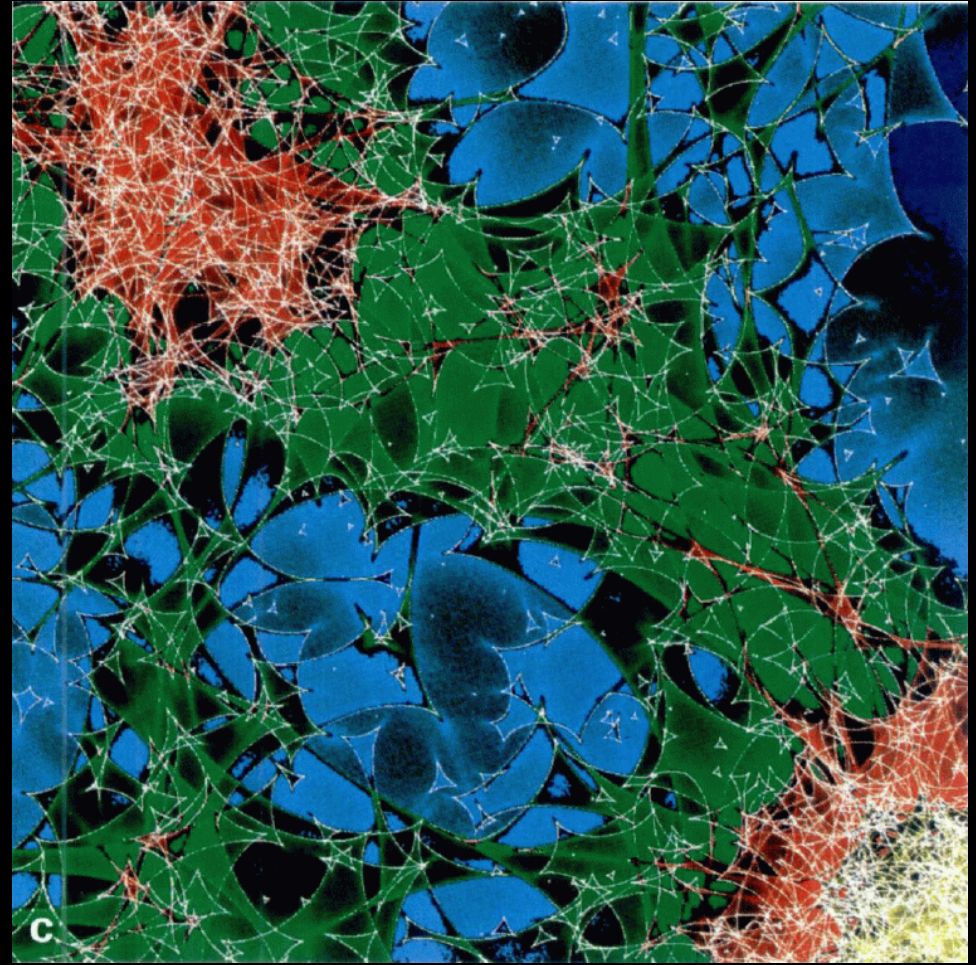
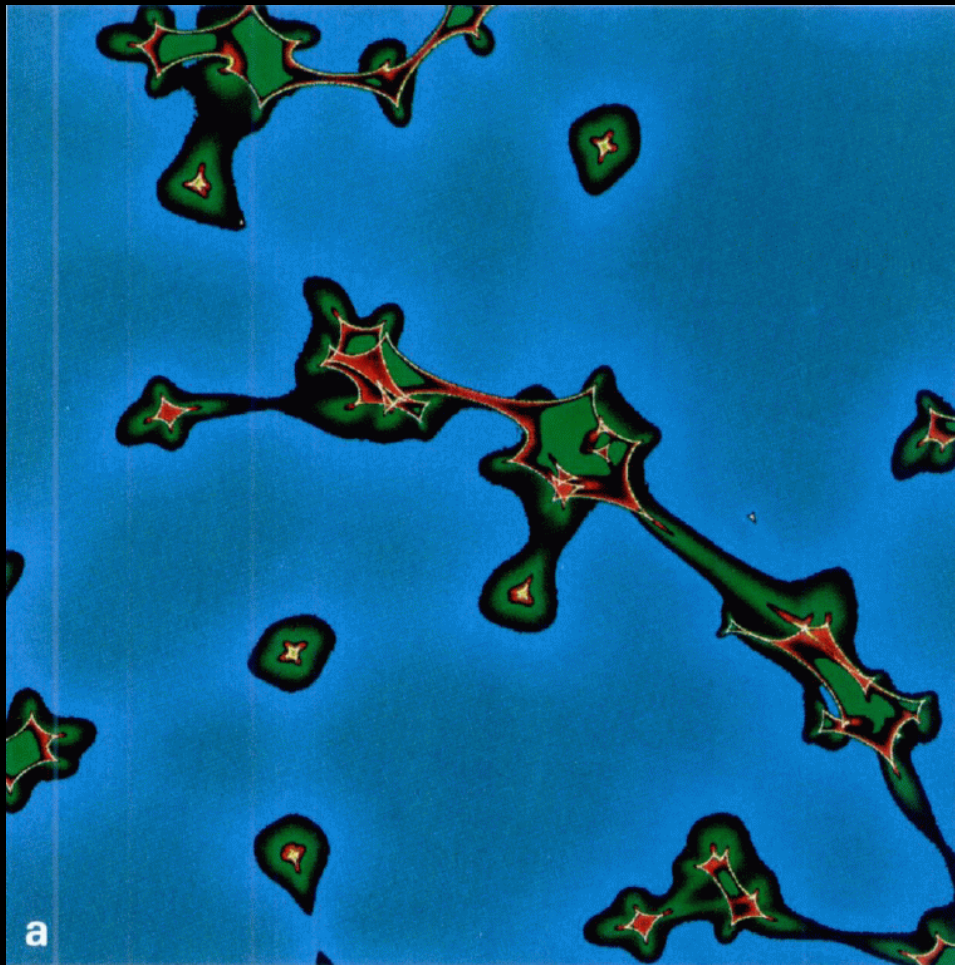
- Well it depends ...
- Assumption breaks down when the source size becomes comparable to the minimum lens-source separation (or planetary Einstein ring)
- The event-rate is increasing, including that of high-magnification events!
- Need to find an optimal path for efficient computational handling of so many events with such a large parameter space



# Magnification Map (ray-shooting)



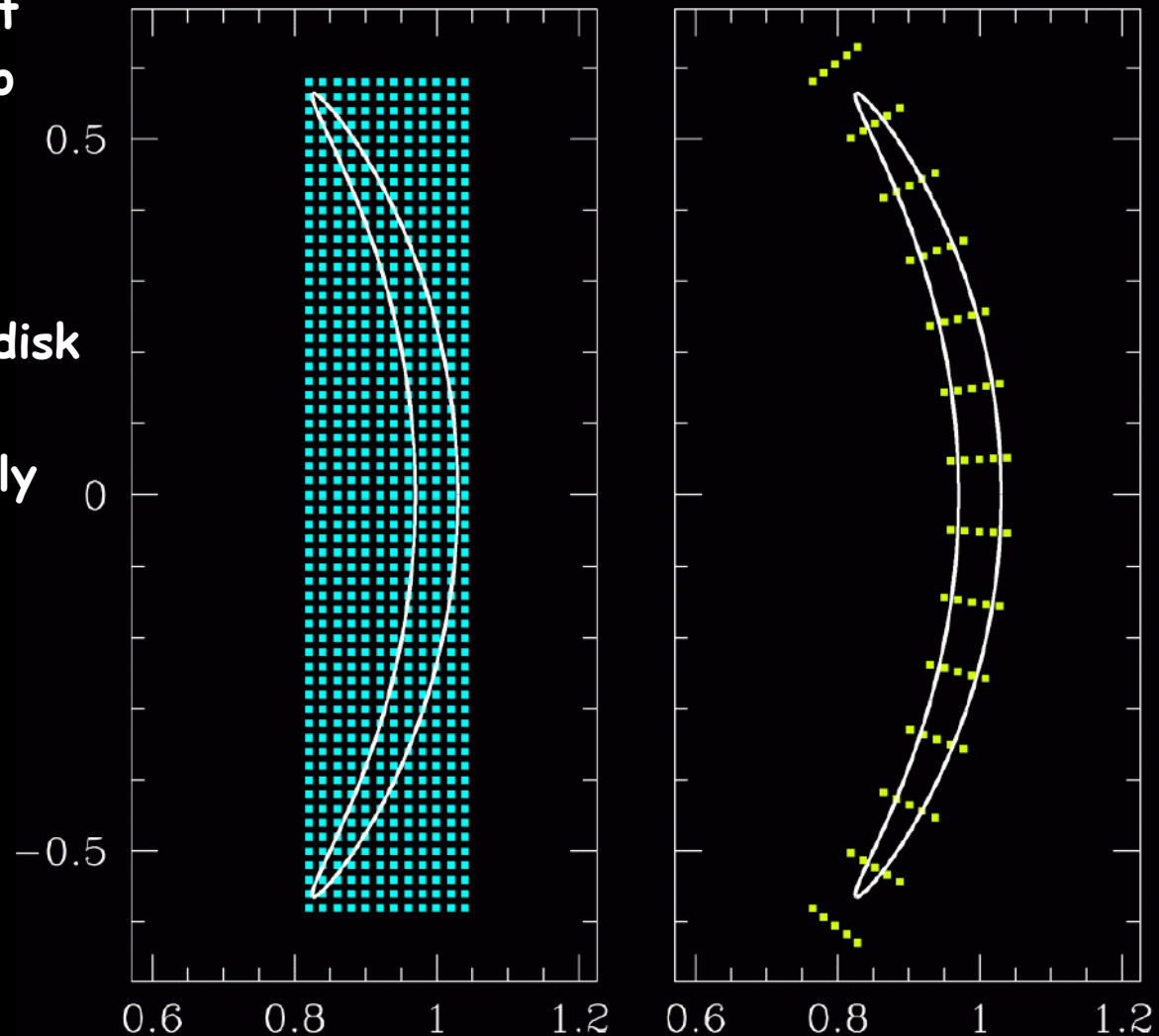
- Brute-force approach
- Handles complicated including multiple lens systems very well (frequently used in quasar lensing modeling)
- Good usage example is Wambsganss, 1997, MNRAS, 284, 172
- Can be computationally expensive – particularly where orbital motion is included!

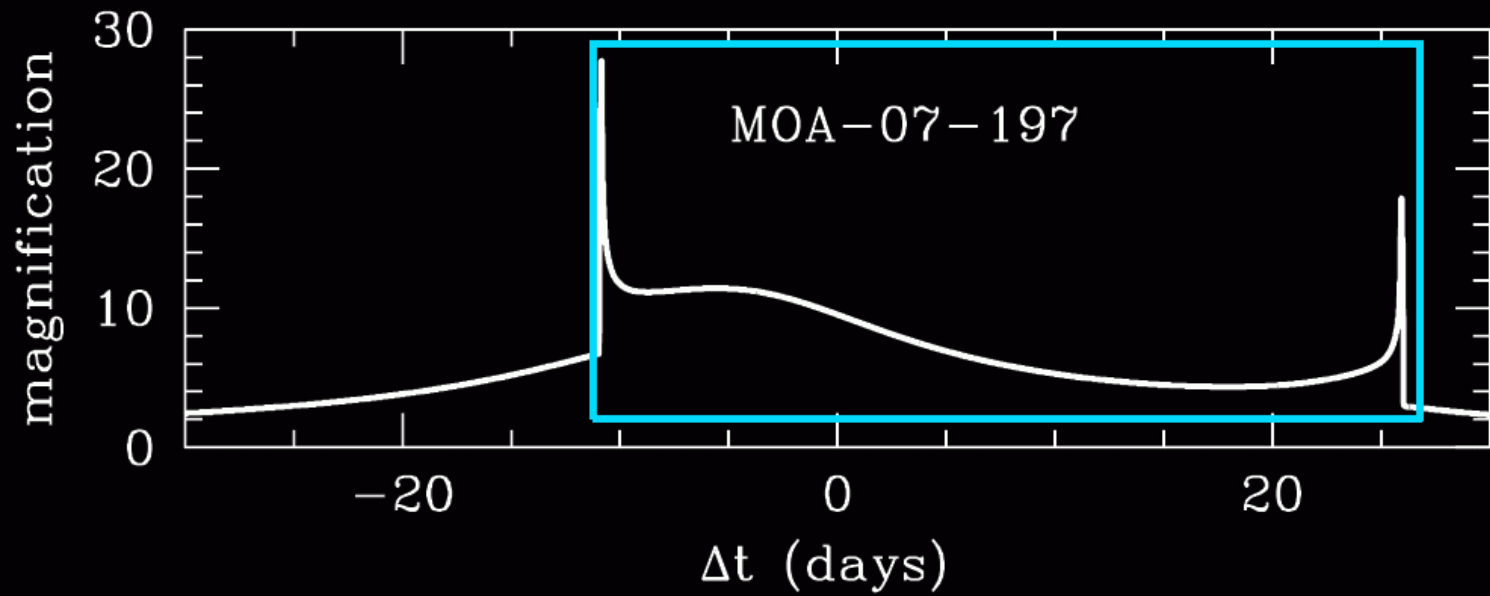
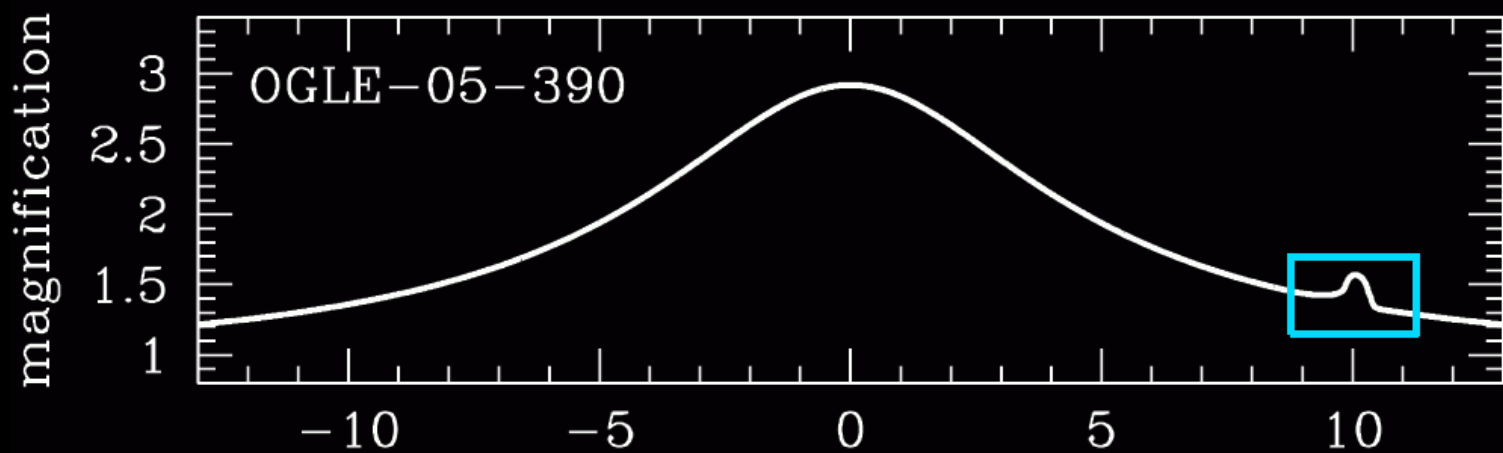
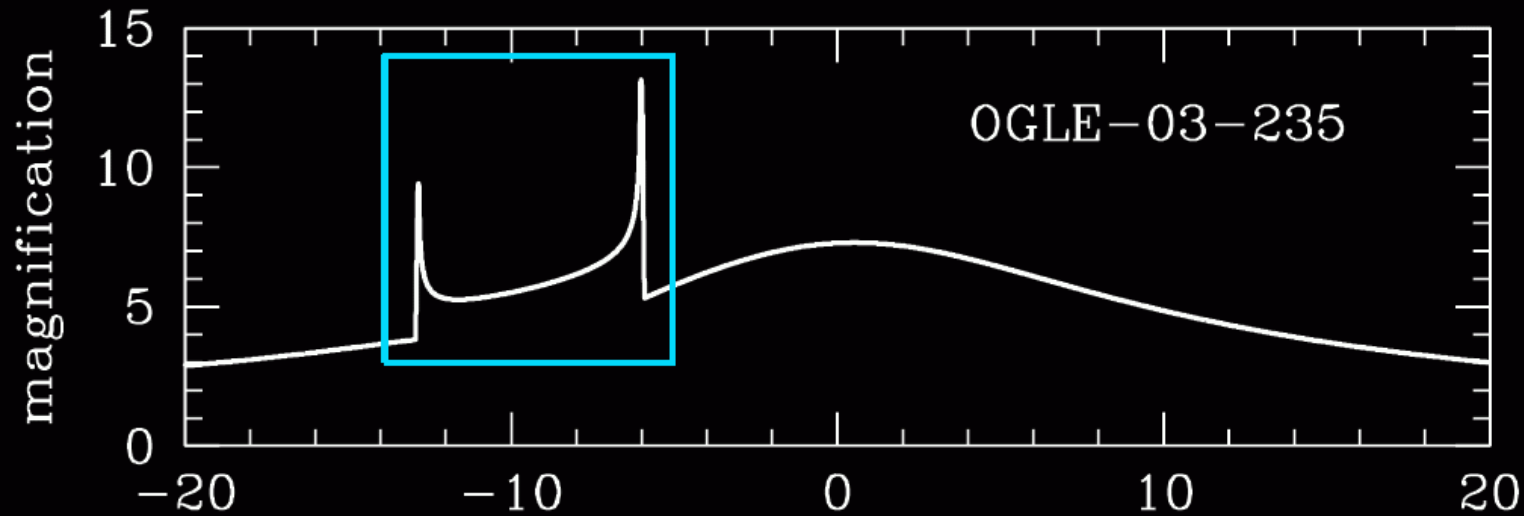


Application to complicated combinations of stellar densities  
Wambsganss, Witt, & Schneider, 1992, A&A, 258, 591

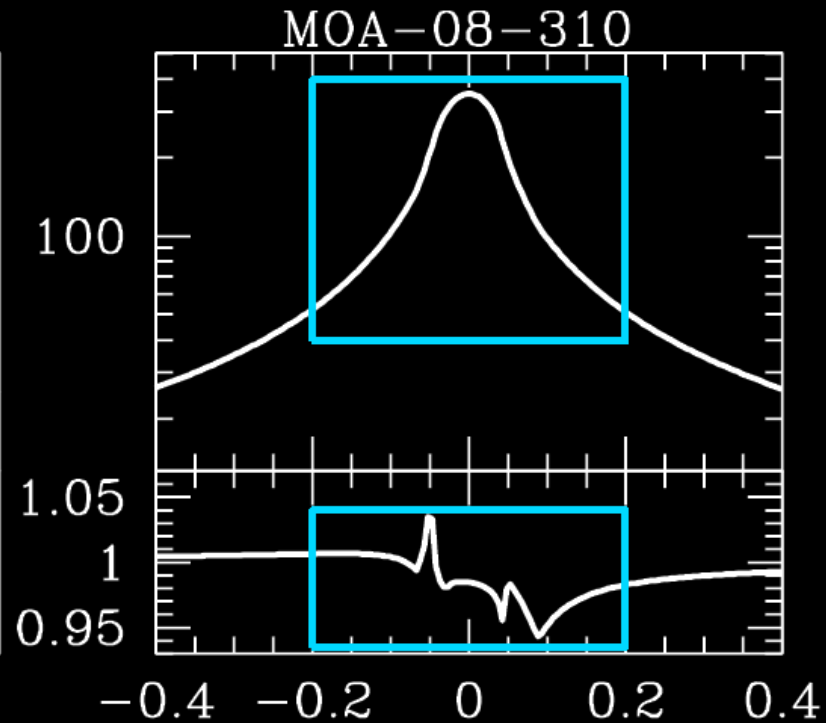
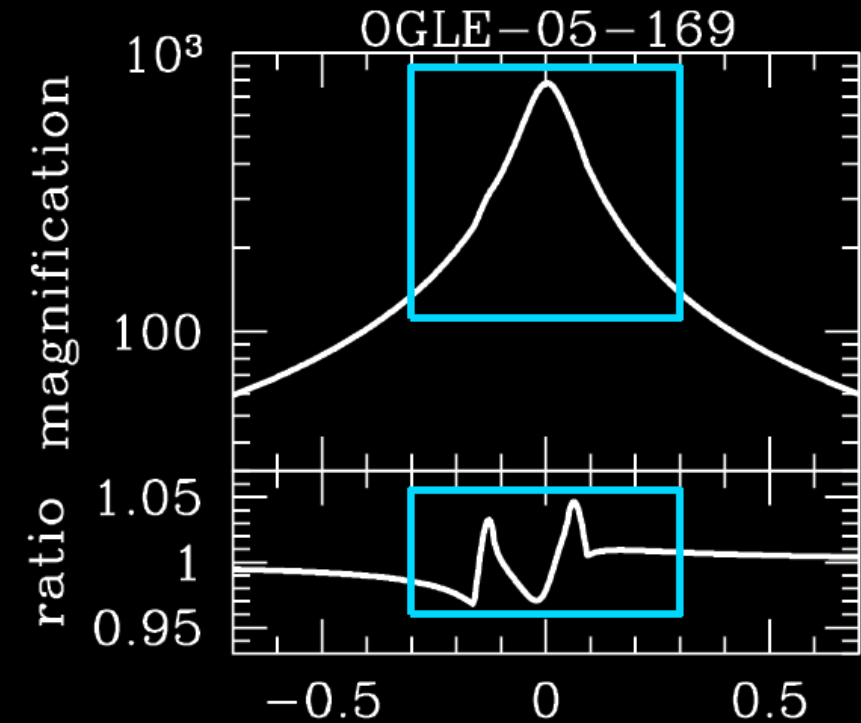
# Magnification Map (ray-shooting)

- The image-centered approach
- Bennett & Rhie, 1996, ApJ, 472, 660; Bennett, 2010, ApJ, 716, 1408
- Use point-source model except when source or image is close to caustic
- Shoot rays from point-source image centers
- Include partial images where disk crosses a caustic
- Polar coordinate system greatly reduces the number of needed grid points compared with Cartesian system

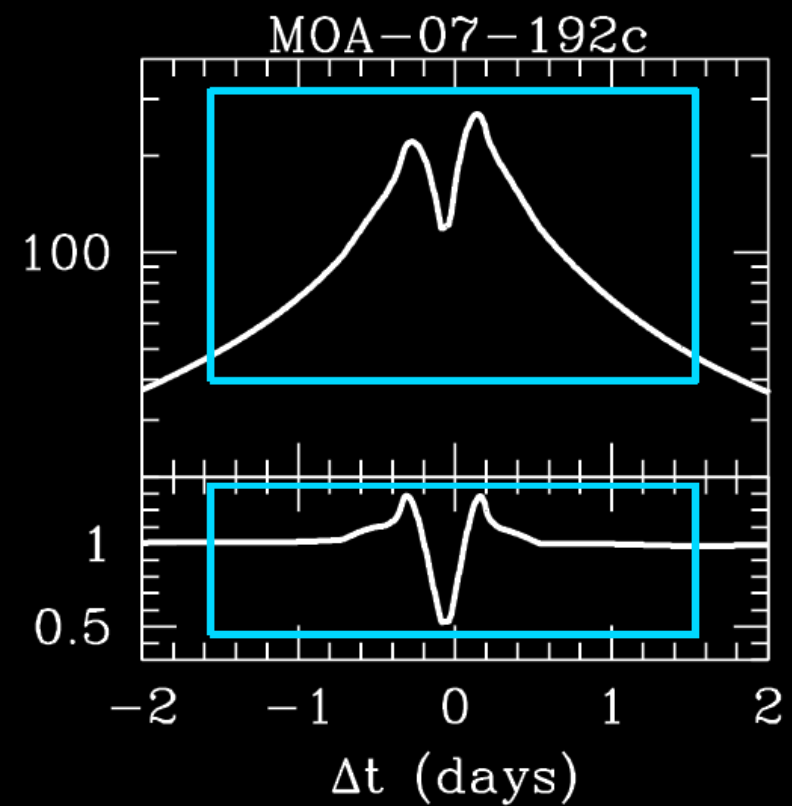
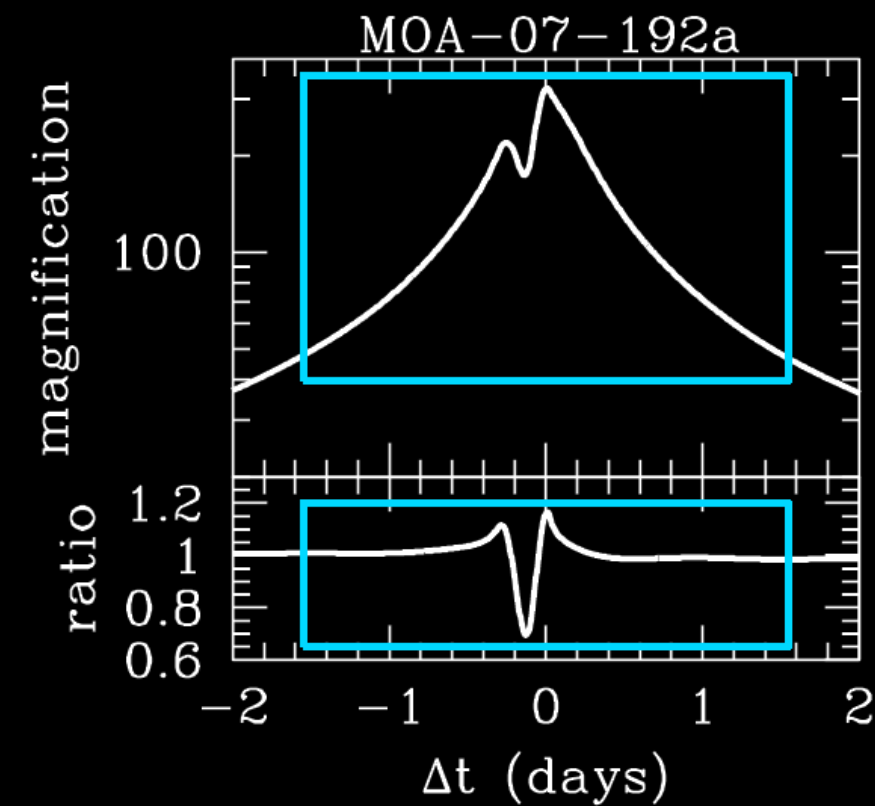




**Light curve  
calculation tests:  
low magnification**



Light curve  
calculation tests:  
high magnification



# Stokes/Green Theorem

- Perform contour integration in the image plane
- Stokes' theorem generalizes integration theorems in vector calculus

$$\oint_{\partial I} \mathbf{L} \cdot d\mathbf{x} = \iint_I dS (\nabla \wedge \mathbf{L}) \cdot \mathbf{n}$$

- Green's theorem is a special case in two dimensions

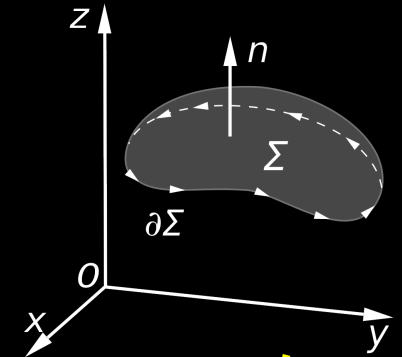
$$\oint_{\partial I} (L_1 dx_1 + L_2 dx_2) = \iint_I dx_1 dx_2 \left( \frac{\partial L_2}{\partial x_1} - \frac{\partial L_1}{\partial x_2} \right)$$



# Stokes/Green Theorem

- Perform contour integration in the image plane
- Stokes' theorem generalizes integration theorems in vector calculus

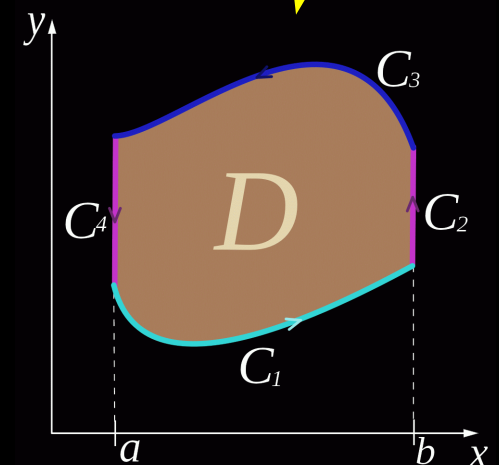
$$\oint_{\partial I} \mathbf{L} \cdot d\mathbf{x} = \iint_I dS (\nabla \wedge \mathbf{L}) \cdot \mathbf{n}$$



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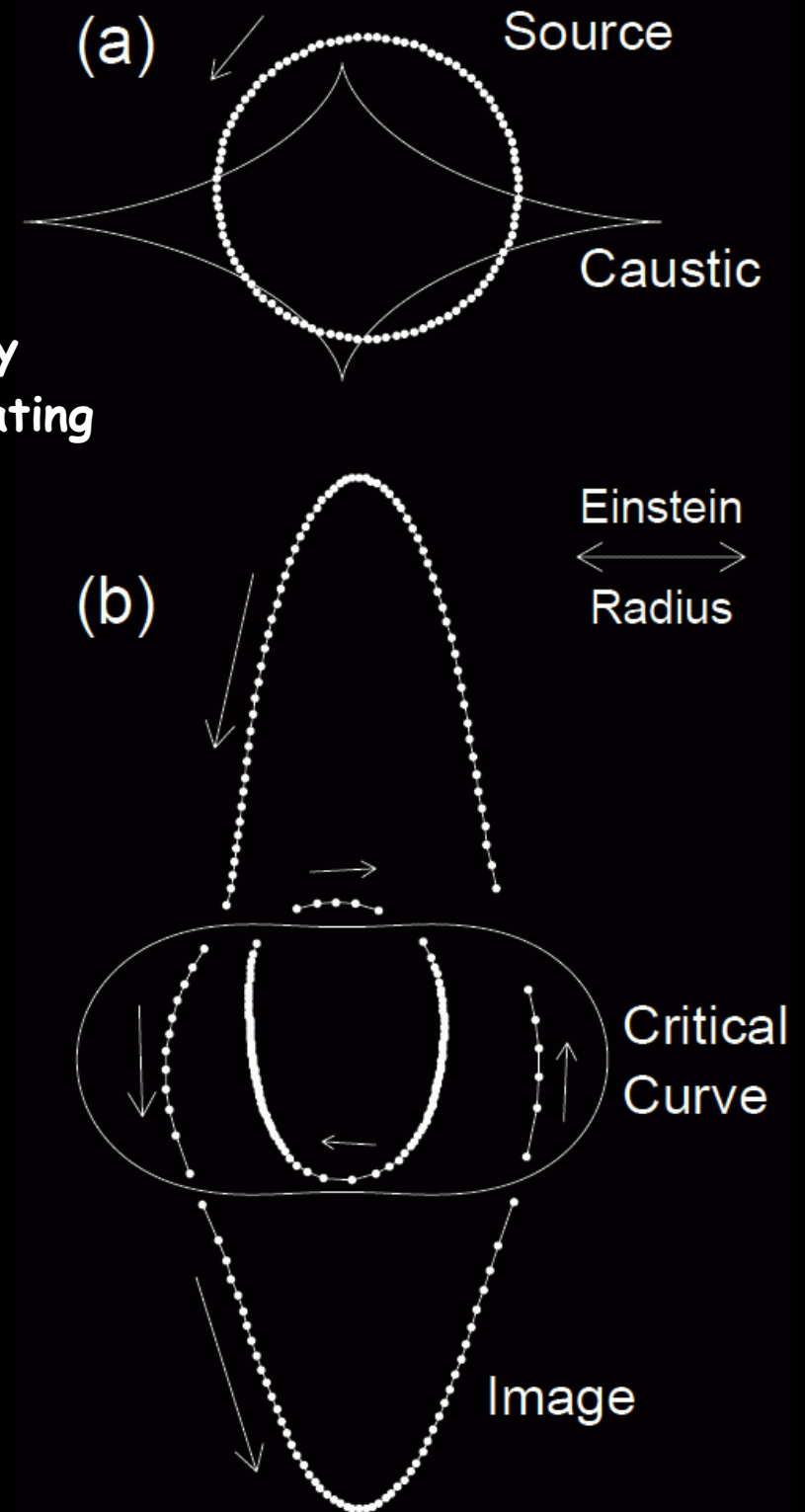
$$\oint_{\partial I} (L_1 dx_1 + L_2 dx_2) = \iint_I dx_1 dx_2 \left( \frac{\partial L_2}{\partial x_1} - \frac{\partial L_1}{\partial x_2} \right)$$

- Fast for uniform sources



# Stokes/Green Theorem

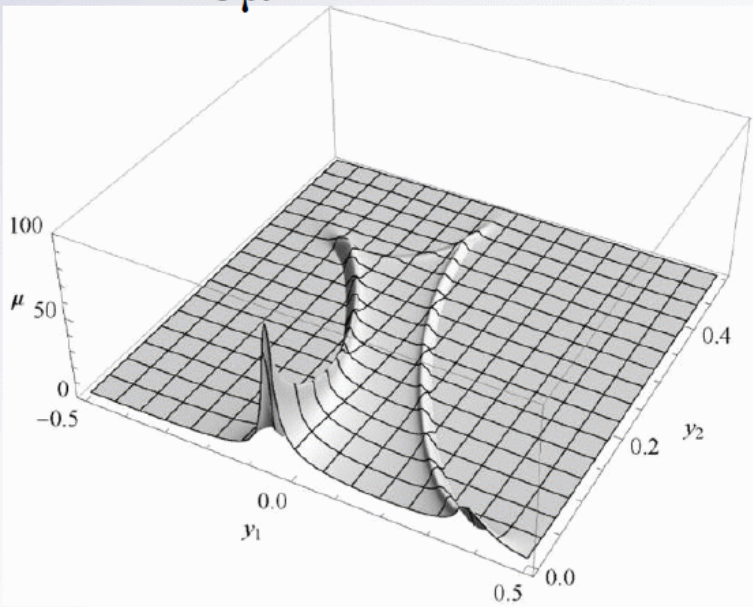
- Calculate area of a microlensing image by a contour integral on the image boundary
- Sample the source boundary by approximating as a polygon
- Invert the lens equation to find image boundaries
- Re-order the points in the sample
- Approximate the contour integral using the sample
- Model limb-darkening by using rings of constant brightness



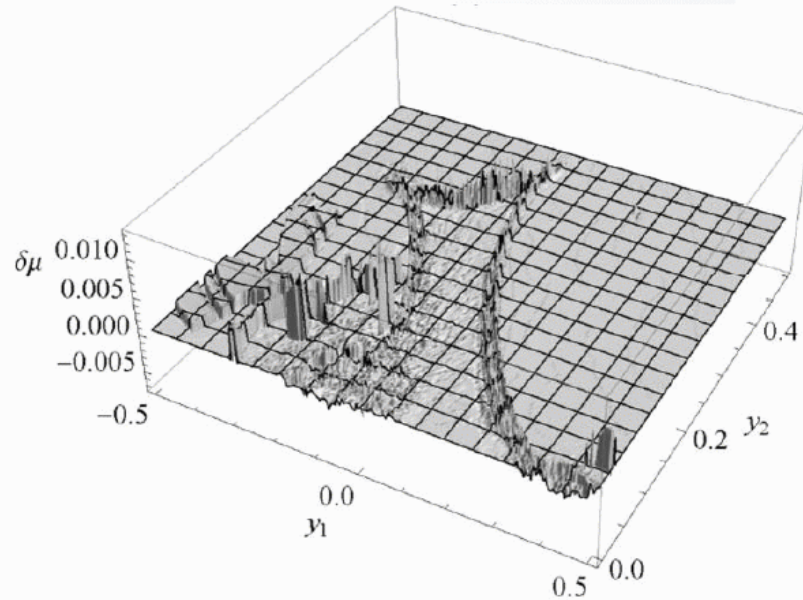
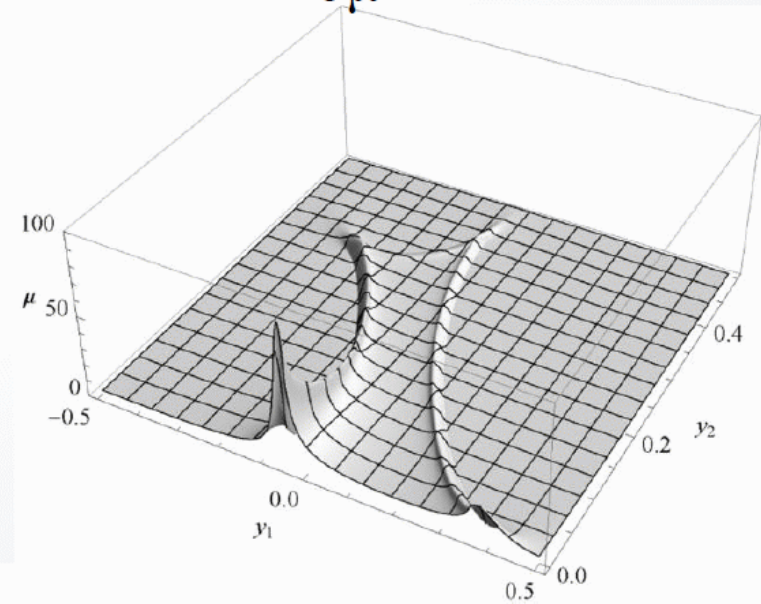
# Stokes/Green Theorem

- Gould & Gaucherel, 1997, ApJ, 477, 580
- Dominik, 1998, A&A, 333, L79  
(Application to microlensing)
- Dong et al., 2006, ApJ, 642, 842  
(Hybridization with inverse ray shooting)
- Dominik, 2007, MNRAS, 377, 1679  
(Adaptive grid search)
- Bozza, 2010, MNRAS, 408, 2188  
(Advanced contour integration)

$\delta\mu=10^{-2}$

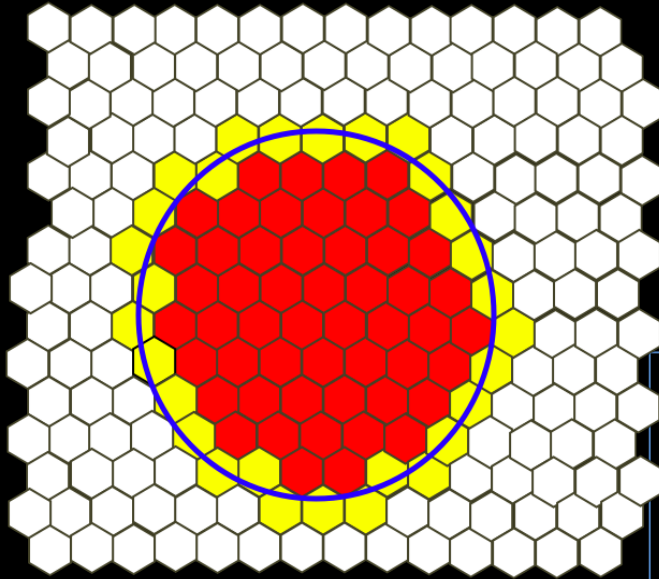


$\delta\mu=10^{-5}$



• Bozza, 2010, MNRAS, 408, 2188  
(Advanced contour integration)

# An Efficient and Robust Algorithm



*Dong, S., et al., 2006, ApJ, 642, 842 (Appendix A)*

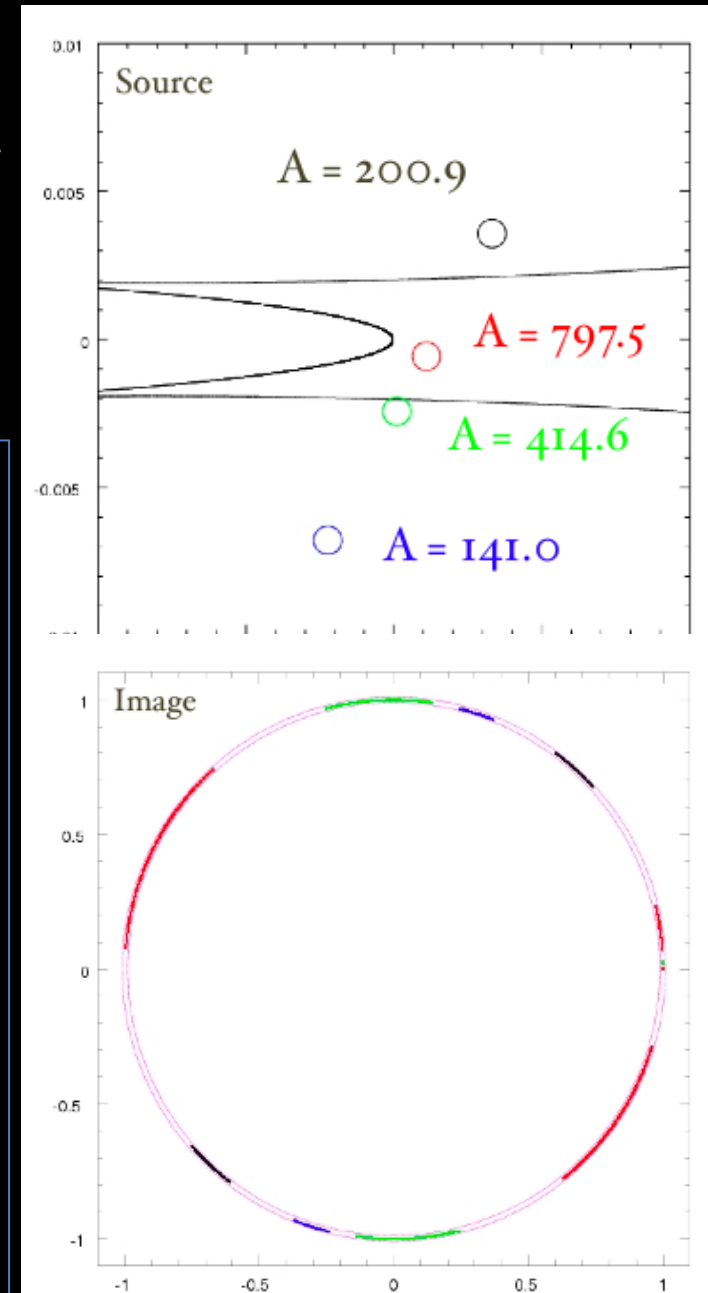
*Dong, S. et al., 2009, ApJ, 698, 1826*

- (Hexagon-cell Magnification) Map-Making (workhorse):

- Shoot rays from a narrow annulus on the image plane - reduce the overhead by orders of magnitude

- On the source plane, a combination of pixels and rays, with enhanced speed while preserving accuracy

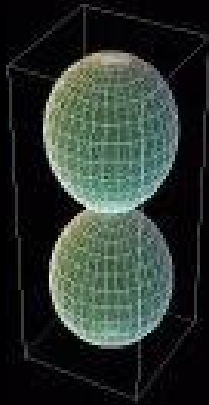
Loop-Linking  
(backup) combines  
contour integration  
and ray-shooting



# Hexadecapole

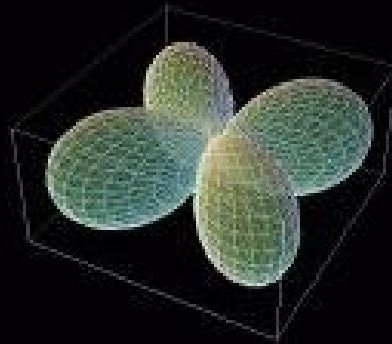
2

Dipole



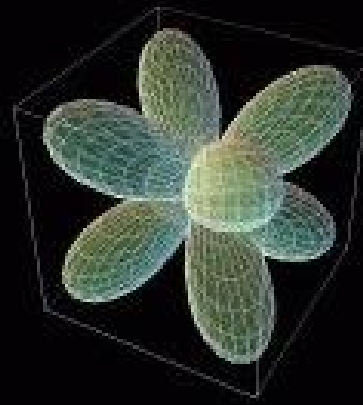
4

Quadrupole



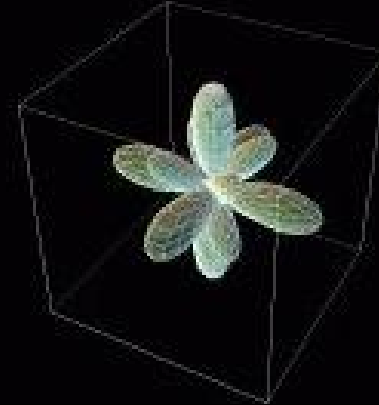
8

Octupole



16

Hexadecapole



- Uses 13 point "grid" in the source plane
- Several times slower than point-source calculations but orders of magnitude faster than finite-source calculations
- Requires more than two source radii away from caustic
- Best when combined with other methods
- Particularly useful where planetary orbital motion is included
- Gould, 2008, ApJ, 681, 1593

# Hexadecapole

$$A_{\text{finite}}(\rho, x_0, y_0) \equiv \frac{\int_0^\rho dw \int_0^{2\pi} d\eta A(w, \eta)}{\pi \rho^2}$$

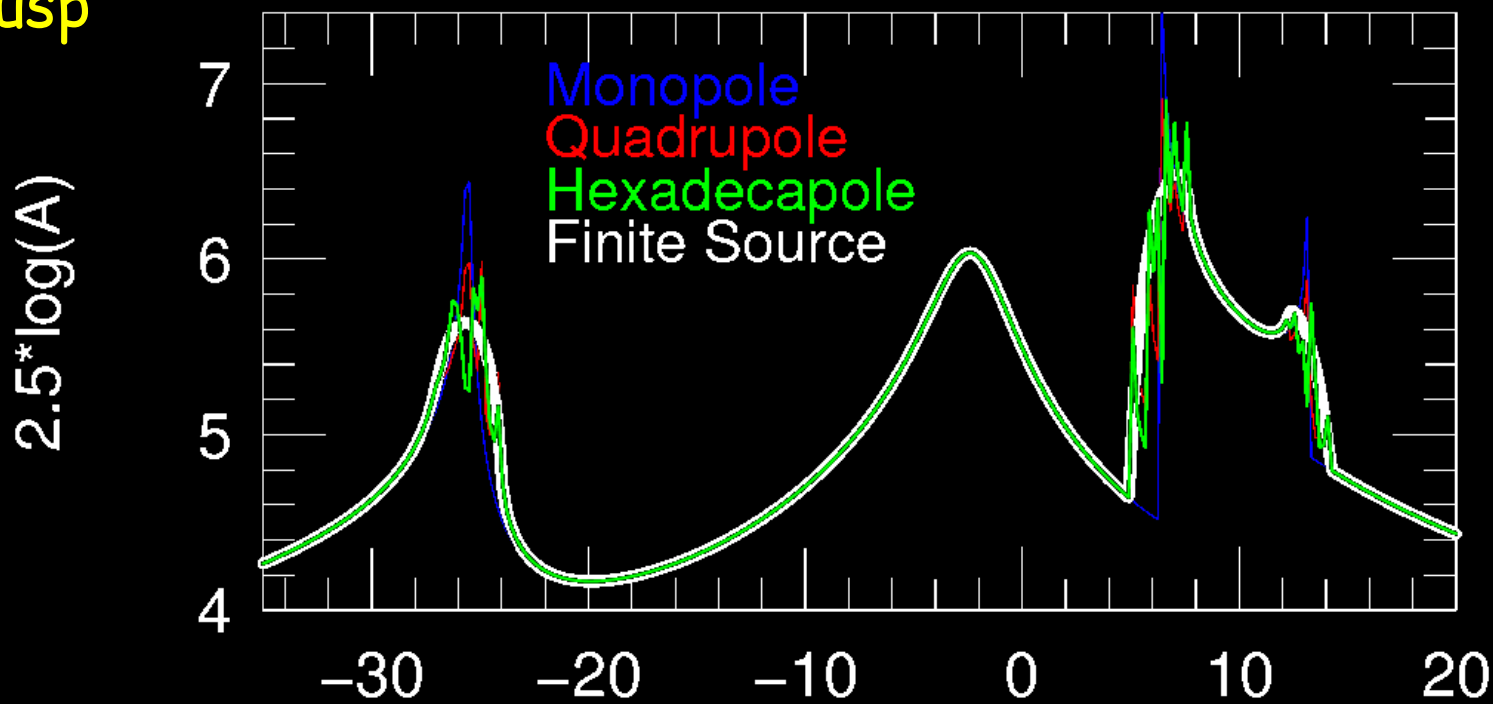
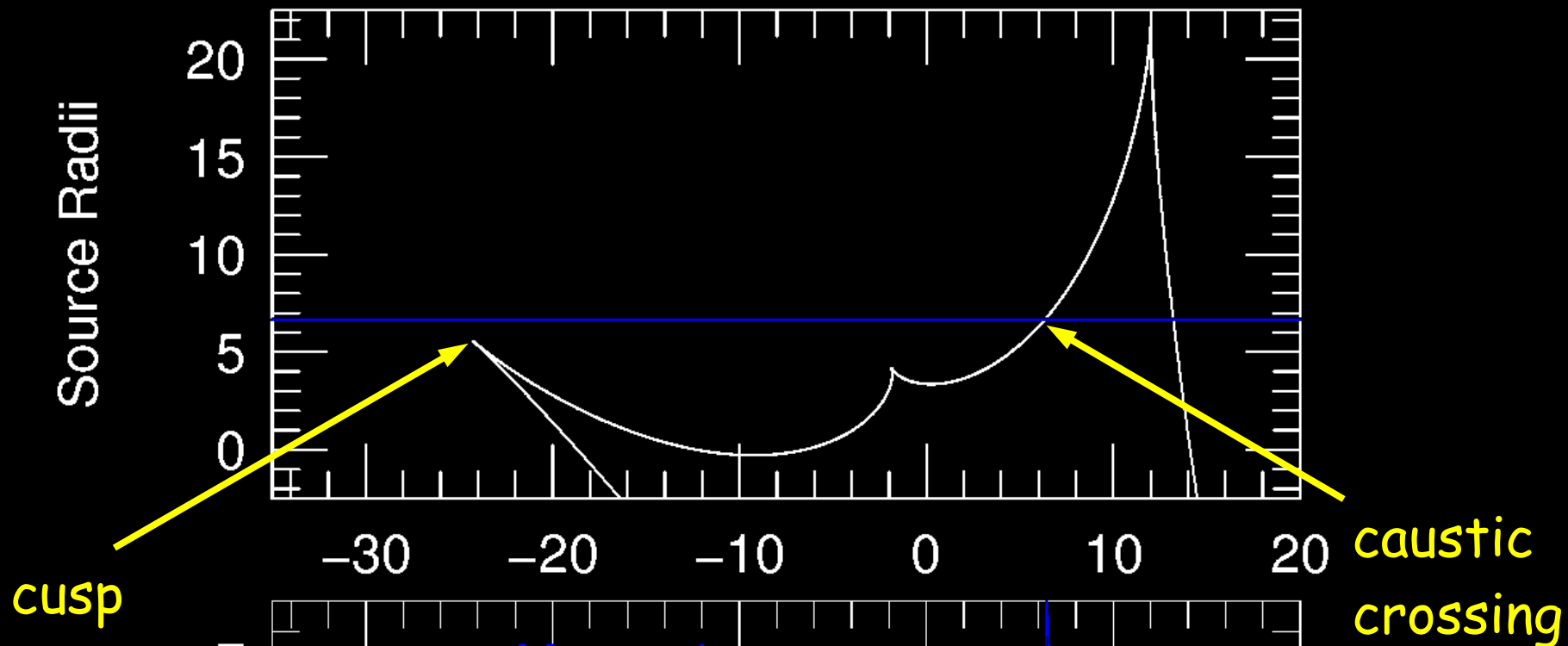
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# Hexadecapole

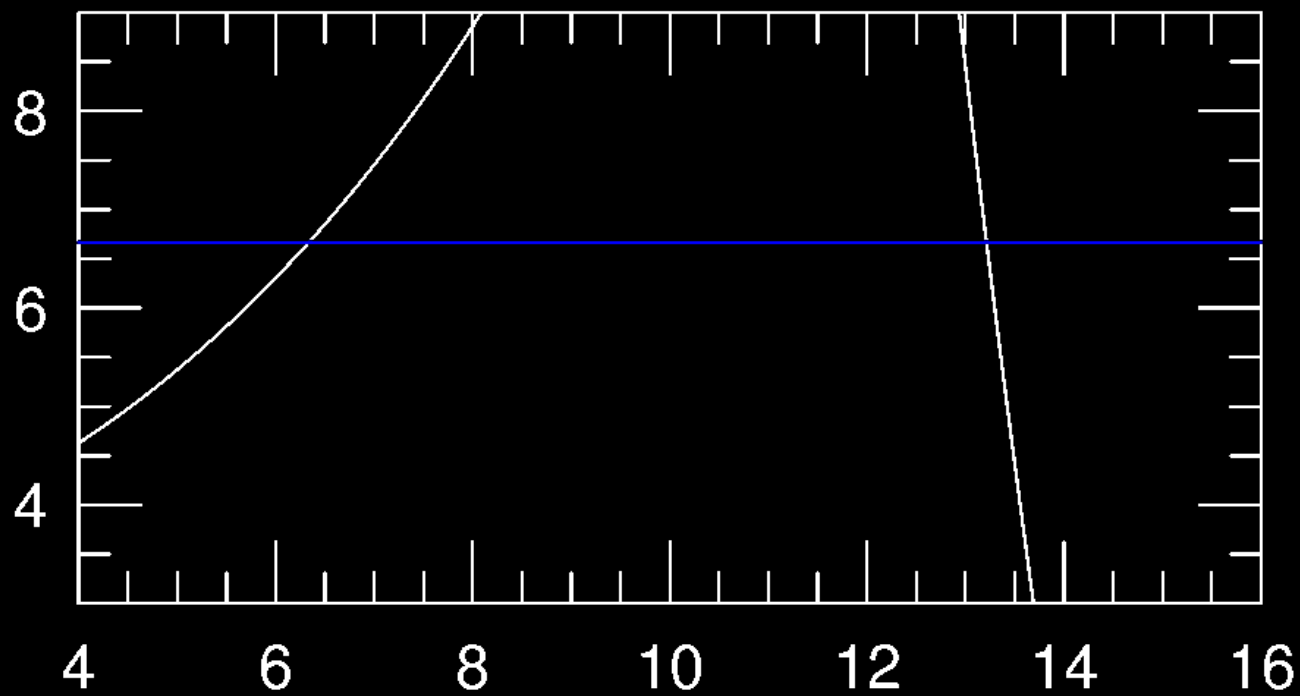
$$A_{w,+} \equiv \frac{1}{4} \sum_{j=0}^3 A \left[ w \cos \left( \phi + j \frac{\pi}{2} \right), w \sin \left( \phi + j \frac{\pi}{2} \right) \right] - A_0$$
$$= A_2 w^2 + \frac{(A_{40} + A_{44})(1 + \cos^2 2\phi) + A_{42} \sin^2 2\phi}{4} w^4,$$

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- Several times slower than point-source calculations but orders of magnitude faster than finite-source calculations
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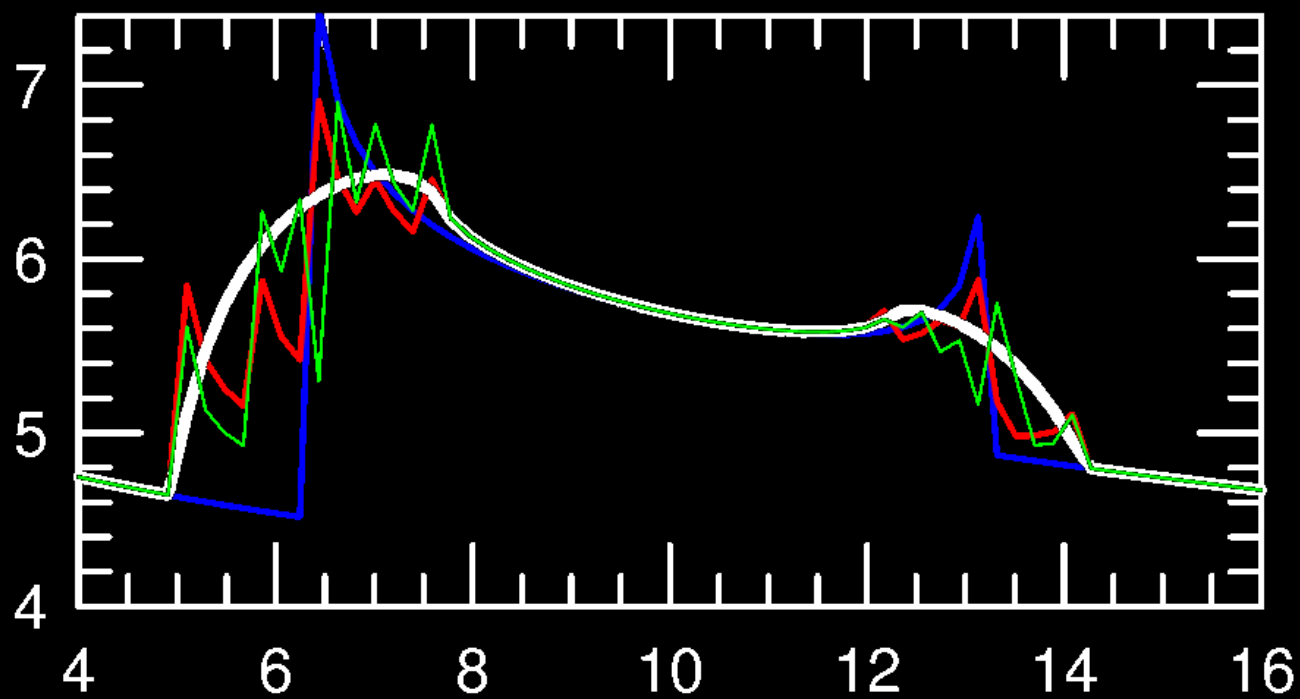




Source Radii



$2.5 \cdot \log(A)$



# MICROLENSING MAGNIFICATION CALCULATIONS WITH POINT-SOURCE AND FINITE-SOURCE EFFECTS

Practical implementation in Hands-on  
session after lunch

(Subo Dong & Jan Skowron)