

# Astrometry and Photometry with CCDs

William van Altena

Yale University

**Michelson Summer Workshop**

**25-29 July 2005**

# References

- **References**

S. Howell, “Handbook of CCD Astronomy”, Cambridge Univ. Press (2000)

A good up-to-date book on the CCD as a detector.

G. Walker, “Observational Astronomy”, Cambridge Univ. Press (1987), Ch. 7.

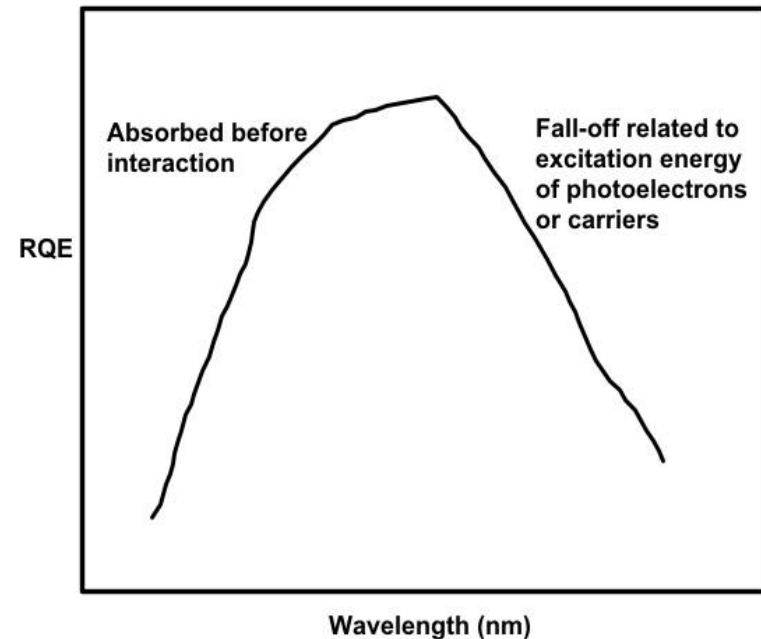
A good general introduction to all aspects of observational astronomy.

M. Newberry 1991, PASP, 102, 122

A detailed discussion of signal-to-noise

# CCD Quantum Efficiency

- Advantages of CCDs
  - High quantum efficiency.
    - Responsive QE
      - $RQE = N(\text{detected})/N(\text{incident})$   
25% > RQE > 95%
    - Detective QE
      - $DQE = (S/N)_{\text{out}}/(S/N)_{\text{in}}$
  - Spatially stable silicon substrate.
  - Reasonable resolution  
Typically 10-20 microns
  - Modest format size (4k x 4k)
- Disadvantage
  - Need to enhance blue response with down-converting phosphors.



# Front-side versus Back-side Illumination

- Front-side illumination
  - The electrical connections interfere with the access of the photons to the sensitive area
    - Potential for systematic “pixel-phase” position errors
    - Reduced sensitive area > lower QE
    - $QE(\text{max}) \sim 25\%$
- Back-side illumination
  - Thick semiconductor
    - Photons are absorbed and charge carriers created too far from the depletion layer
    - Thinning the backside to  $\sim 10$  microns can yield  $QE \sim 95\%$ !

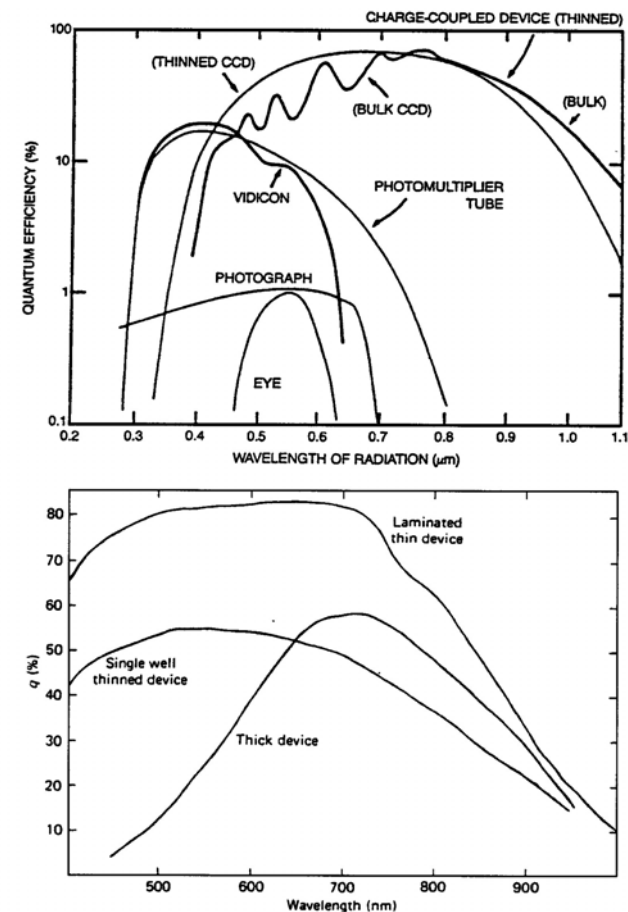


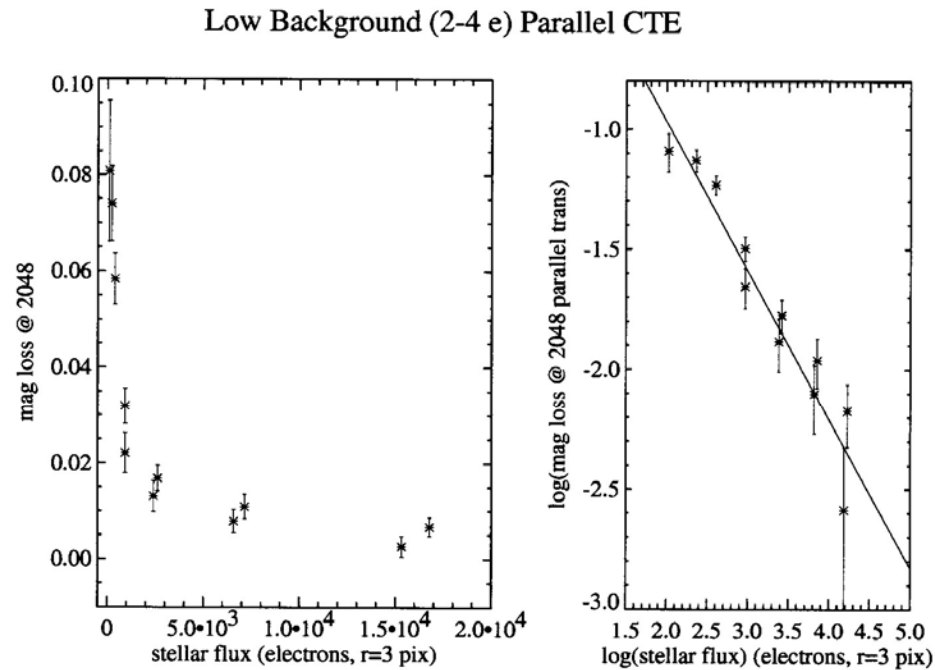
Fig. 3.2. The top panel shows QE curves for various devices indicating why CCDs are a quantum leap above all previous imaging devices. The failure of CCDs at wavelengths shorter than about  $3,500 \text{ \AA}$  has been completely eliminated via thinning or coating of the devices. The bottom plot shows representative QE curves for both thick and thinned devices. Laminated and single-well devices are two processes used to produce mechanically stable thinned CCDs.

# Charge-Transfer Efficiency

- Reiss (STScI-ACS 2003-009)
  - If the transfer of charge is not 100% from pixel to pixel during the read out, then CTE loss occurs.
  - CTE loss occurs from radiation damage, temperature of CCD and scene characteristics (# counts, extent of image, local and global background)
    - CTE                      1k chip                      2k                      4k
    - 0.999975                      3%                      5%                      10%
    - 0.9999975                      0.3%                      0.5%                      1%
  - Minimization of CTE loss:
    - “pre-flashing” chip
    - “charge injection” in front of image readout direction

# CTE versus Pixel Flux

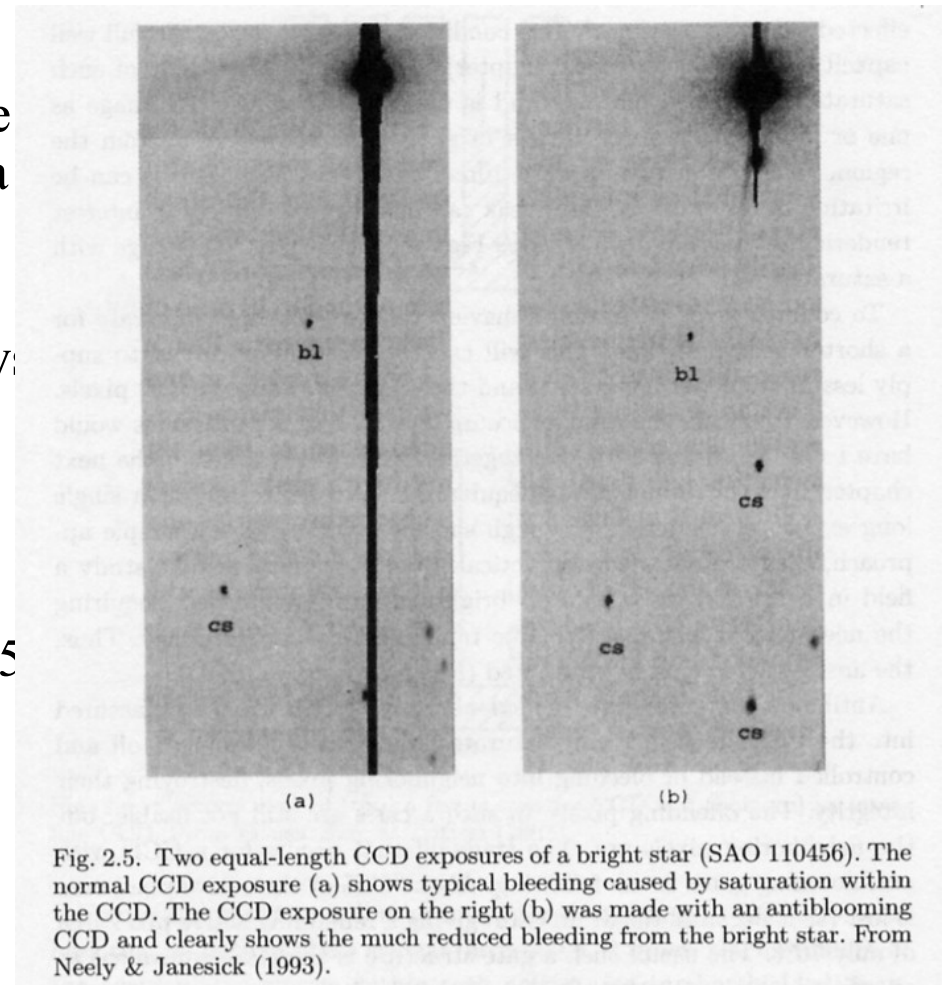
- Charge-transfer efficiency effects on Astrometry and Photometry



**Figure 4:** The dependence of parallel transfer CTE loss versus stellar flux contained within an  $r=3$  pixel aperture (all images at low sky level). As shown the relation is strong and suggests a power-law relation which was utilized in the correction formulae. The panel on the right shows the same in log-log space.

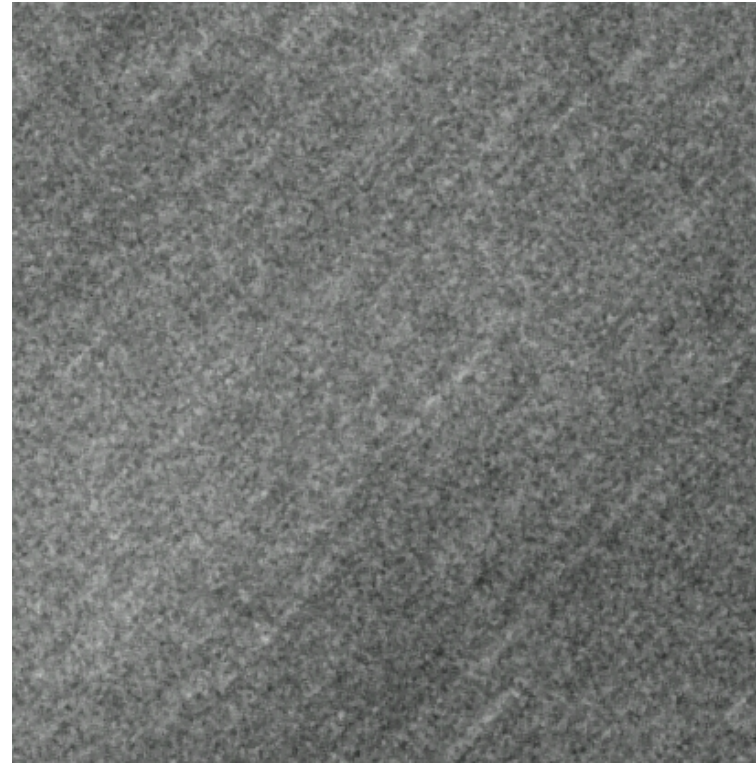
# Saturation Effects

- CCD Full- Well/Saturation level
  - Full well is proportional to the volume of a pixel, i.e. the area times the thickness of the depletion region.
  - Typical: Manufacturer says full well = 250,000  $e^-$ ; using a 16-bit A/D converter ( $2^{16} - 1 = 65535$  bits)
    - Gain =  $250,000/65535 = 3.81 e^- / \text{ADU}$  (analog-to-digital units)



# How the Flat-Field Effects Astrometry

- Types of Flat fields
  - Median Flat-Field (MFF) from several exposures in each filter.
    - Median gets rid of cosmic ray hits and bad data.
  - Super-sky flat - median average all exposures during a night.
    - Random star locations in field-of-view average out in median.
    - Generally low S/N due to low sky background.
  - Dome-diffuser flat - plexiglass diffuser in front of objective or corrector in a Schmidt.
    - Good for wide fov.
    - High S/N since it is a dome flat and bright lights.
    - Zhou, Burstein, et al. 2004, AJ 127, 3642.



S-S flat taken with wire objective grating ( $45^\circ$ ) in place. Diagonal streaks are due to partial overlapping of the stars in the deep and dense exposures. Fluctuations in the S-S flat are  $\sim 8\%$ .



# Micro- and Macro-Noise

- Micro-noise and Macro-noise - useful concepts
  - Macro:  $\sigma^2 = (1/N)\sum_{i=1}^N (\langle S \rangle - S_i)^2$
  - Micro:  $\sigma^2 = [1/2*(N-1)]\sum_{i=1}^{N-1} (S_{i+1} - S_i)^2$
- Macro-noise includes the large-scale non-uniformities (or errors) while Micro-noise includes only the point-to-point fluctuations
  - Astrometry is limited primarily by Micro-noise
    - Noise in defining the image center
  - Photometry is limited primarily by Macro-noise
    - Zero-point variations in the scale over the field

# Signal-to-Noise

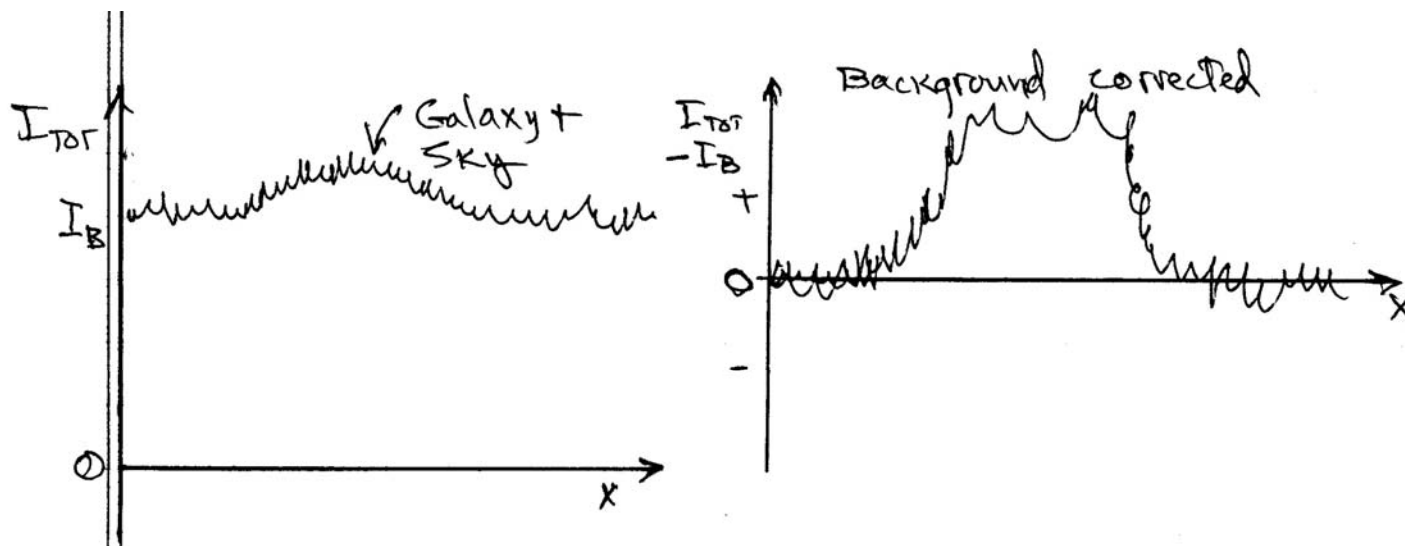
- References
  - Newberry 1991, PASP 103, 122.
  - Newberry 1994, CCD Astronomy
  - Howell, “Handbook of CCD Astronomy”, p. 53
- S/N sets the fundamental limit on our ability to measure the signal from the target.
- Bias, Dark and Flat-field corrections all contribute to degrading the S/N in the measured signal.

# CCD Photometry

- References
  - “Astronomical CCD Observing and Reduction Techniques”, S. B. Howell, ed., 1992, *ASP Conf. Series 23*.
  - Stetson, P. 1990, *PASP* 102, 932.
  - DaCosta, G., 1992, *ASP Conf. Series 23*, 90.

# Surface Photometry

- Surface Photometry
  - Generally trying to determine the surface brightness of a galaxy that is very faint in the presence of a “bright” sky.
  - Goal is to trace the galaxy out to, say,  $\leq 1\%$  of sky level.
    - $I(x,y) = \text{Galaxy}(x,y) + \text{Sky}(x,y)$
    - Need  $\text{Sky}(x,y)$  to better than 1% accuracy, or  $\sigma_s(x,y) \leq 0.01 * S(x,y)$
    - Macro noise is critical here, due to possible poor flat-field.



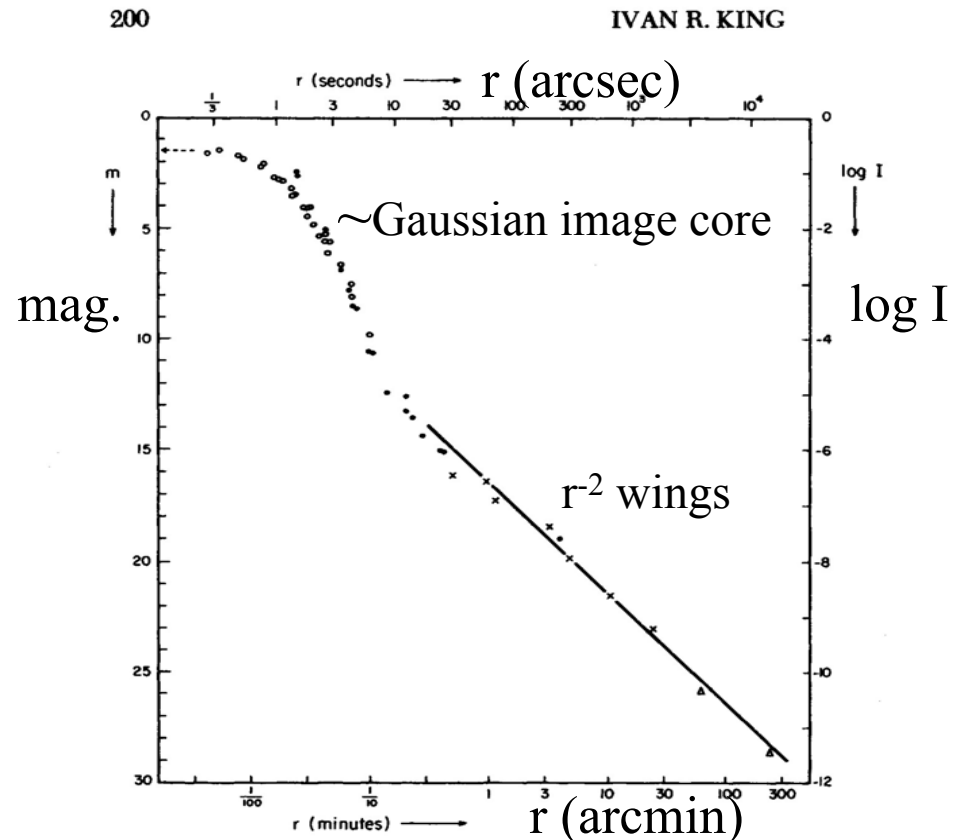
# Sky Background

- Sky background and its error
  - $\langle S \rangle$  = corrected mean sky
  - $\sigma_m$  = macro-noise in the sky
    - $\sigma_m$  = Poisson + Read Noise + Large scale background fluctuations.
    - $\sigma_m = [S + N_r^2 + \dots]^{0.5}$
  - $\sigma(\langle S \rangle)$  = error of mean sky
  - $X$  = % error desired in  $\langle S \rangle$ 
    - $X = \sigma(\langle S \rangle) / \langle S \rangle$ , e.g. 1%
      - $\sigma(\langle S \rangle) = \sigma_m / N^{0.5}$ , or  $X = \sigma_m / [N^{0.5} * \langle S \rangle]$
      - $N = \{\sigma_m / [X * \langle S \rangle]\}^2$ 
        - »  $\langle S \rangle = 100$  counts,  $N_r = 5$  counts,  $X = 1\%$ , then  $N=121$  pixels
        - »  $\langle S \rangle = 100$  counts,  $N_r = 5$  counts,  $X = 0.1\%$ , then  $N=12,100$  pixels

# Stellar Profile

- King, PASP 83, 199, 1971
  - Stellar profile observed through atmosphere.
    - Sky background near star is affected by presence of stellar wings

**FIG. 1** — Surface brightness, in magnitude per square second, in the image of a star of magnitude zero. Open circles are derived from 60-inch Cassegrain images, closed circles from diameters of NPS stars on the Palomar Observatory Sky Survey (POSS), and crosses from other stars on POSS. Straight line is inverse-square law found by de Vaucouleurs. Triangles are from sky brightness near the sun.

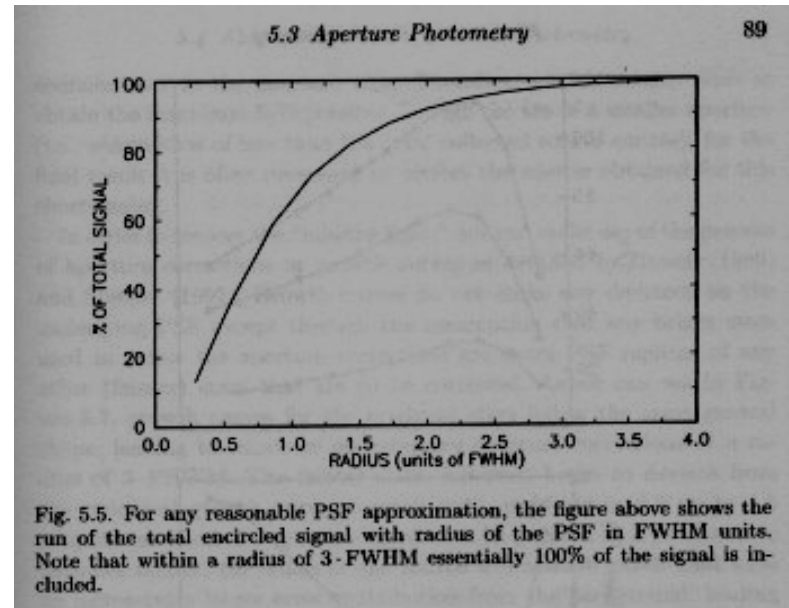


# Aperture Photometry

- Aperture photometry
  - DaCosta, G. ASP Conf. Series **23**, 90, 1992.
    - Select an aperture of radius  $r$  that contains the image and sum all of the pixels that fall within the aperture
      - $I_{\text{sum}} = I_* + \langle S \rangle$
      - What radius should be used to include all of the stellar flux?
        - » Remember that the King stellar profile extends many arcsec.
      - What about the sky within the aperture?

# Aperture Corrections

- Stetson, P. PASP **102**, 932, 1990
- Large  $r$  to include all star light, but:
  - Includes other stars
  - Adds sky noise
- All stars have the same psf, so all psfs scale with the # photons
  - psf is constant, except for
    - optical aberrations
    - seeing variations over field (short exposures)
  - CCD is linear, except for
    - saturation and CTE
  - $dm = -2.5 \log_{10} I_2/I_1$ 
    - $dm = -2.5 \log X$ 
      - Constant aperture correction, **dm**.



“Optimal” aperture = 1.35 FWHM  
Ap. corr. = 0.2 mag for a Gaussian

Star image must be accurately centered in aperture for this to work.

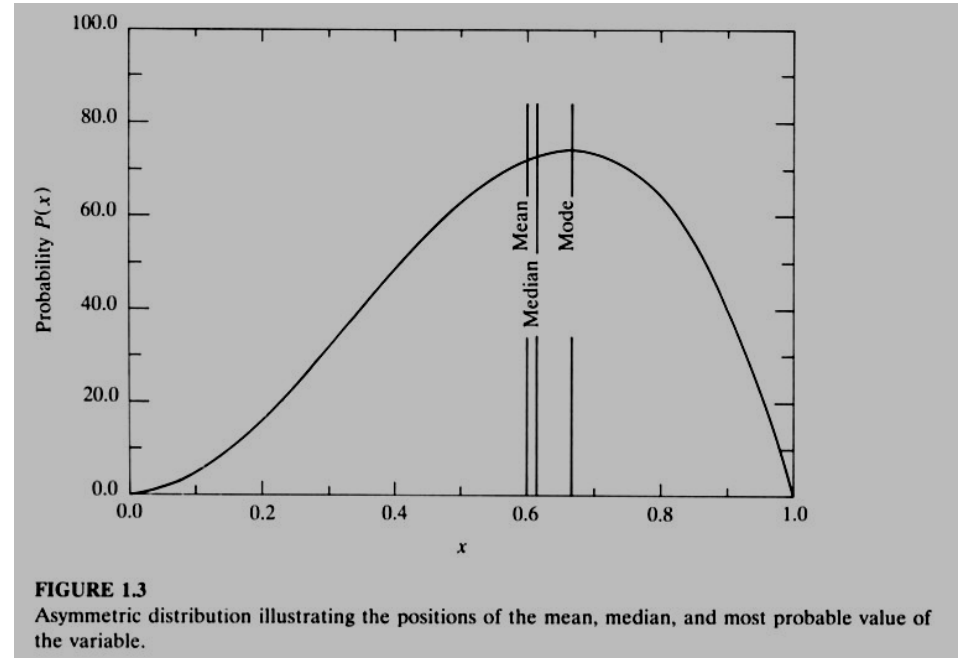


# Photometric Errors

- Sky background
  - Select many star-free spots and average the results
  - Probability plot analysis of background and scale to aperture
- Error in photometry
  - $S_o$  = total stellar signal, but noise and read-out noise are per pixel
    - $\sigma^2 = S_o + n(S_s + N_r^2)$ 
      - $\text{mag} = m_o - 2.5 \log_{10} S_o$
      - $\sigma_m = \sigma_{mo} - 2.5 (\log_{10} e) \sigma(S_o) / S_o$
      - $\sigma_m = \sigma_{mo} - 1.086 (1/S_o) \sigma(S_o)$
      - $\sigma_m = \sigma_{mo} - 1.086 (1/S_o) \{S_o + n(S_s + N_r^2)\}^{0.5}$ 
        - » = Zero pt. Error + Poisson + Sky + Read-out noise
        - »  $S_o$  is fixed, so the larger  $n$ , the greater the contribution of the **sky** and **read-out** noise is to the error.

# Sky Noise

- Sky noise
  - Ideally the histogram of pixel values is unimodal and we determine the mode.
  - Mode is usually difficult to determine and poorly defined.
  - Kendal & Stuart, p.40, 1977
    - $\text{mean-mode} = 3(\text{mean-median})$



# Fitting PSFs

- King found the stellar profile to be approximately Gaussian in the core:
  - $G(x) = (\sigma_x \sqrt{2\pi})^{-1} \exp\{-0.5[(x - x_c)/\sigma_x]^2\}$ 
    - $\int_x G(x) = 1.0$
    - $G(x = x_c) = 0.3989$
    - $G(x = x_c \pm \sigma_x) = 0.2420$ : 0.607 height at  $x = x_c$
    - $G(x = x_c \pm \text{fwhm}/2) = 0.1995$ : one-half the height at  $x = x_c$
  - Gaussian doesn't fit in wings, so other functions are added
    - Modified Lorentzian:  $L(x) = C * \{1 + (x^2/\sigma^2)^\beta\}^{-1}$
    - Moffat function:  $M(x) = C * \{1 + (x^2/\sigma^2)\}^{-\beta}$ 
      - $C = \text{constant}$
  - Not even those are perfect so a table of corrections ( $H(x,y)$ ) is added to give the final model PSF:
    - $\text{PSF}(x,y) = [a * G(x,y) + b * L(x,y) + c * M(x,y)] * [1 + H(x,y)]$
    - DAOPHOT and IRAF: See Stetson, P. in PASP **102**, 932, 1990.

# Photometric Passbands

- Passbands and filters
  - UBVRI: Bessell PASP 102, 1181, 1990
  - JHKLM: Bessell & Brett PASP 100, 1134, 1988
  - IR: Astrophys. Quant. IV (AQ4), A. N. Cox, ed. Ch. 7.1-7.7
  - Visual: AQ4 Ch 15.3
  - Asiago Database on photometric systems (ADPS)
    - <http://www.pd.astro.it/Astro/ADPS/>

# Photometric Surveys

Band	Sloan (NG cap)	DENIS (South)	2MASS (All sky)	UCAC (All sky)	NPM (N) SPM (S)
u	22.0				
g	22.2				
r	22.2				
i	21.3	18.5			
z	20.5				
B					18.5
V					17.5
R				16.5	
J		16.5	15.8		
H			15.1		
K <sub>s</sub>		14.0	14.3		

# Image Centering

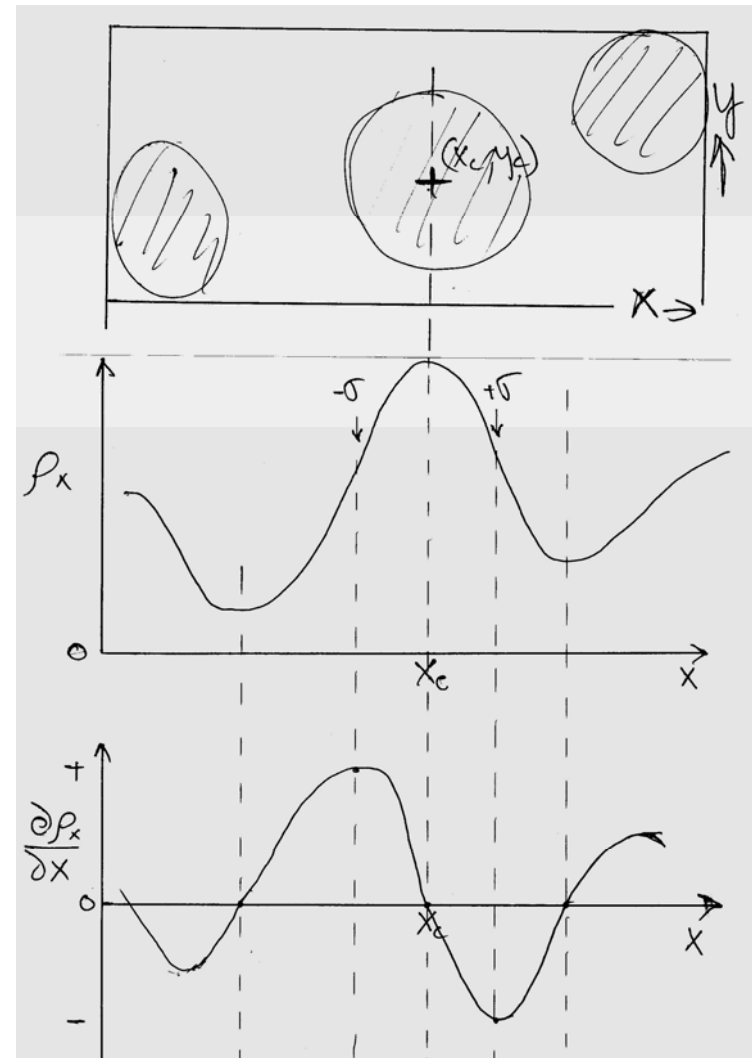
- References
  - van Altena and Auer: in “Image Processing Techniques in Astronomy”, p. 41, 1975
  - Auer and van Altena: AJ 83, 531, 1978
  - Lee and van Altena: AJ 88, 1683, 1983

# Centroids versus Image Centers

- Given an intensity distribution,  $S(x,y)$ 
  - The **Centroid, center of mass** or **1st moment** of the distribution.
    - See: van Altena and Auer: in “Image Processing Techniques in Astronomy”, p. 41, 1975
    - $\langle x \rangle = \sum_{x,y} \{x_i * [S(x,y) - B]\} / \sum_{x,y} [S(x,y) - B]$ ,
    - $\langle y \rangle = \sum_{x,y} \{y_i * [S(x,y) - B]\} / \sum_{x,y} [S(x,y) - B]$ 
      - where B is the assumed sky background around the image.
      - The centroid is very sensitive to the adopted sky background, but it also works well for very faint images.
  - The **Image Center**
    - See: Auer and van Altena: AJ 83, 531, 1978
    - and Stetson in DAOPHOT manuals.
    - The Marginal distributions are defined by:
      - $\rho_x(x) = N_y^{-1} \sum_y S_o(x,y)$
      - $\rho_y(y) = N_x^{-1} \sum_x S_o(x,y)$

# Marginals and Image Centers

- Derivative-search centers
  - Take the x-marginal,  $\rho_x(x)$ , in the middle panel.
    - Image crowding noted where  $\rho_x(x)$  increases at edges of the diagrams.
  - The derivative of the x-marginal,  $\rho'_x(x)$ , is in lower panel.
    - Peaks at  $\pm$  gaussian radius,  $R_x = \text{FWHM} / 2.36$
    - Zeros at image center and inflection points in  $\rho_x(x)$  that indicate image crowding.





# Univariate Gaussian

$$\rho_x(x) = a_x + b_x(x-x_c) + [1 + c_x(1.5t - t^3)]h_x \exp(-0.5t^2)$$

- Univariate Gaussian
  - $t = (x-x_c) / R_x$
  - $h_x = N_x / [R_x \sqrt{2\pi}]$
  - $R_x =$  Gaussian radius
  - $b_x =$  sloping background
  - $c_x =$  skewness of image
  - Generally take  $b_x = c_x = 0.0$ , since there is usually a high degree of correlation between the *odd* terms and this degrades the image center precision, i.e. use a symmetric Gaussian for the fit.

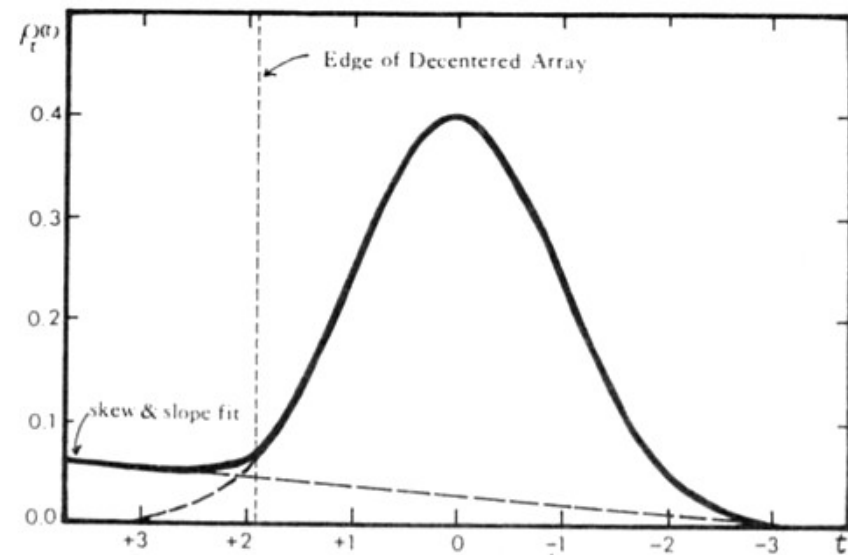
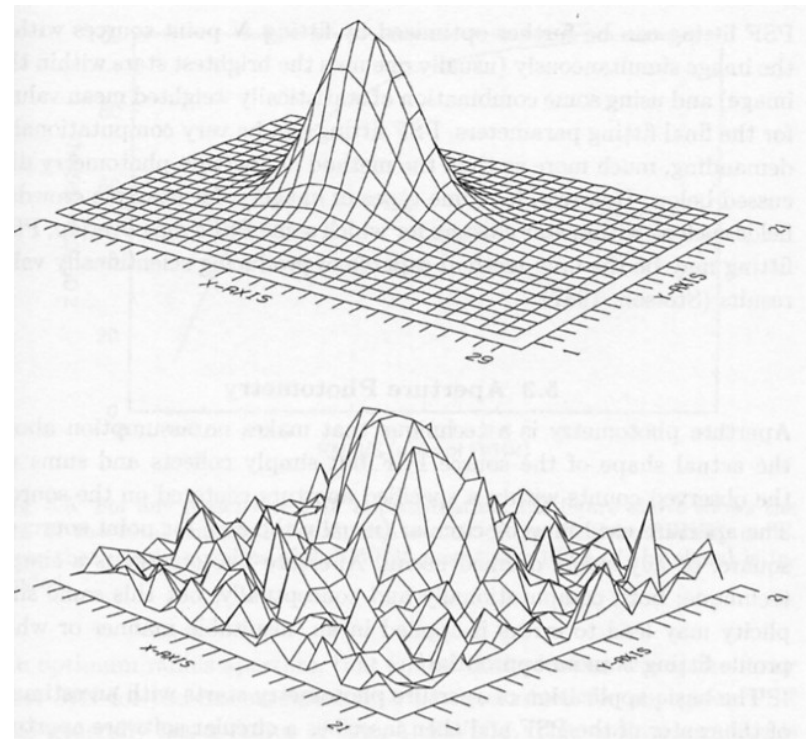


FIG. 2. Illustration of how a poorly centered array can yield a spurious, but apparently good fit with a skewed Gaussian and a sloping background when the true image is symmetric and the background flat.

# Bivariate Gaussian

$$F(x,y) = D_0 \exp(-0.5r^2 / R^2) + B$$

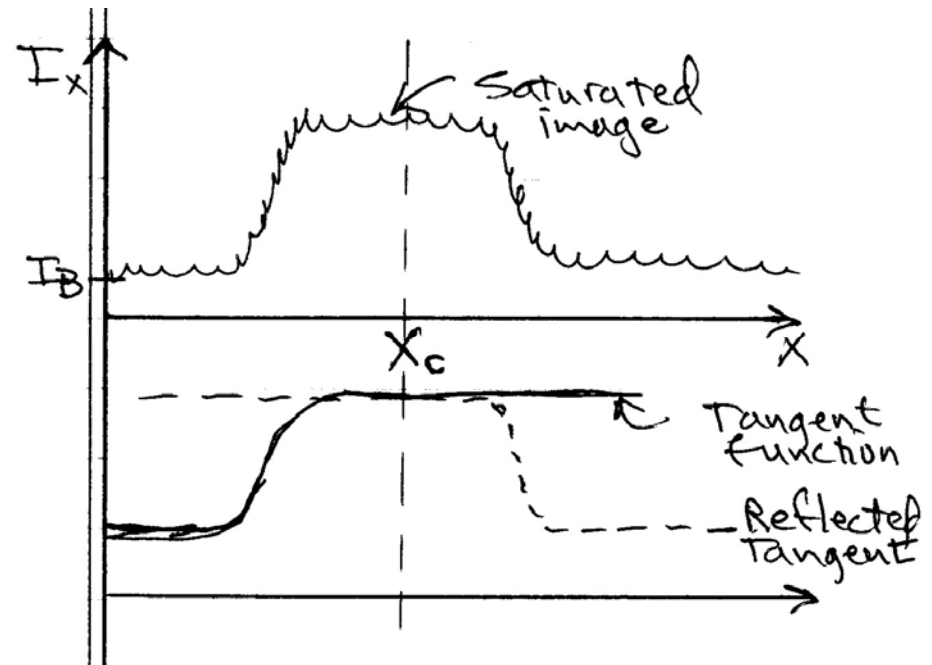
- Bivariate Gaussian
  - Lee and van Altna: AJ 88, 1683, 1983
  - $D_0$  = image height at center
  - $r^2 = (x-x_c)^2 + (y-y_c)^2$
  - $R$  = Gaussian radius
  - $B$  = background
- Precision
  - #1 Bivariate
  - #2 Univariate
  - #3 Centroid
- Convergence
  - Inverse order of precision, i.e. the centroid is most stable, especially for faint images.



# Dealing with Saturated Images

$$F(x) = a[2\tan^{-1}bc]^{-1} \{ \tan^{-1} [b(x-x_0+c)] - \tan^{-1} [b(x-x_0-c)] \} + d_0$$

- Saturated photographic image usually has a flat top and Gaussian fits poorly
- Two arctangent functions fit very nicely
- Stock, J. ~1997
  - $a$  = image height above background
  - $d_0$  = background
  - $c$  is proportional to image width
  - $b$  = scale factor for image slope and gradient
- Winter (Ph.D. thesis 1995) - a generalized Lorentz profile also works well.



# Image Center Precision Estimators

- Precision of the image center
  - $h = N_x / [R_x \sqrt{(2\pi)}]$  = central height of the Gaussian
  - $R = \text{FWHM} / 2.36 = \text{Gaussian radius}$
  - $N = \text{integral under the univariate function}$
  - $\sigma_1 = \text{standard error of fit to the univariate function, i.e. the dispersion of the intensities around the best fit to the marginal distribution.}$
  - $\sigma_h = \text{standard error of the central height of the Gaussian}$
  - $\varepsilon = (2/\pi^{1/2}) (\sigma_1/h) R^{1/2}$
  - $\varepsilon = (2\pi^{1/4}) (\sigma_1/N) R^{3/2}$
  - $\varepsilon = \sqrt{2} (\sigma_h/h) R$  based on photometric precision
  - $\varepsilon = \sqrt{2} R (S/N)^{-1}$  based on photometric precision

# Summary of Astrometry & Photometry (1/2)

- Maximum Photometric precision
  - Emphasis here is on the total number of counts
  - Fitting function is not too important, since a look-up table must be used to correct to psf.
  - Stetson: PASP 102, 932, 1990
  - DAOPHOT and IRAF manuals
  - DaCosta ASP Conf. Series 23, 90, 1992
- Under-sampled images, e.g. HST
  - Anderson & King: PASP 112, 1360, 2000
  - Anderson & King: PASP 115, 113, 2003
  - Druckier, et al. AJ 125, 2559, 2003

# Summary of Astrometry & Photometry (2/2)

- Maximum Astrometric precision
  - Emphasis must be placed on the image-profile gradients
  - Use functional fits to the image with weighting by the derivatives of the function
    - Auer and van Altena: AJ 83, 531, 1978
    - Lee and van Altena: AJ 88, 1683, 1983
  - For saturated images
    - Stock (1977) arctangent functions allow for saturation in photographic images
- Quick and Dirty Astrometry
  - Centroids
    - van Altena and Auer: in “Image Processing Techniques in Astronomy”, p. 41, 1975