

DIRECTLY IMAGING RUNAWAY ACCRETION

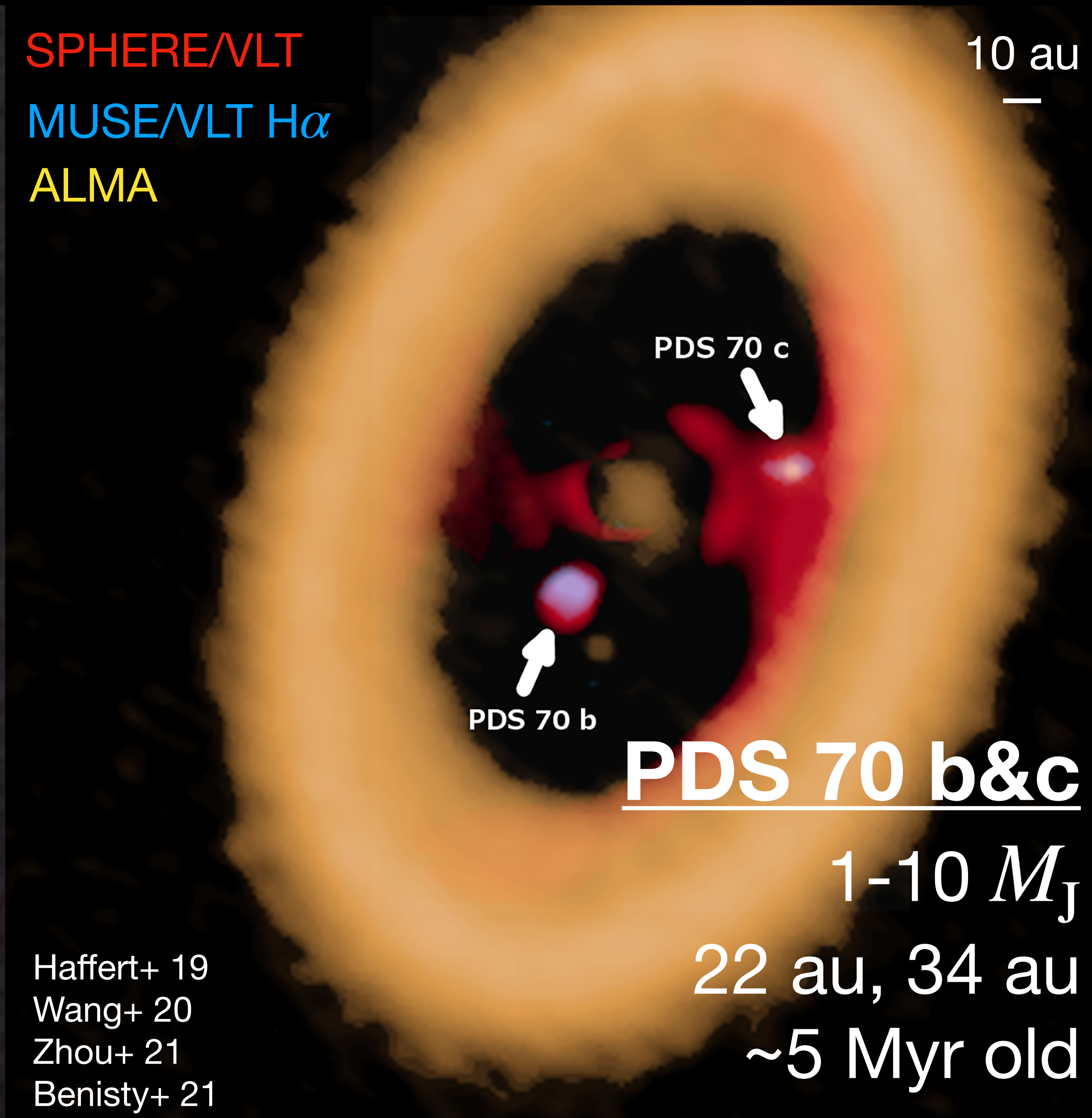


Nick Choksi
51 Pegasi b fellow, Caltech

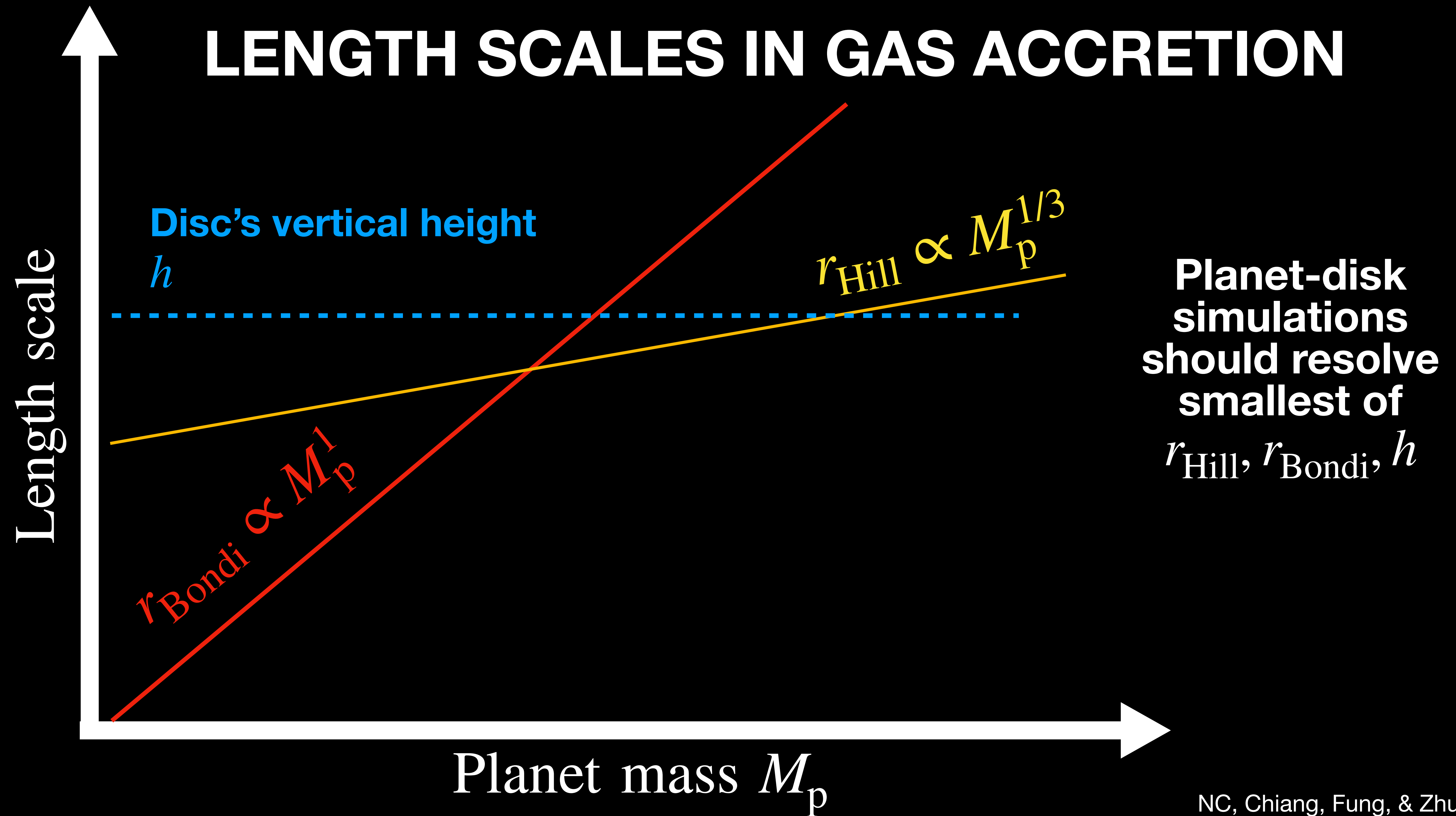
van Capelleveen+ 25, VLT/SPHERE

NC & Chiang 2022
NC, Chiang, Fung, & Zhu 2023
NC & Chiang 2025
NC & Gao in prep

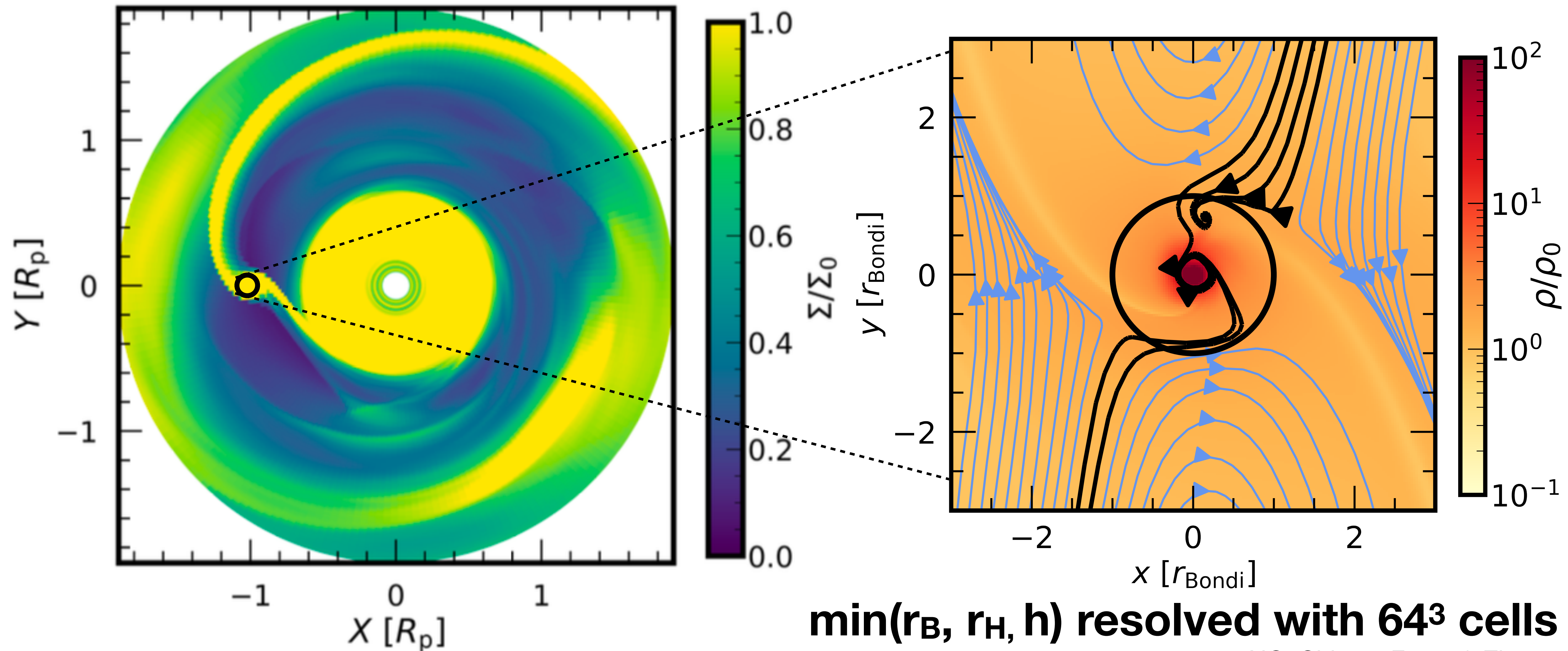
OBSERVED ACCRETING PLANETS



LENGTH SCALES IN GAS ACCRETION

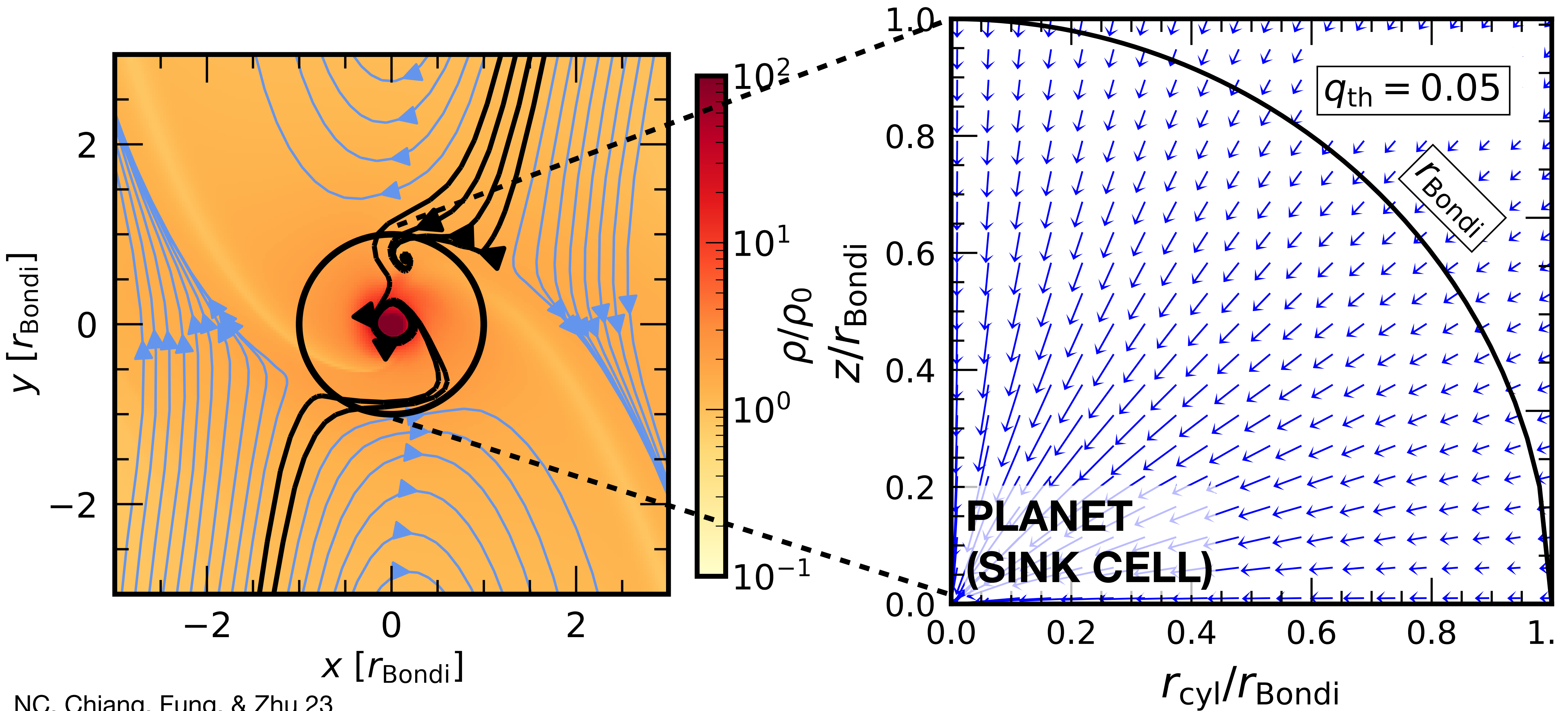


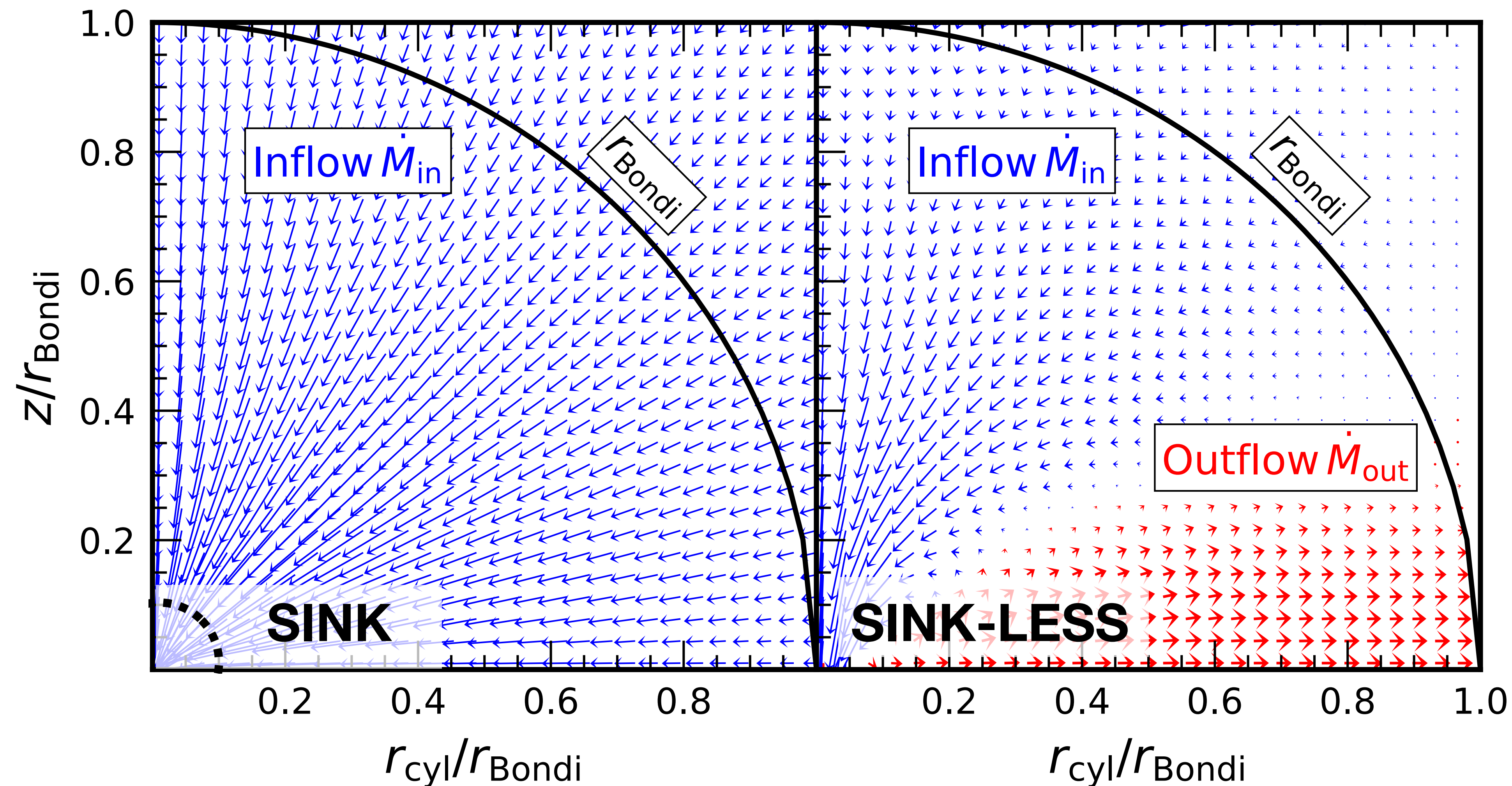
3D GLOBAL PLANET-DISK SIMULATIONS



TOP-DOWN

VERTICAL SLICE



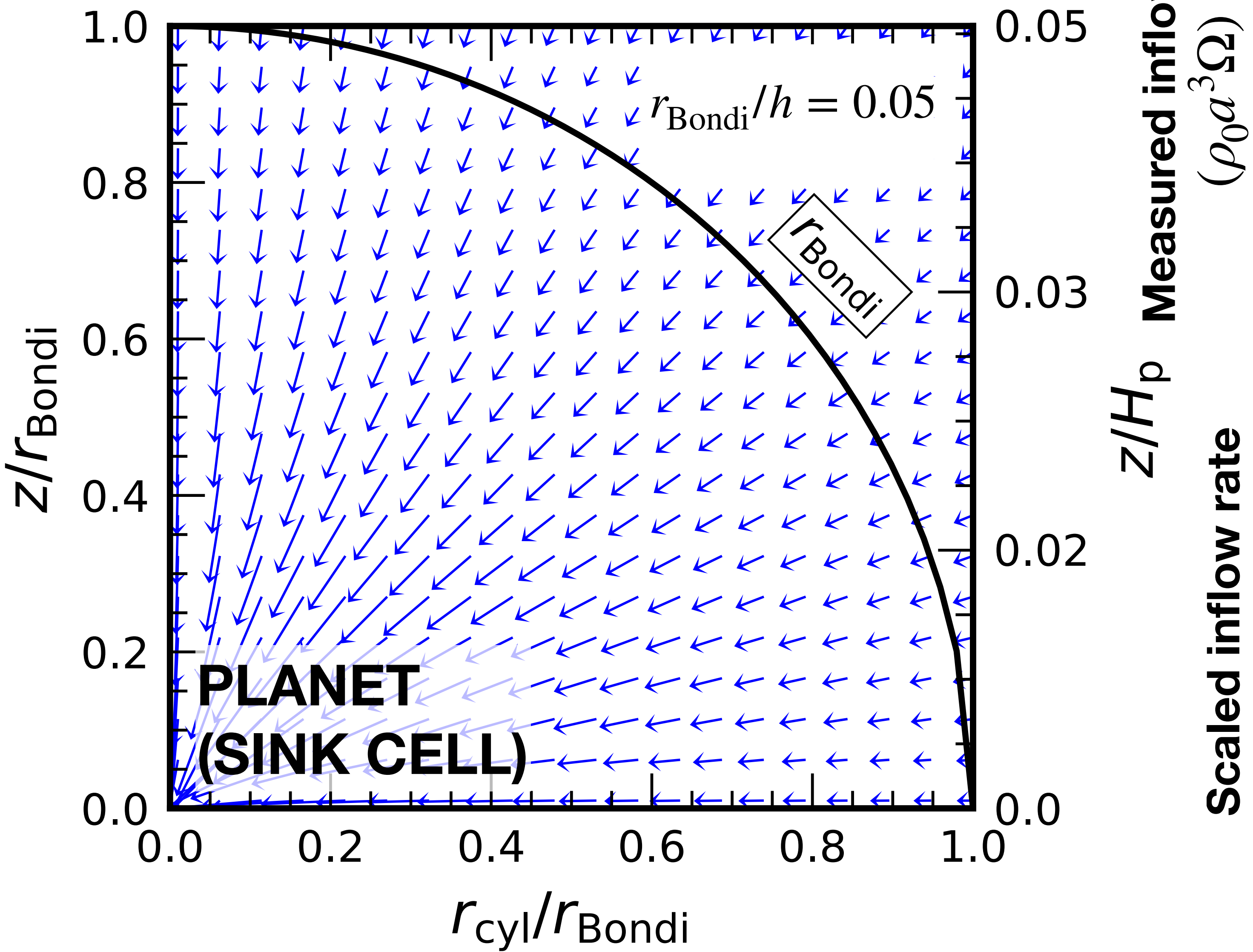


WHAT CAN WE
ROBUSTLY
MEASURE
FROM
HYDRO
SIMS?

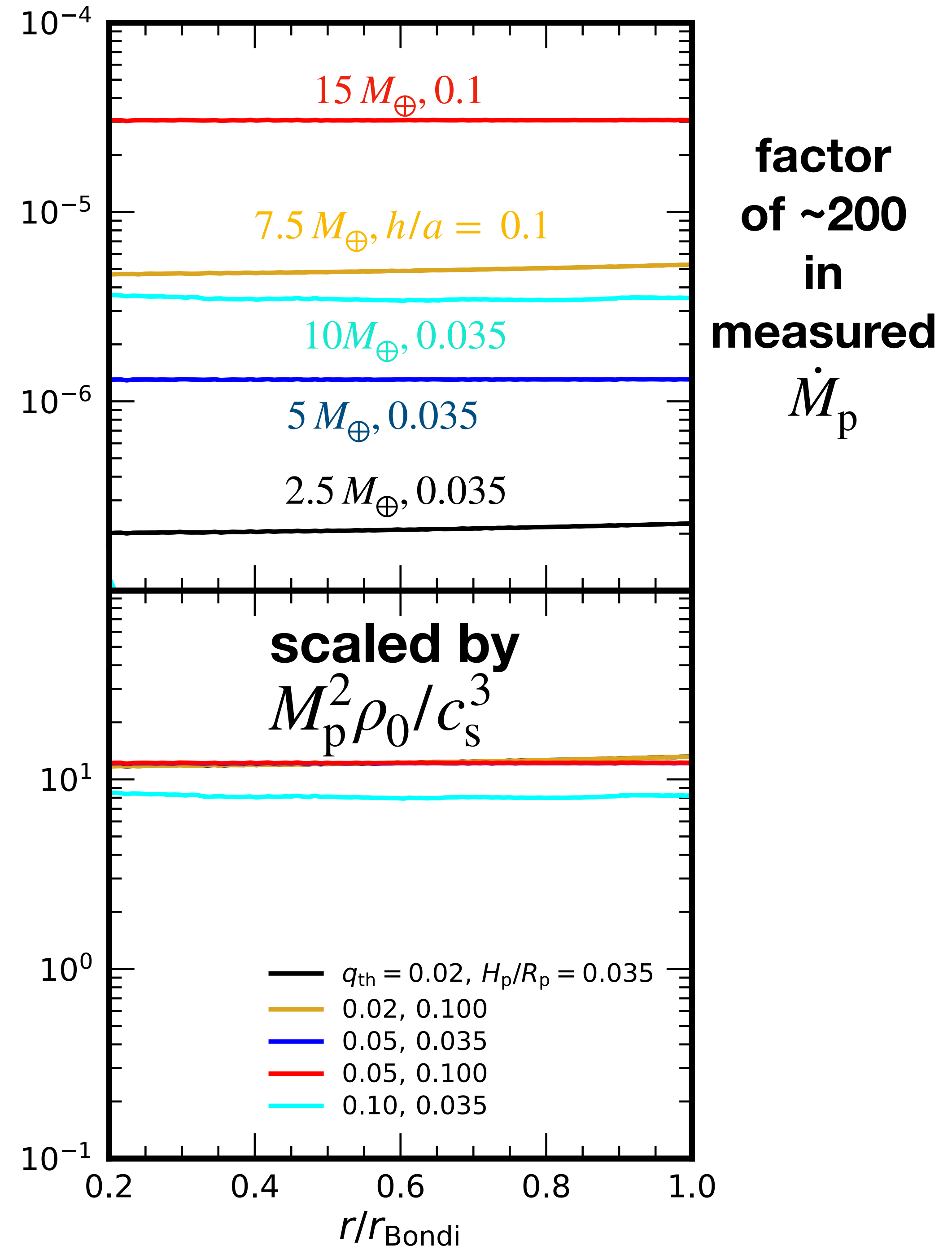
$\dot{M}_{\text{in}}$ does
not depend
on boundary
condition

“decretion disks”
are
artefacts
of boundary
Condition

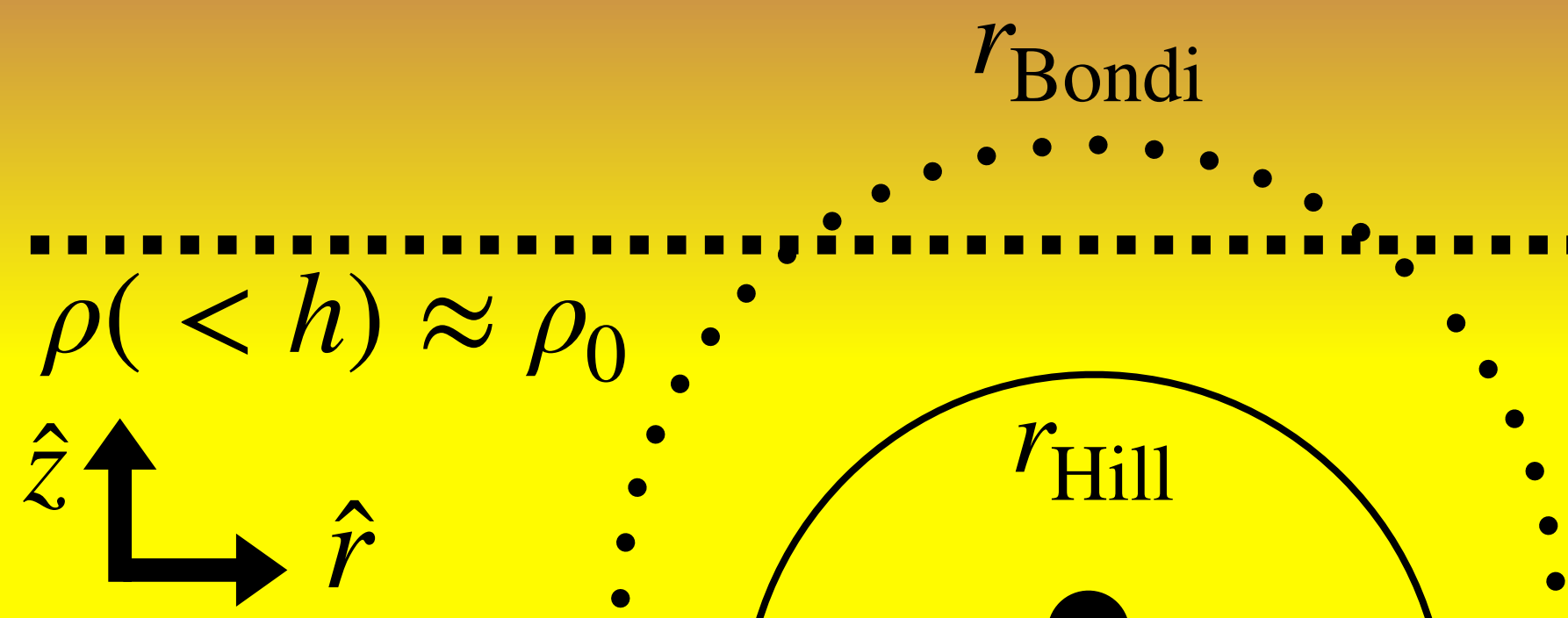
$$\dot{M}_{\text{net}} = \dot{M}_{\text{in}} - \cancel{\dot{M}_{\text{out}}} \quad \dot{M}_{\text{net}} = \dot{M}_{\text{in}} - \dot{M}_{\text{out}} \quad \boxed{= 0}$$



Measured inflow rate
 $(\rho_0 a^3 \Omega)$
 z/H_p

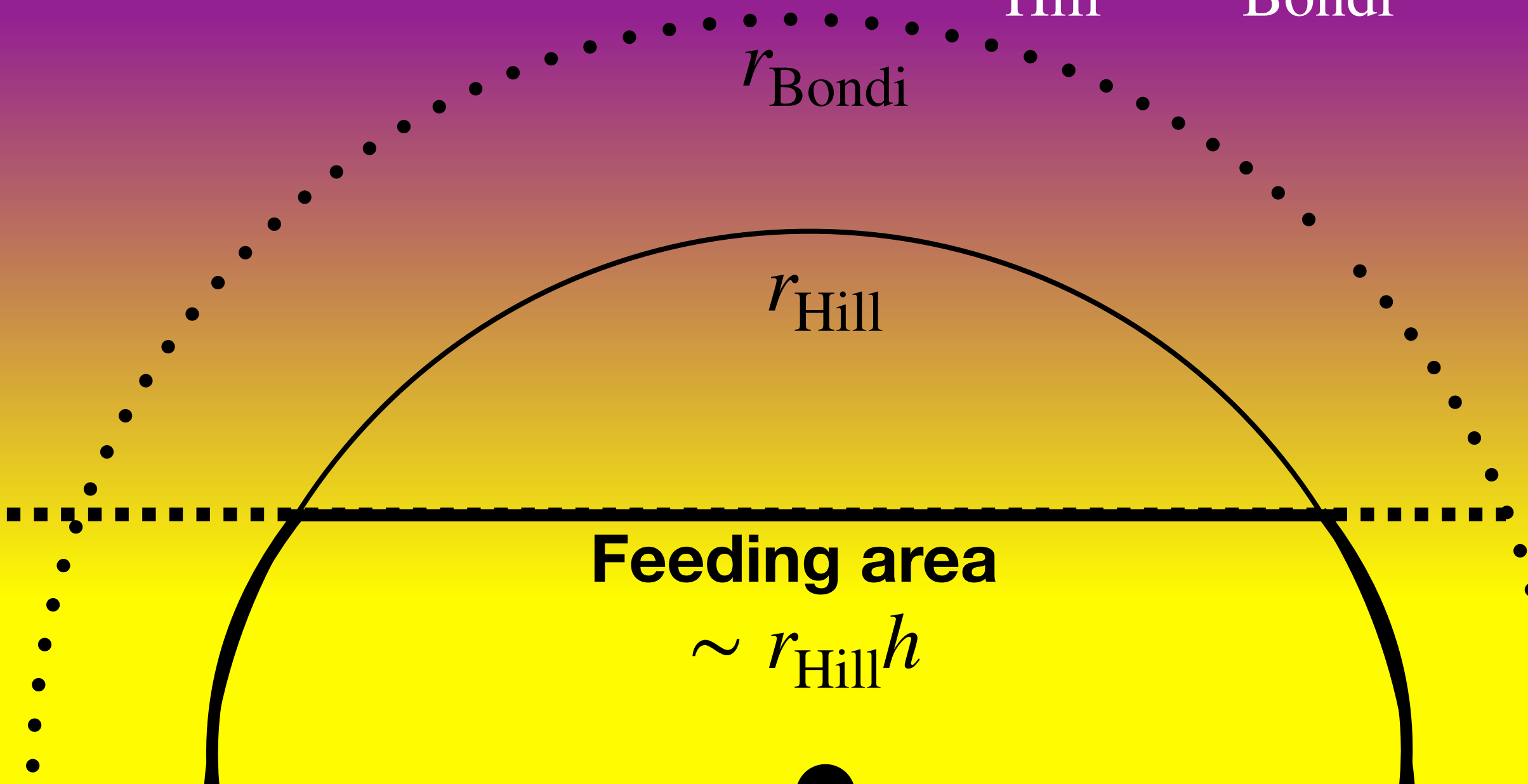


Intermediate masses : $r_{\text{Hill}} < h < r_{\text{Bondi}}$

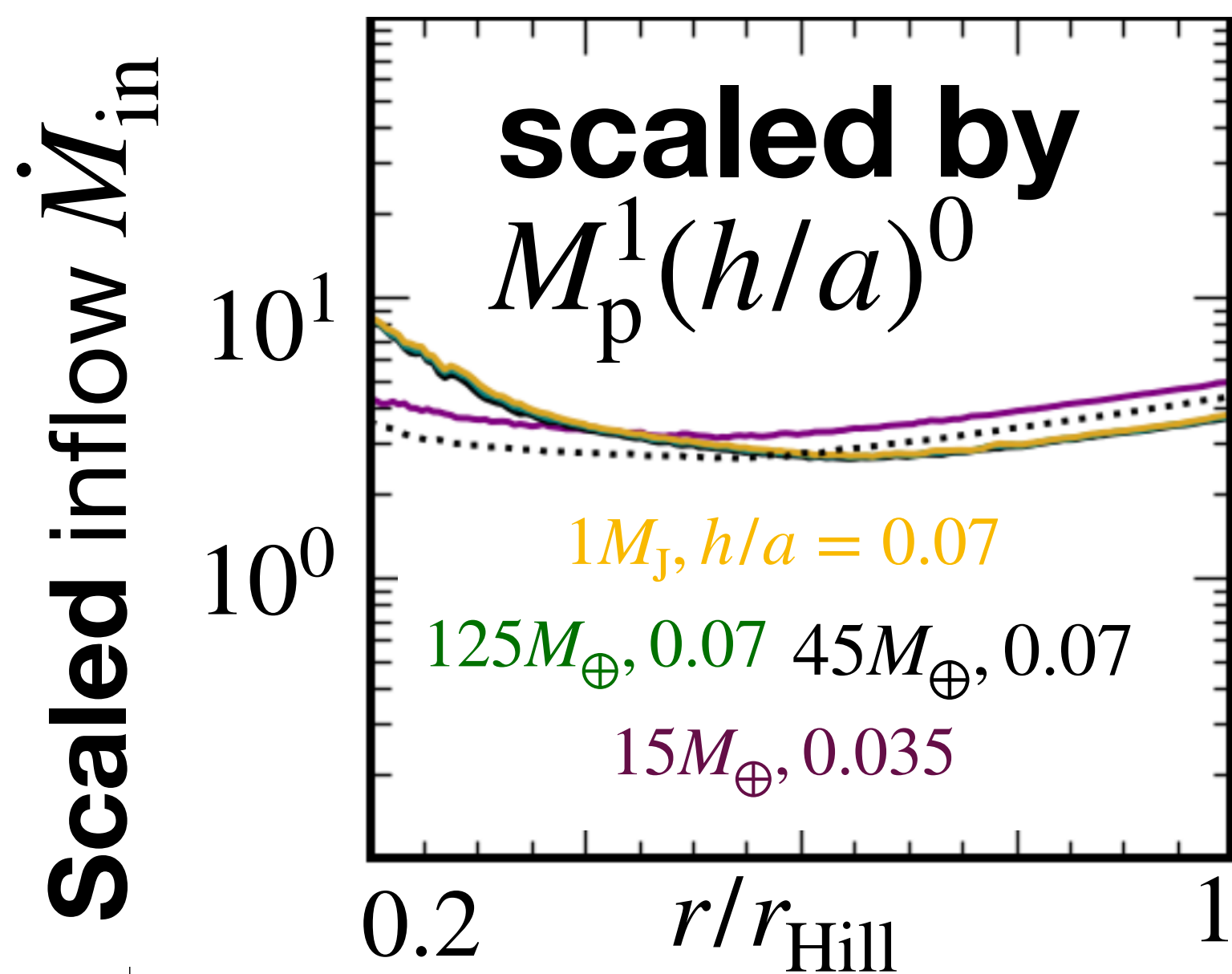


$$\dot{M} \sim \rho_0 \times (3/2)\Omega r_{\text{Hill}} \times r_{\text{Hill}}^2 \propto M_{\text{p}}^1$$

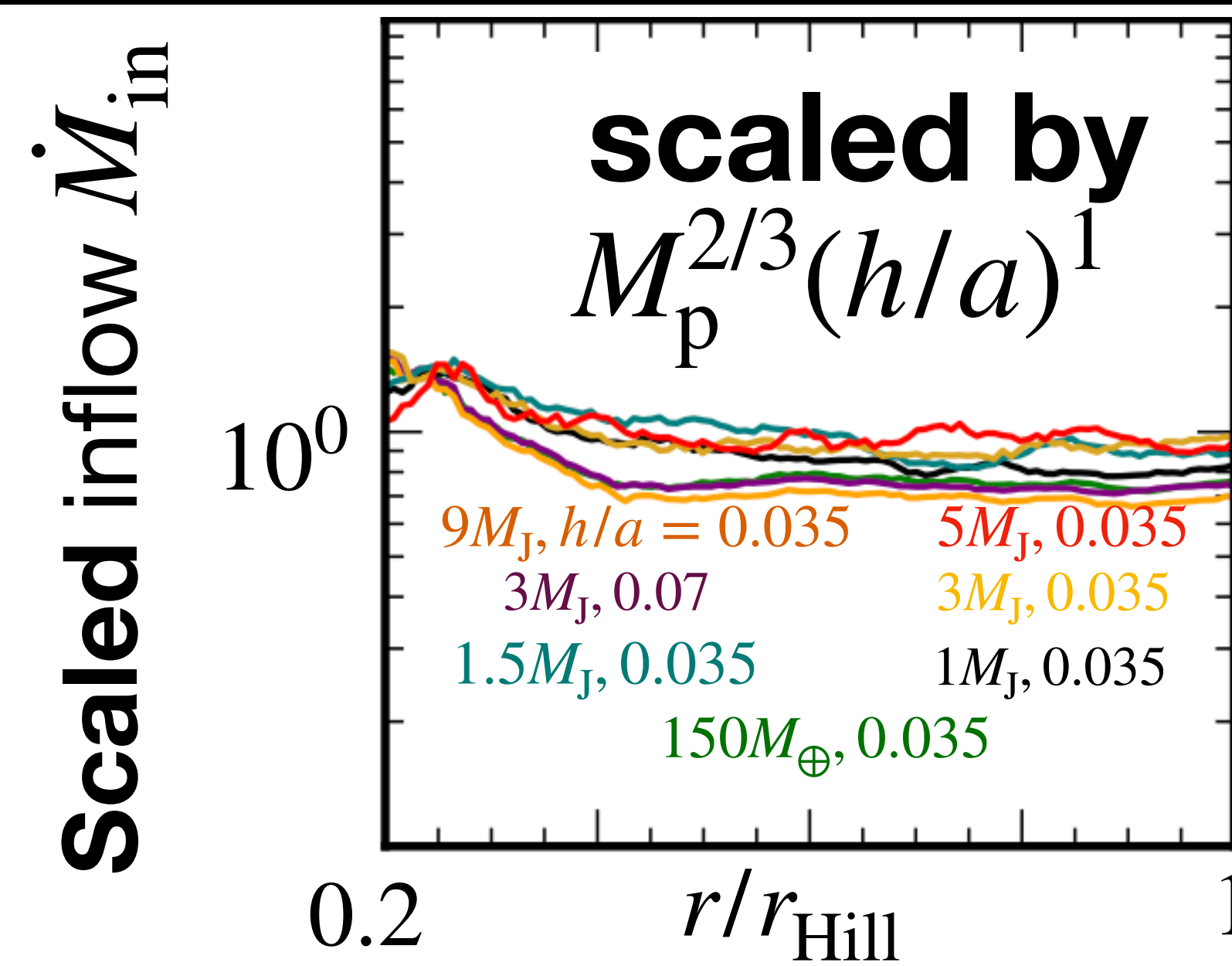
Most massive : $h < r_{\text{Hill}} < r_{\text{Bondi}}$

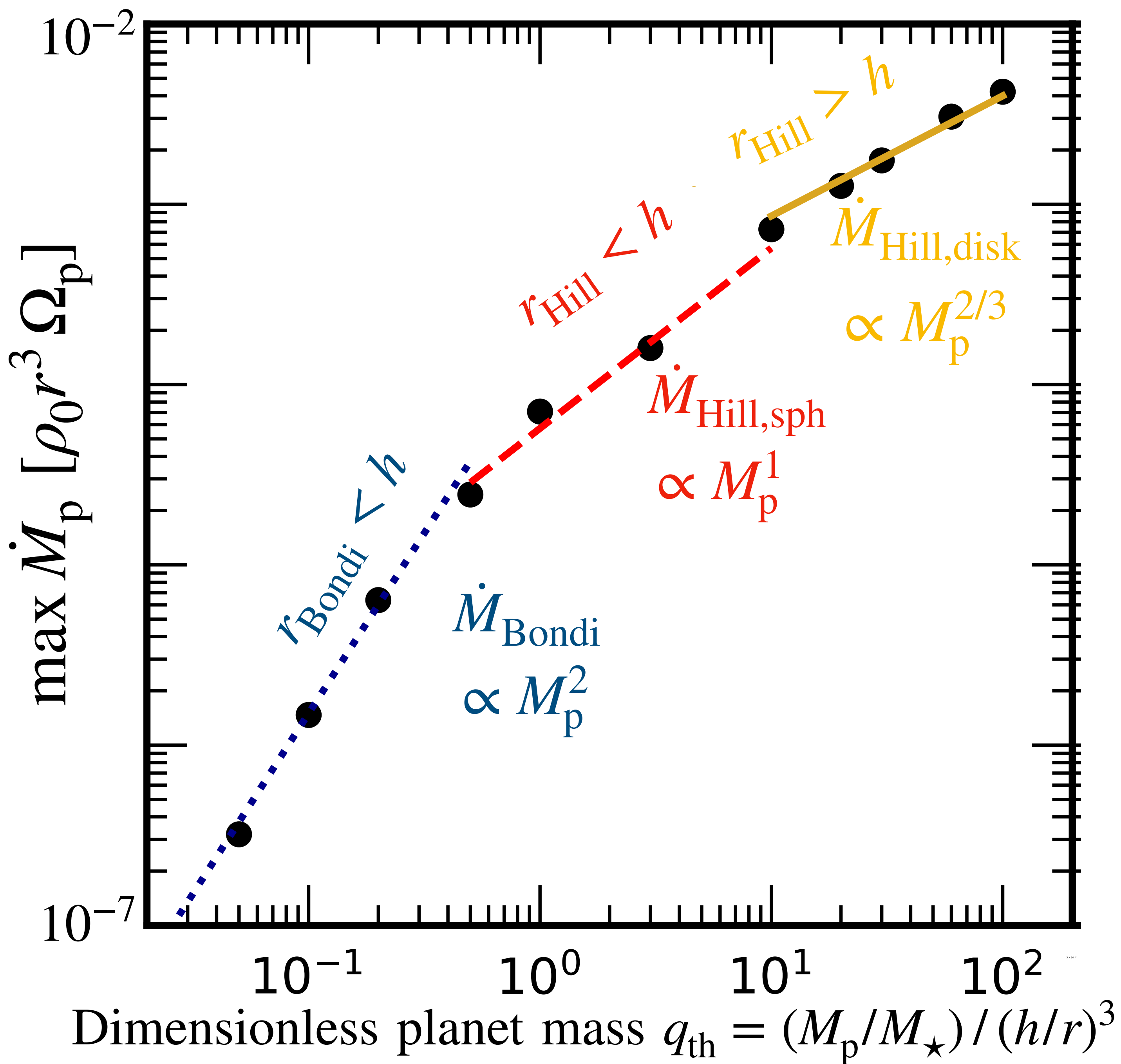


$$\dot{M} \sim \rho_0 \times (3/2)\Omega r_{\text{Hill}} \times r_{\text{Hill}} h \propto M_{\text{p}}^{2/3}$$



NC, Chiang, Fung, & Zhu 23





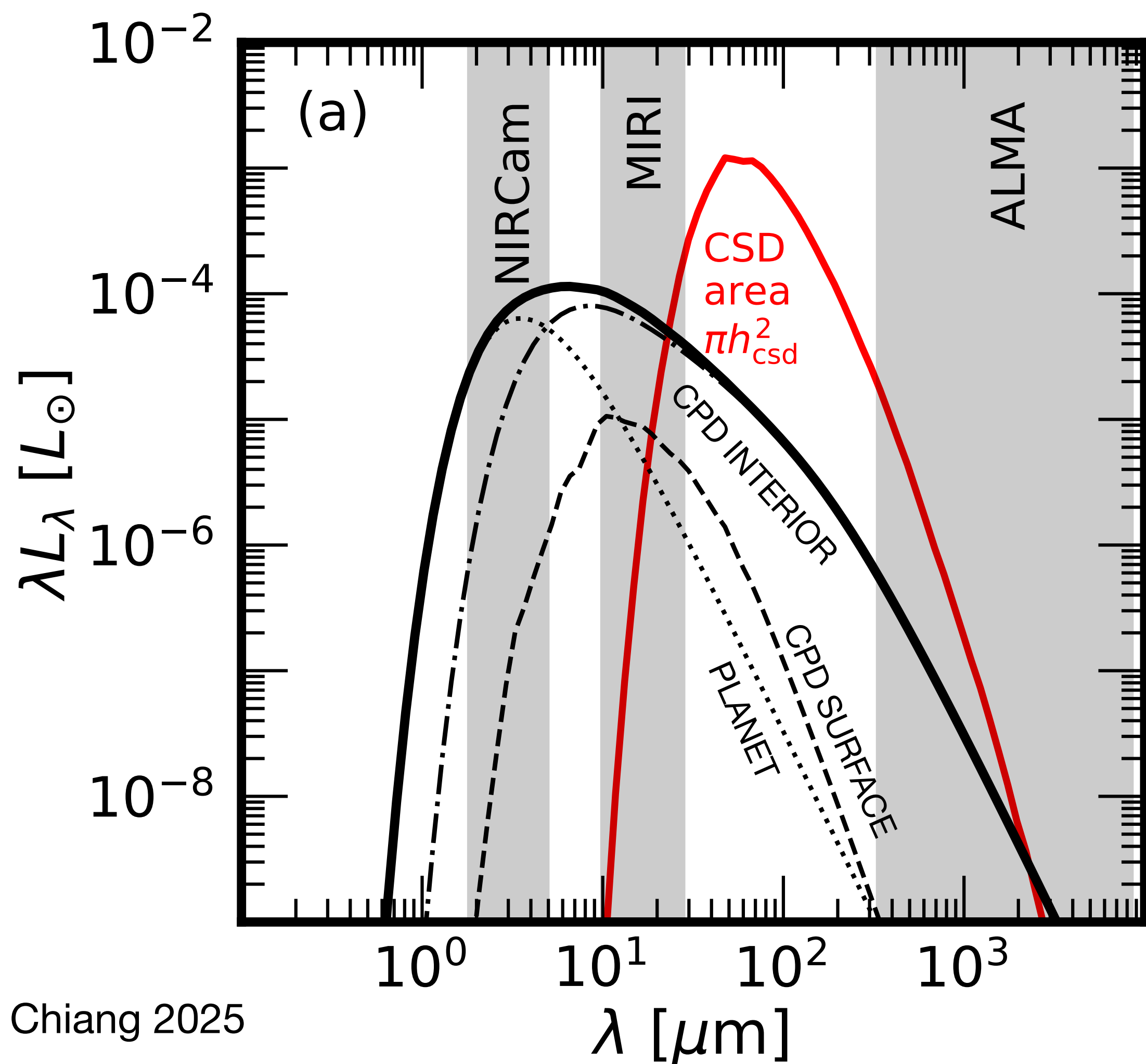
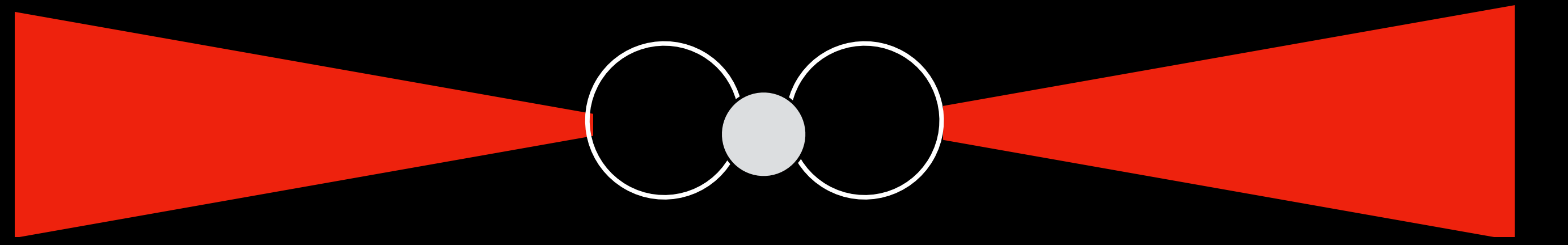
THREE REGIMES OF RUNAWAY GAS ACCRETION

$$\dot{M}_{\text{Hill,disk}} \sim r_{\text{Hill}} h \times \frac{3}{2} \Omega r_{\text{Hill}} \times \rho_0$$

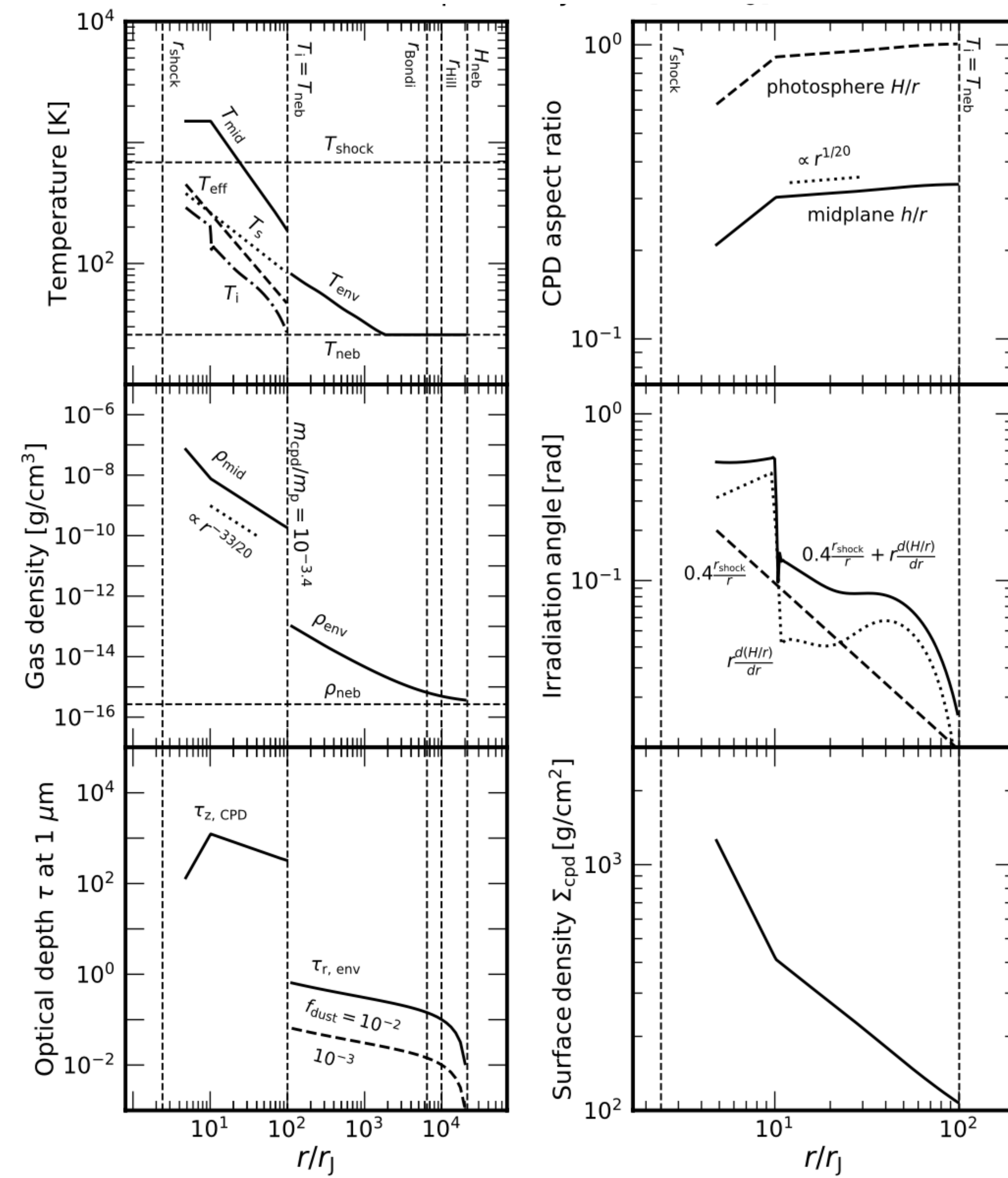
$$\dot{M}_{\text{Hill,sph}} \sim r_{\text{Hill}}^2 \times \frac{3}{2} \Omega r_{\text{Hill}} \times \rho_0$$

$$\dot{M}_{\text{Bondi}} \sim r_{\text{Bondi}}^2 \times c_s \times \rho_0$$

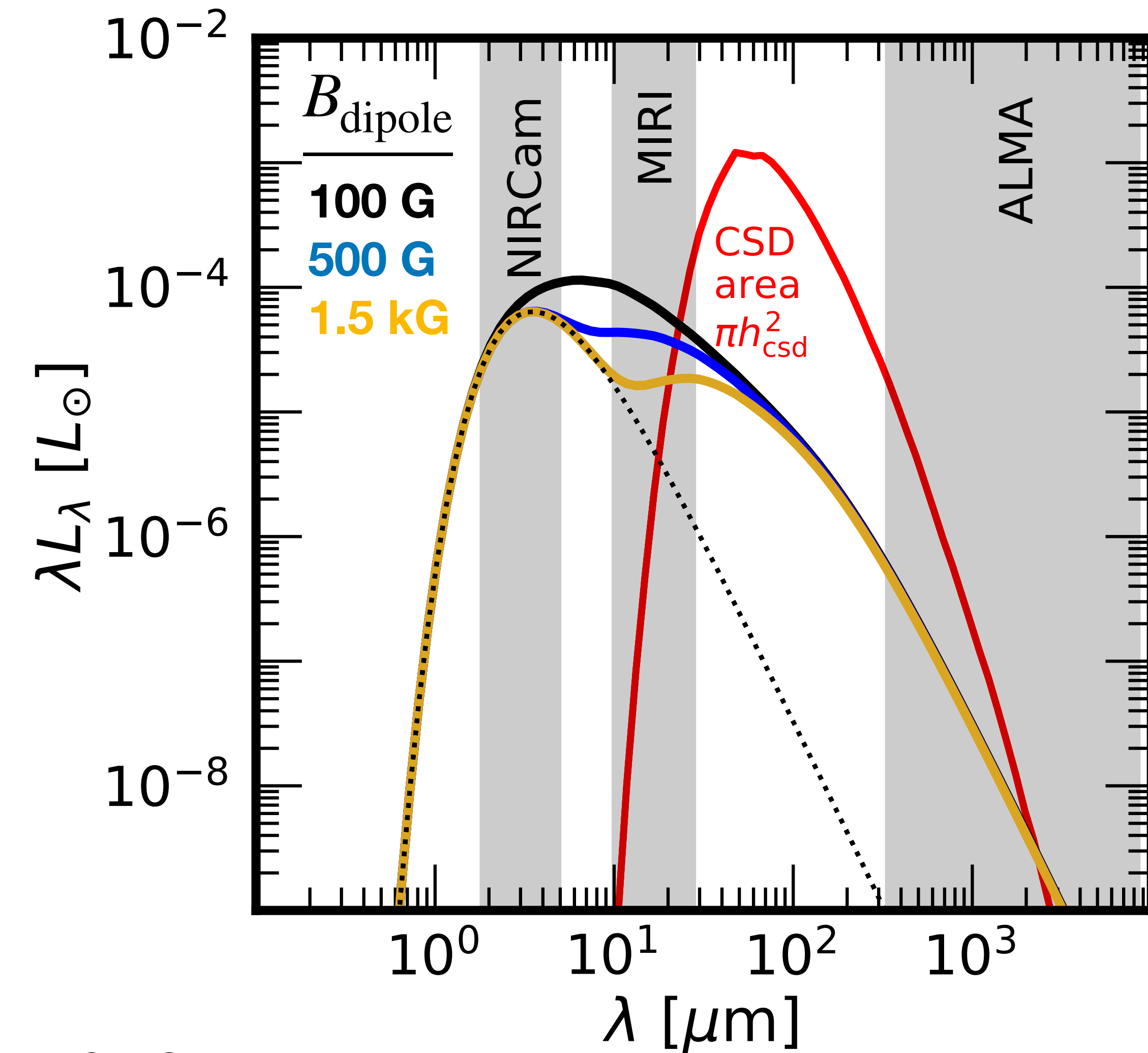
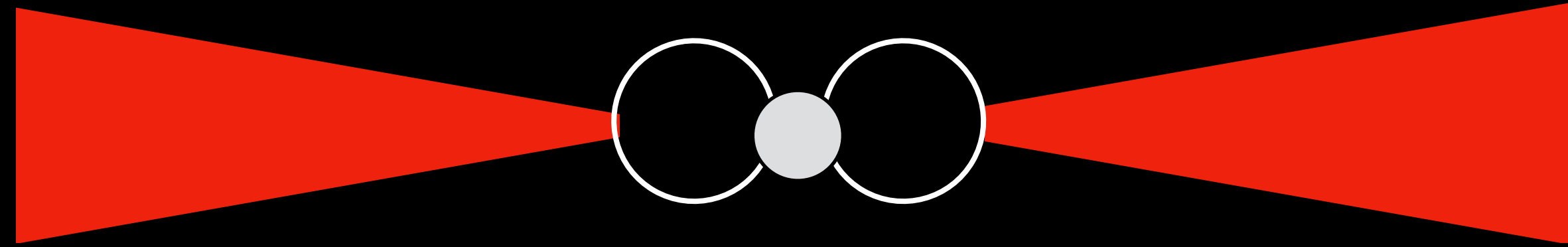
SEDs OF CIRCUMPLANETARY DISKS



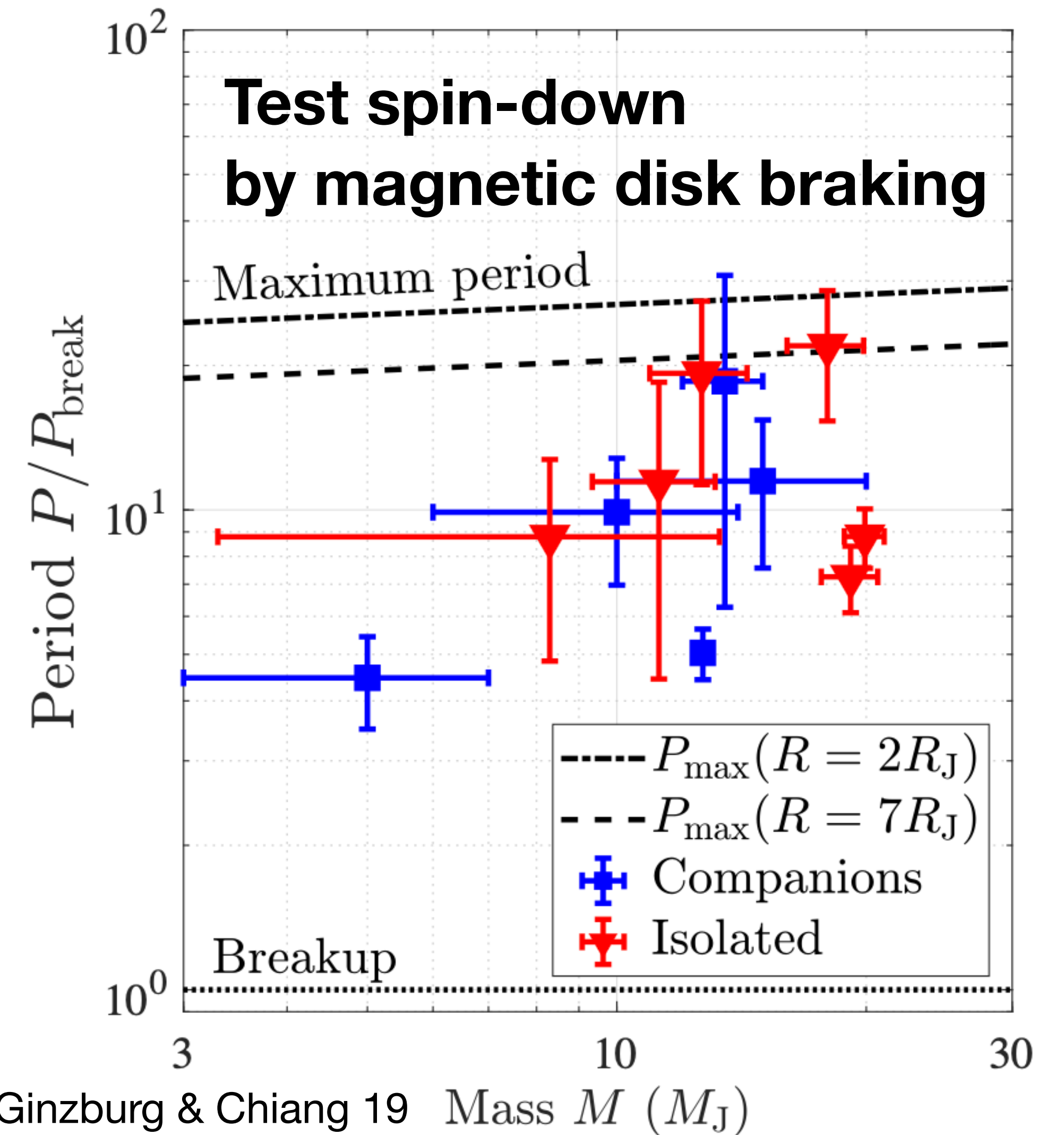
CPDs heated by irradiation and dissipation



SEDs OF MAGNETICALLY TRUNCATED CPDs

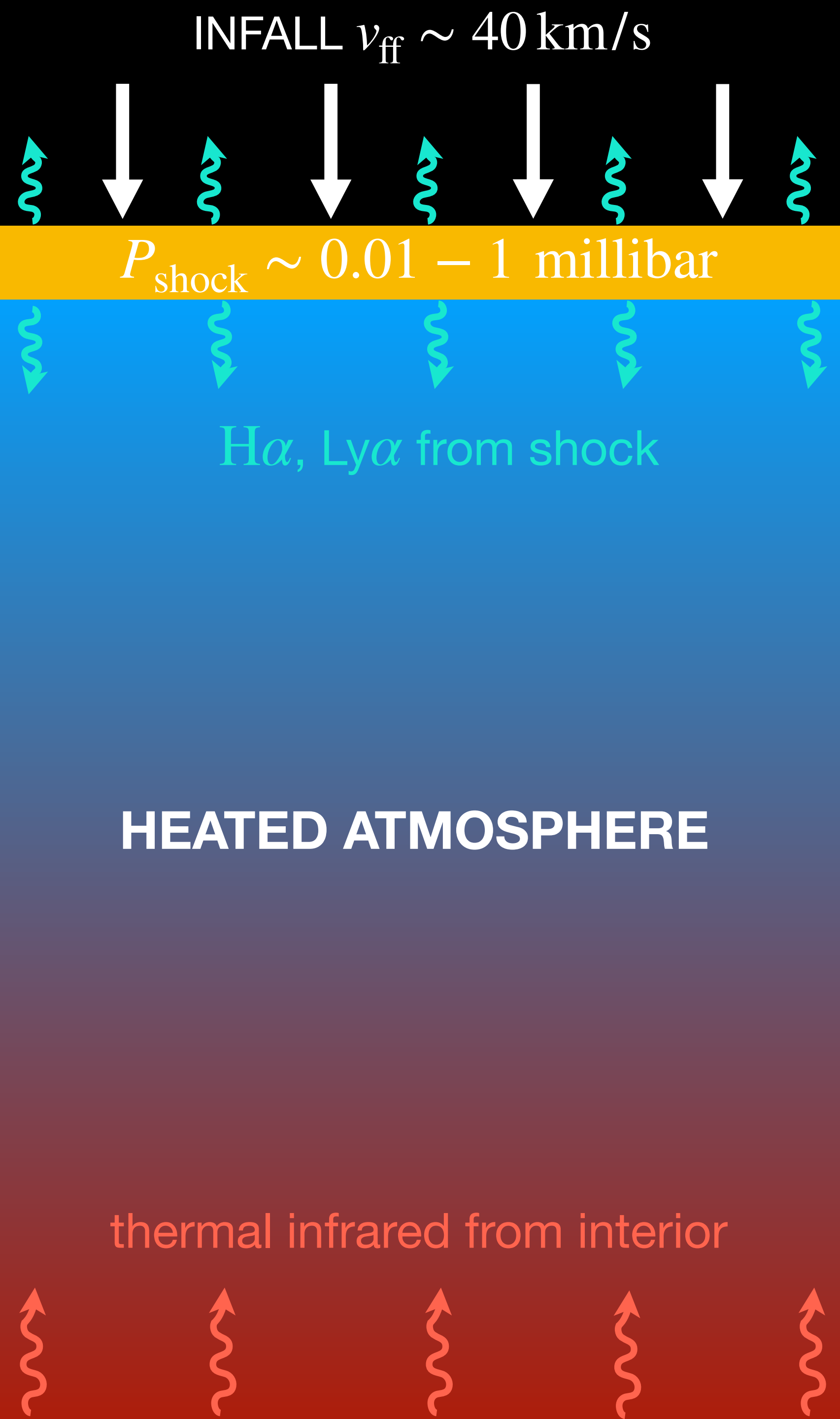


NC & Chiang 2025



Bryan+ 18

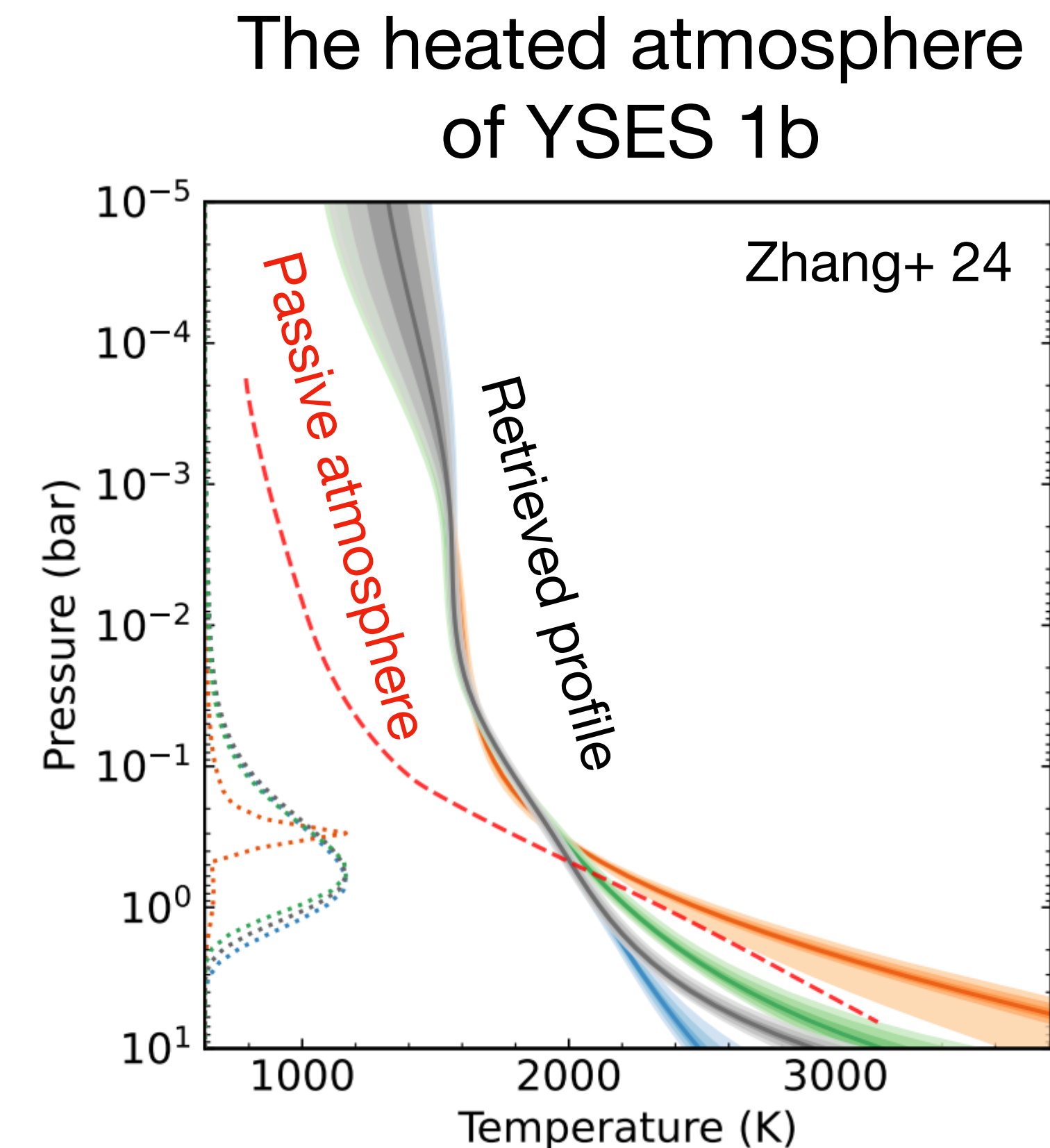
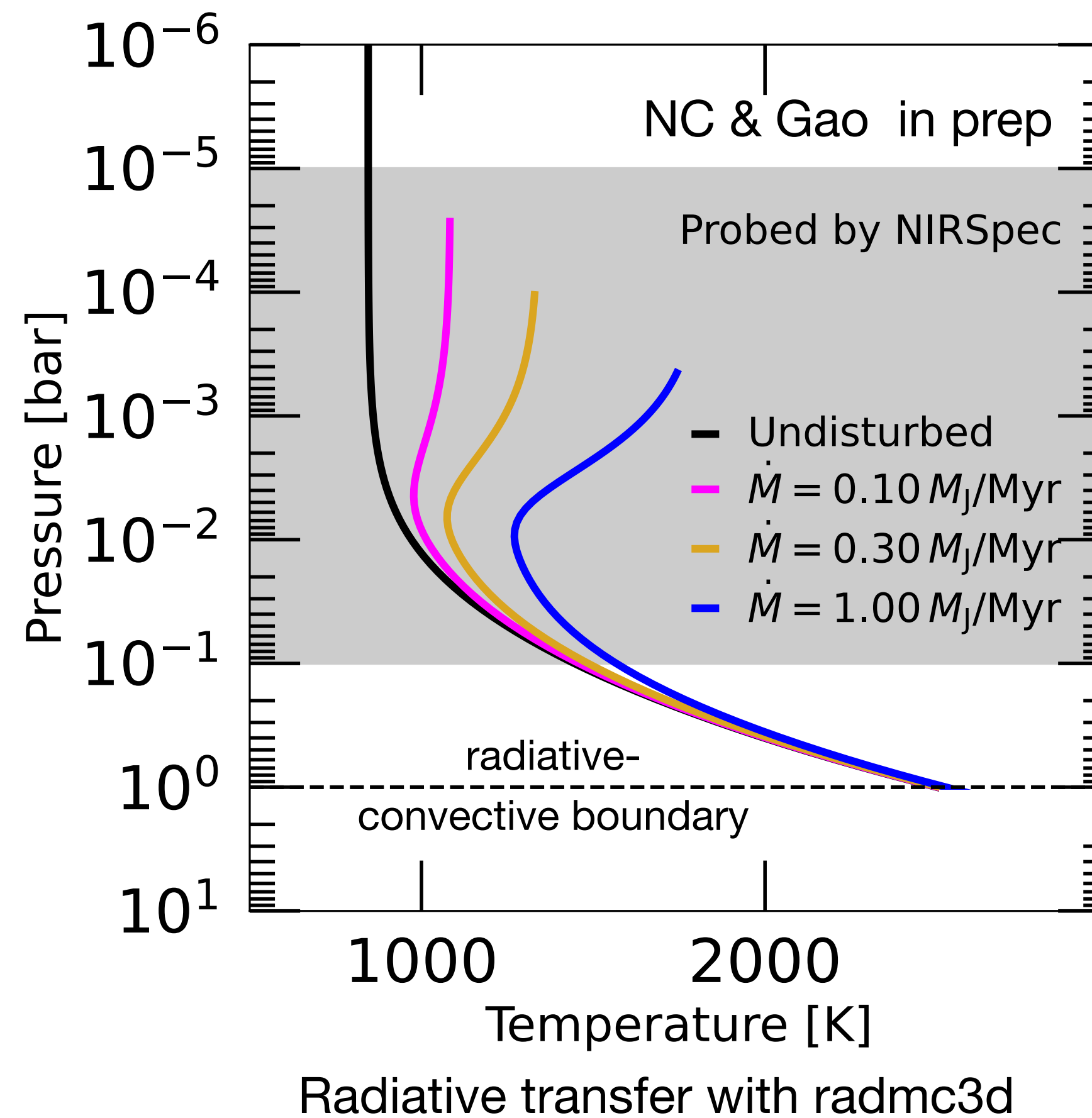
Batygin 18, Ginzburg & Chiang 19



ACCRETIONALLY HEATED ATMOSPHERES

$$T_{\text{atm}}^4(\tau) \approx \frac{3}{4} T_{\text{int}}^4 \left[\frac{2}{3} + \tau \right] + \frac{3}{4} \frac{G M \dot{M}}{4 \pi R^3 f_{\text{fill}} \sigma} q e^{-q\tau}$$

$$\frac{dP}{d\tau} = g / \kappa_{\text{IR}} \quad q \equiv \kappa_{\text{shock}} / \kappa_{\text{IR}} \quad \text{e.g. Guillot 2010}$$



SUPPLY-LIMITED ACCRETION IN PDS 70

Accretion rate from hydro

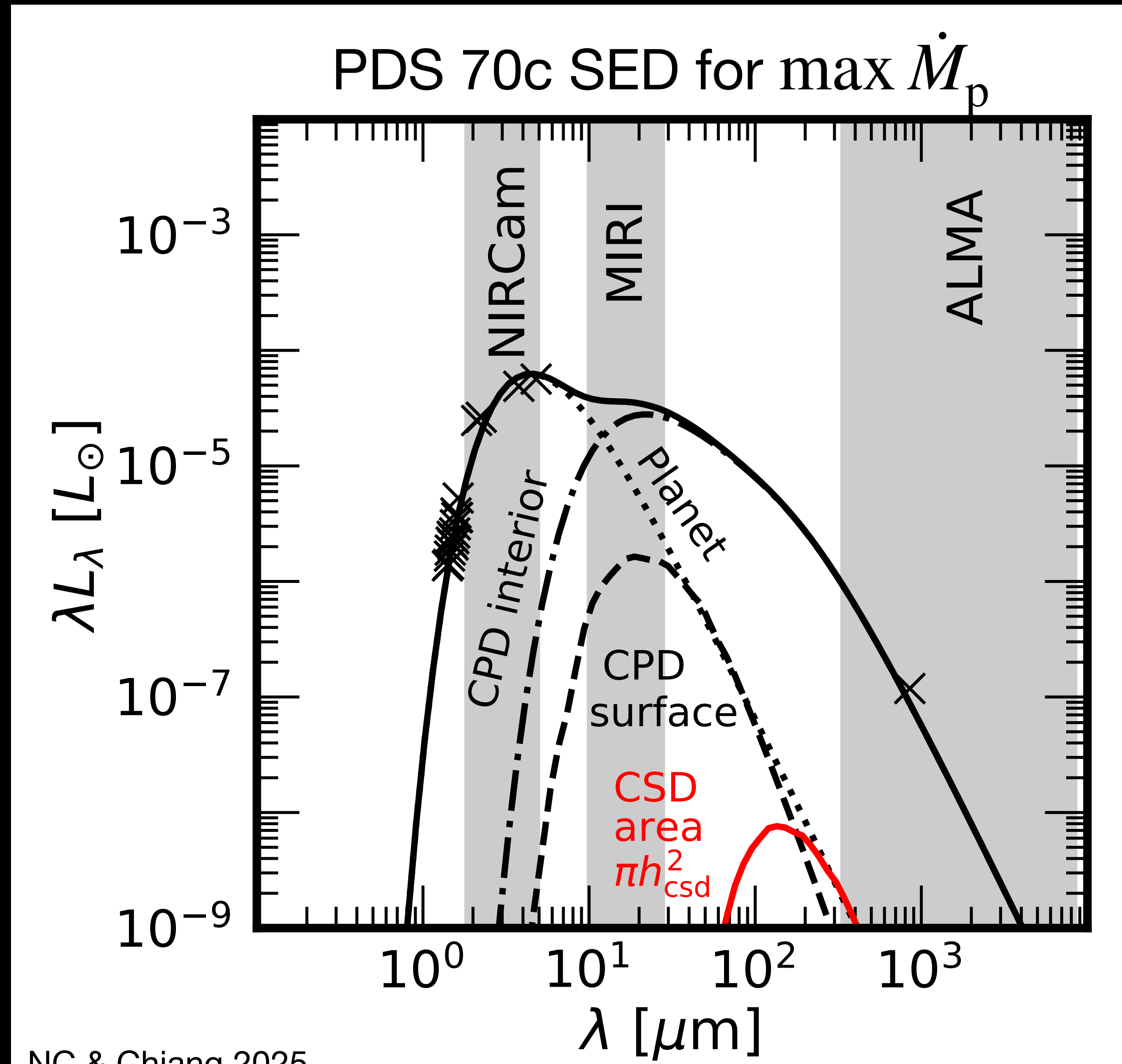
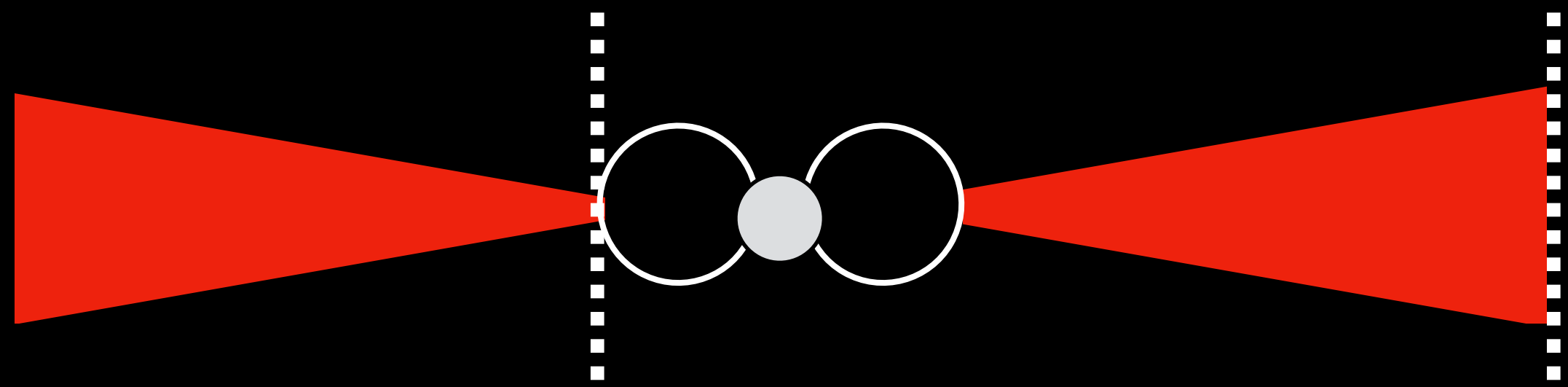
$$\begin{aligned}\max \dot{M}_p &\sim r_{\text{Hill}} h \times (3/2) \Omega r_{\text{Hill}} \times \rho_{\text{gas}} \\ &\sim 1 M_J / \text{Myr} \\ &\sim M_p / t_{\text{age}}\end{aligned}$$

from ALMA

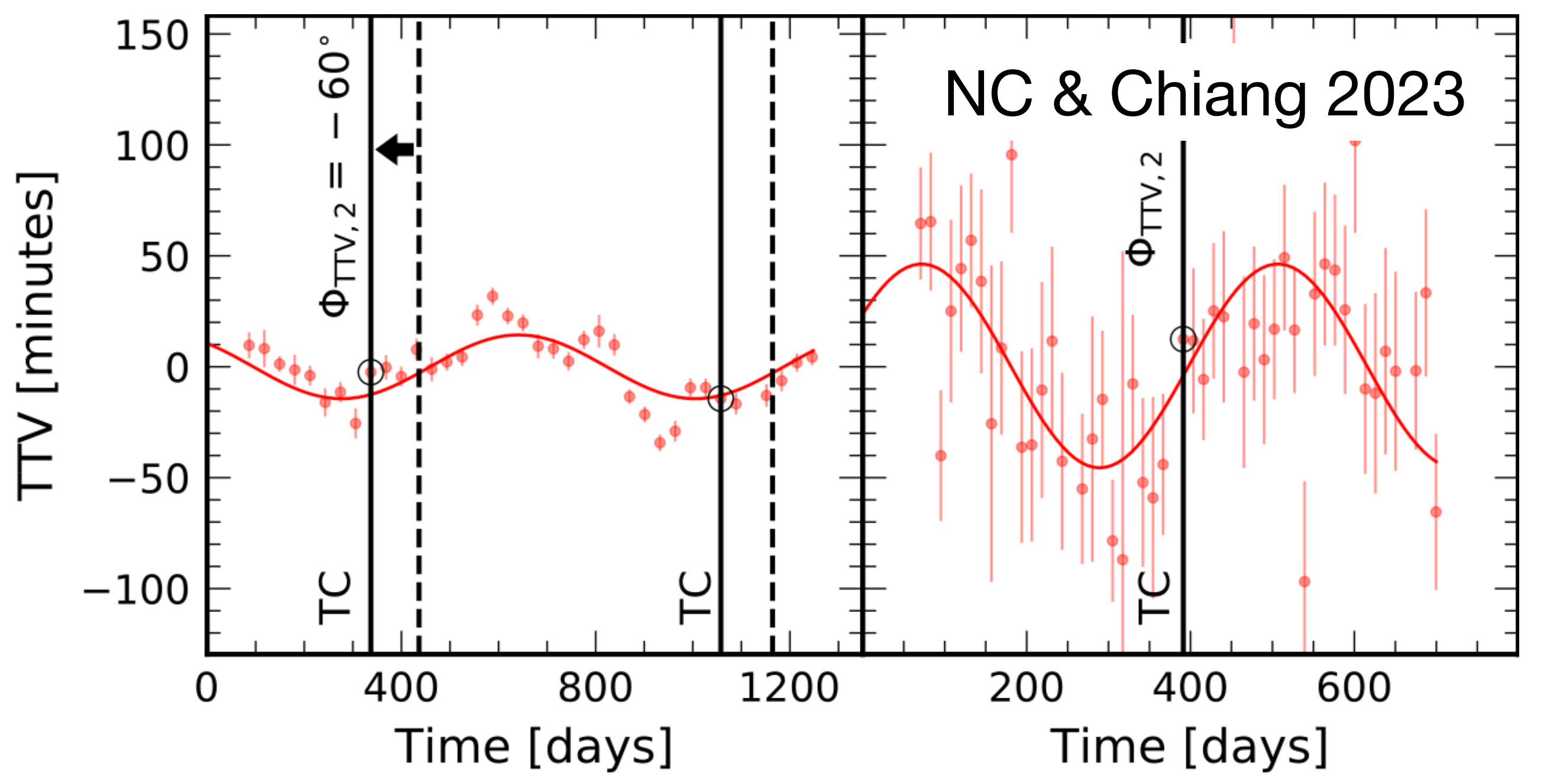
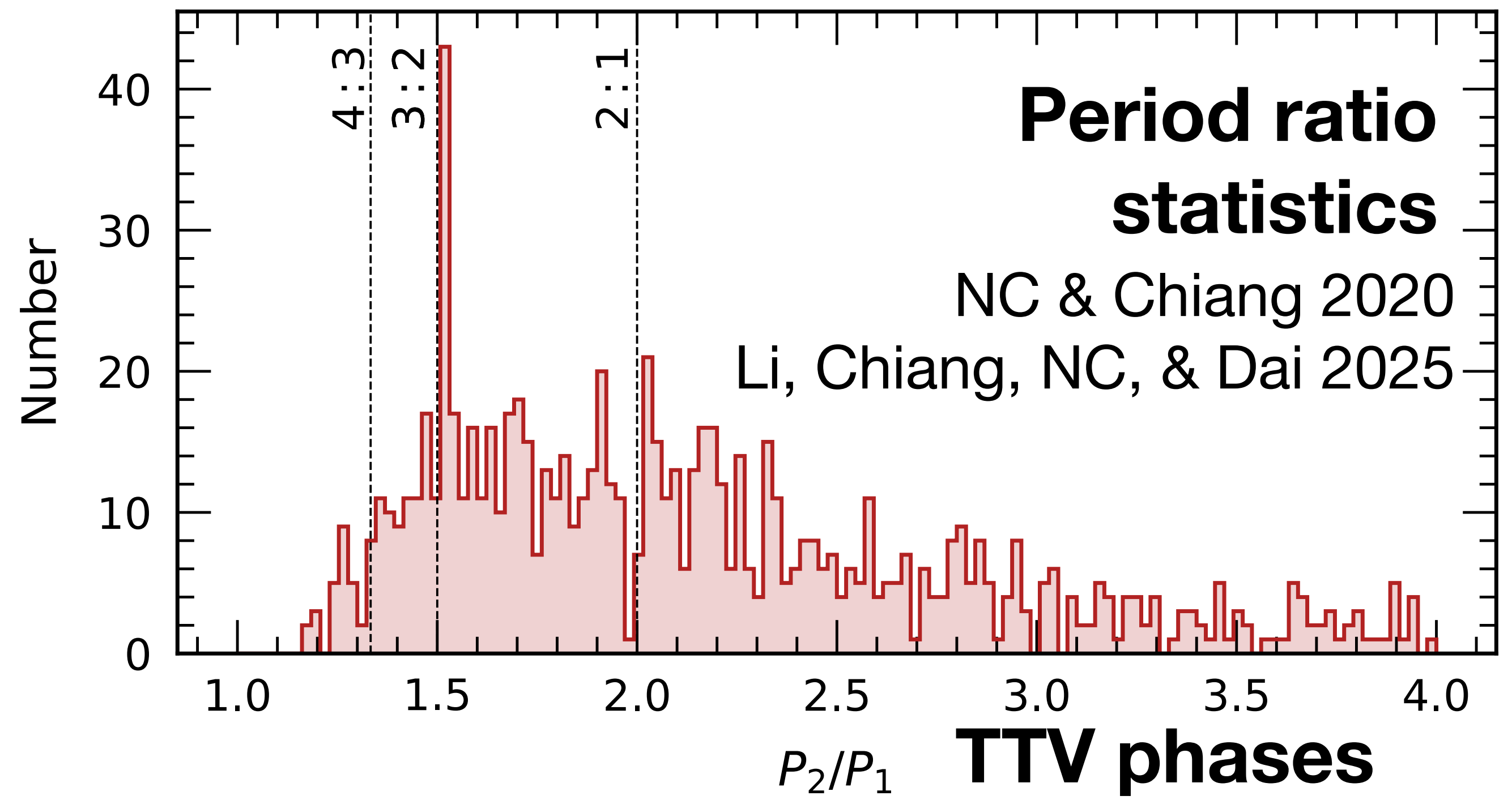
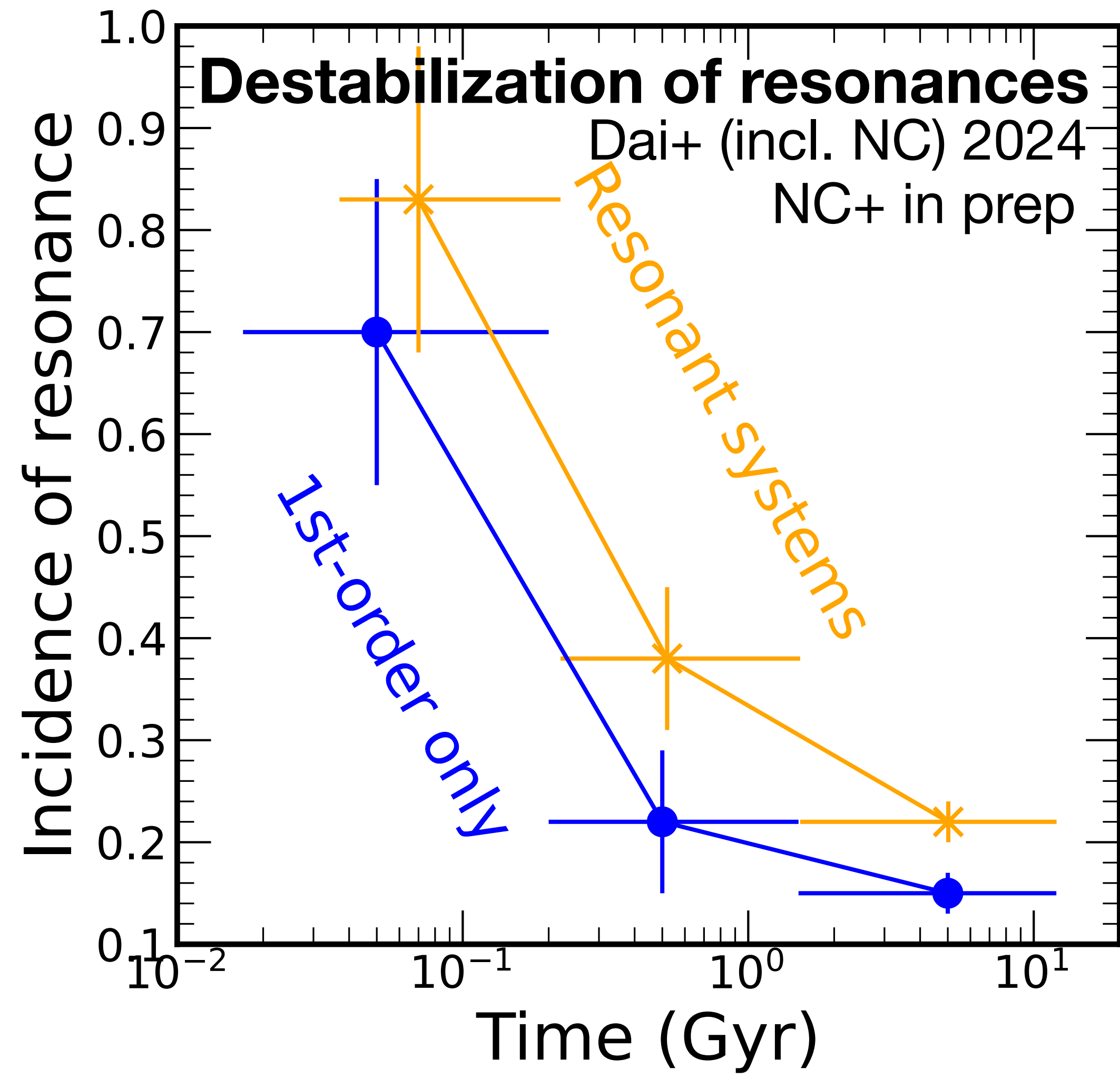
cf. 100x lower $\dot{M}_p$ from $\text{H}\alpha$
e.g. Haffert+ 19

Circumplanetary disk from SED

$$\begin{aligned}&> 20 R_J \\ (B_{\text{dipole}} &> 200 \text{ G}) \quad \sim 0.1 R_{\text{Hill}}\end{aligned}$$

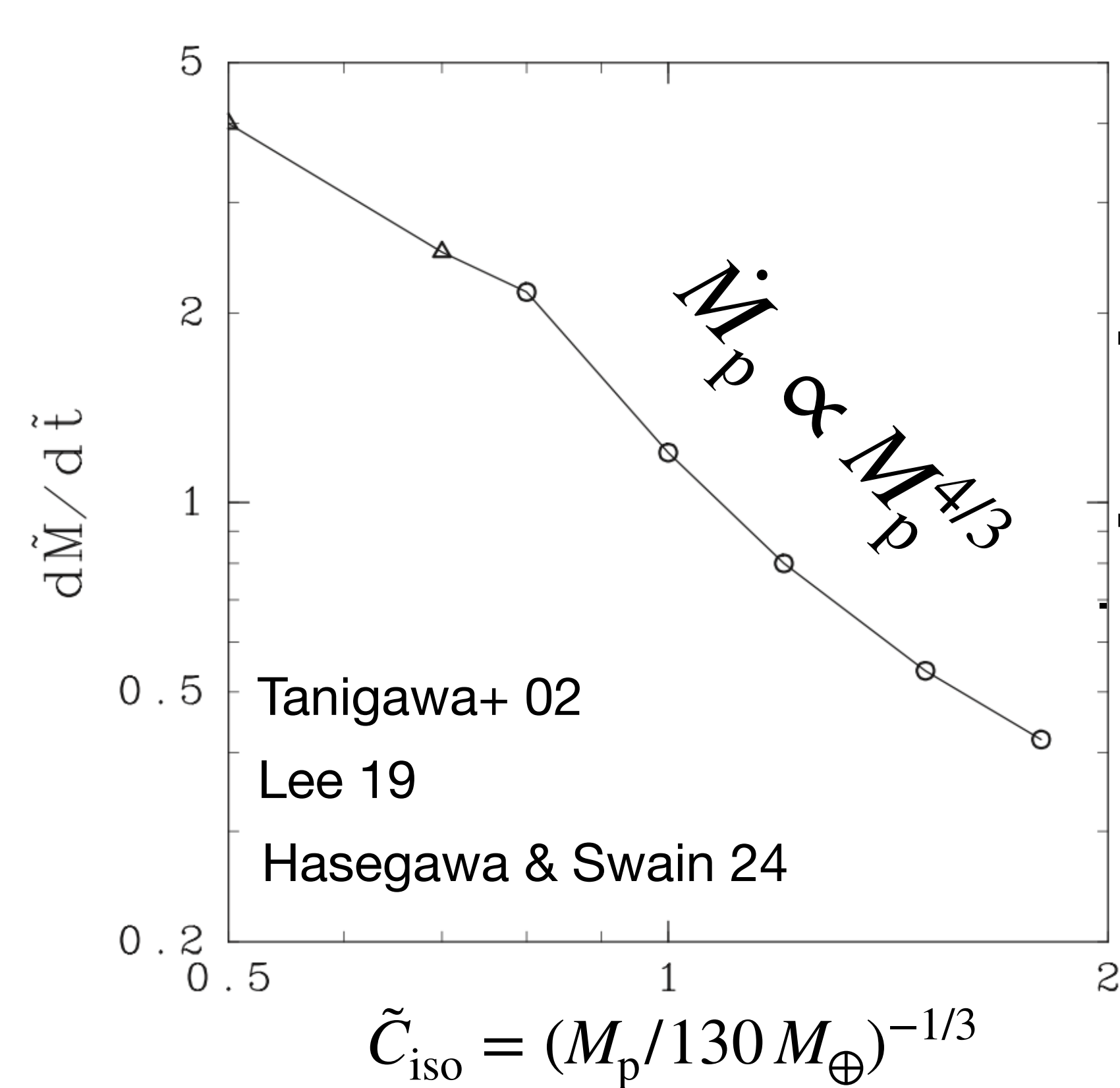


RESONANCES AS SIGNPOSTS OF SUPER-EARTH ASSEMBLY

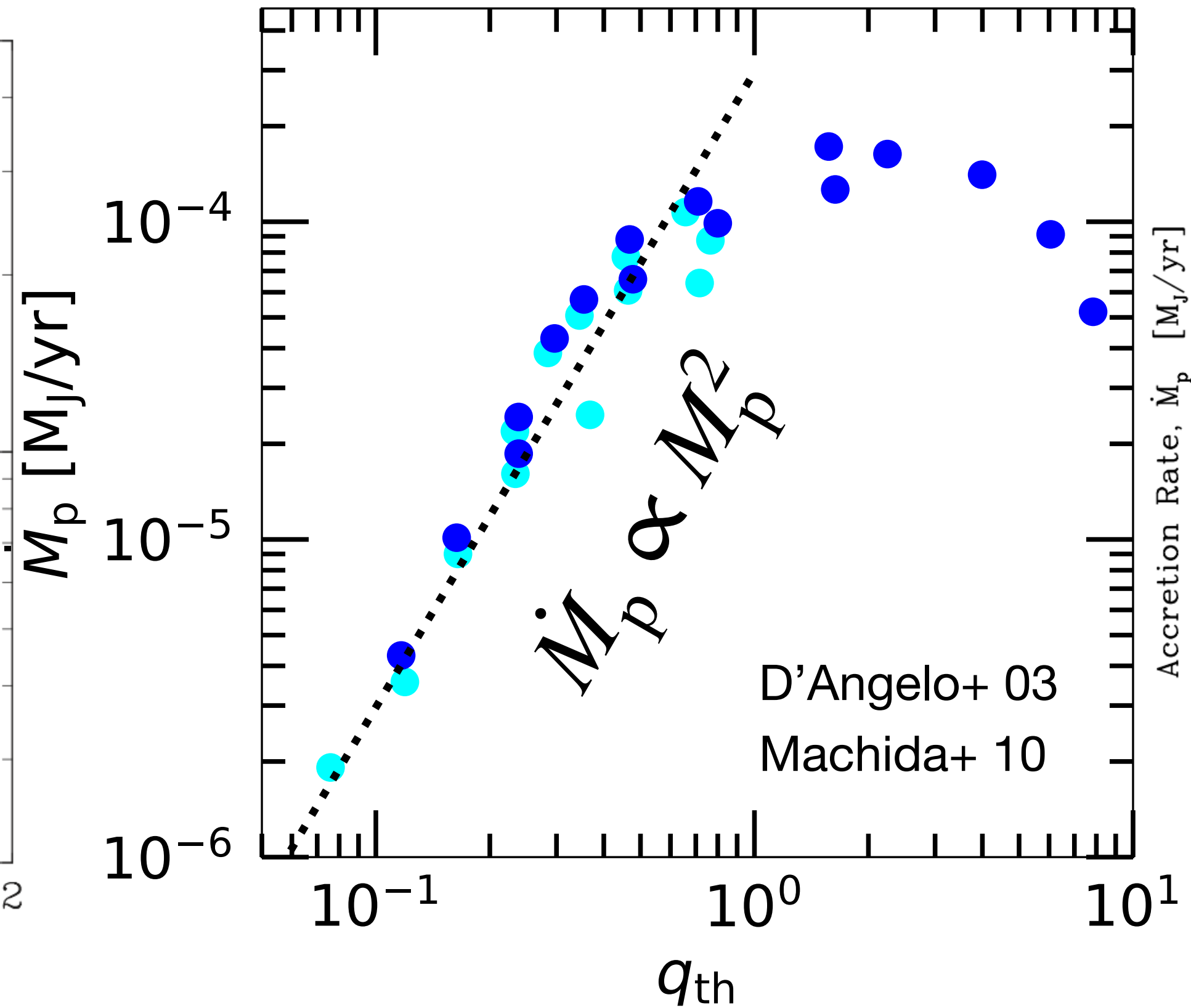


EXTRA SLIDES

HOW FAST DO PLANETS ACCRETE GAS?

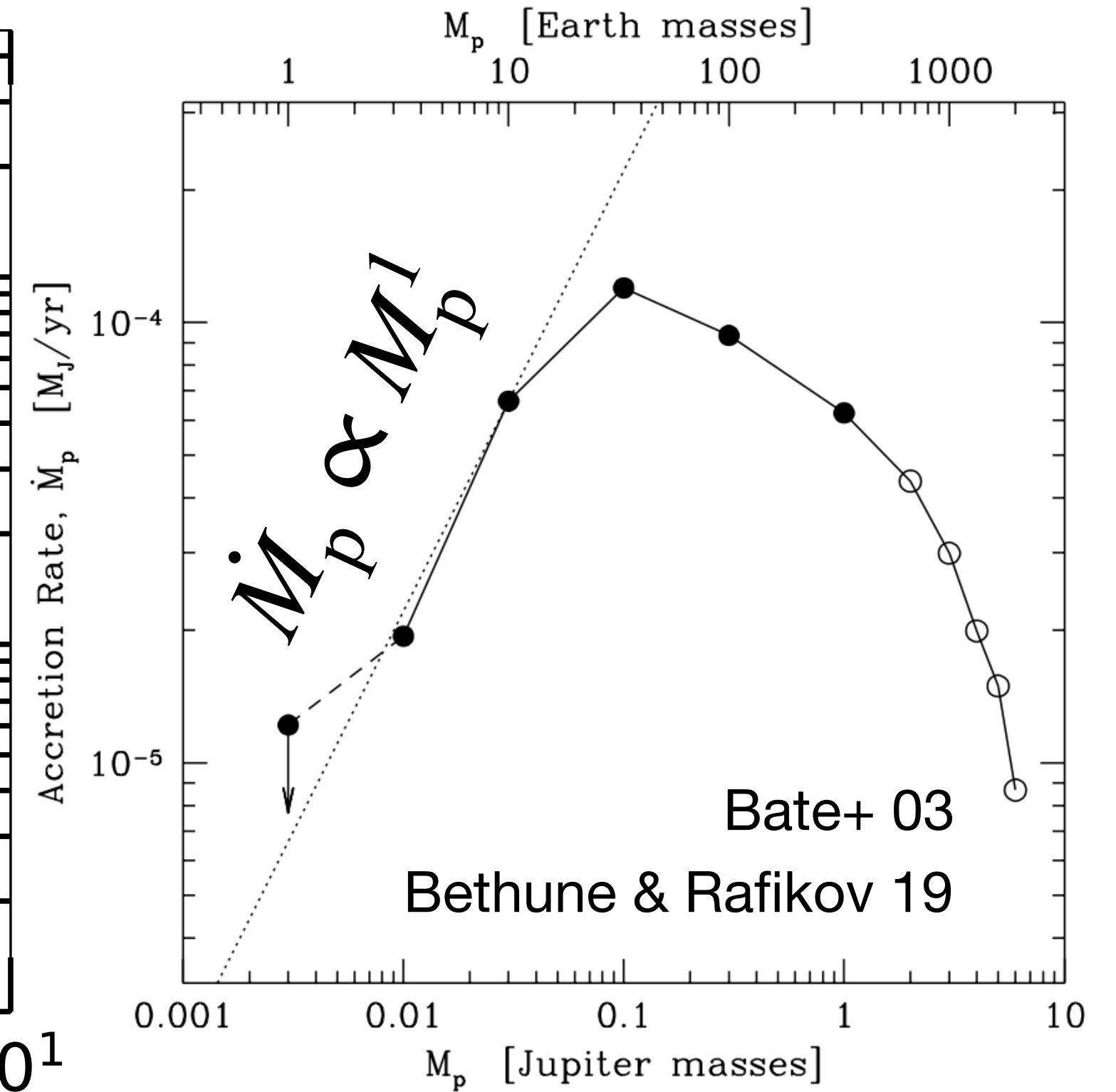


2D isothermal shearing box



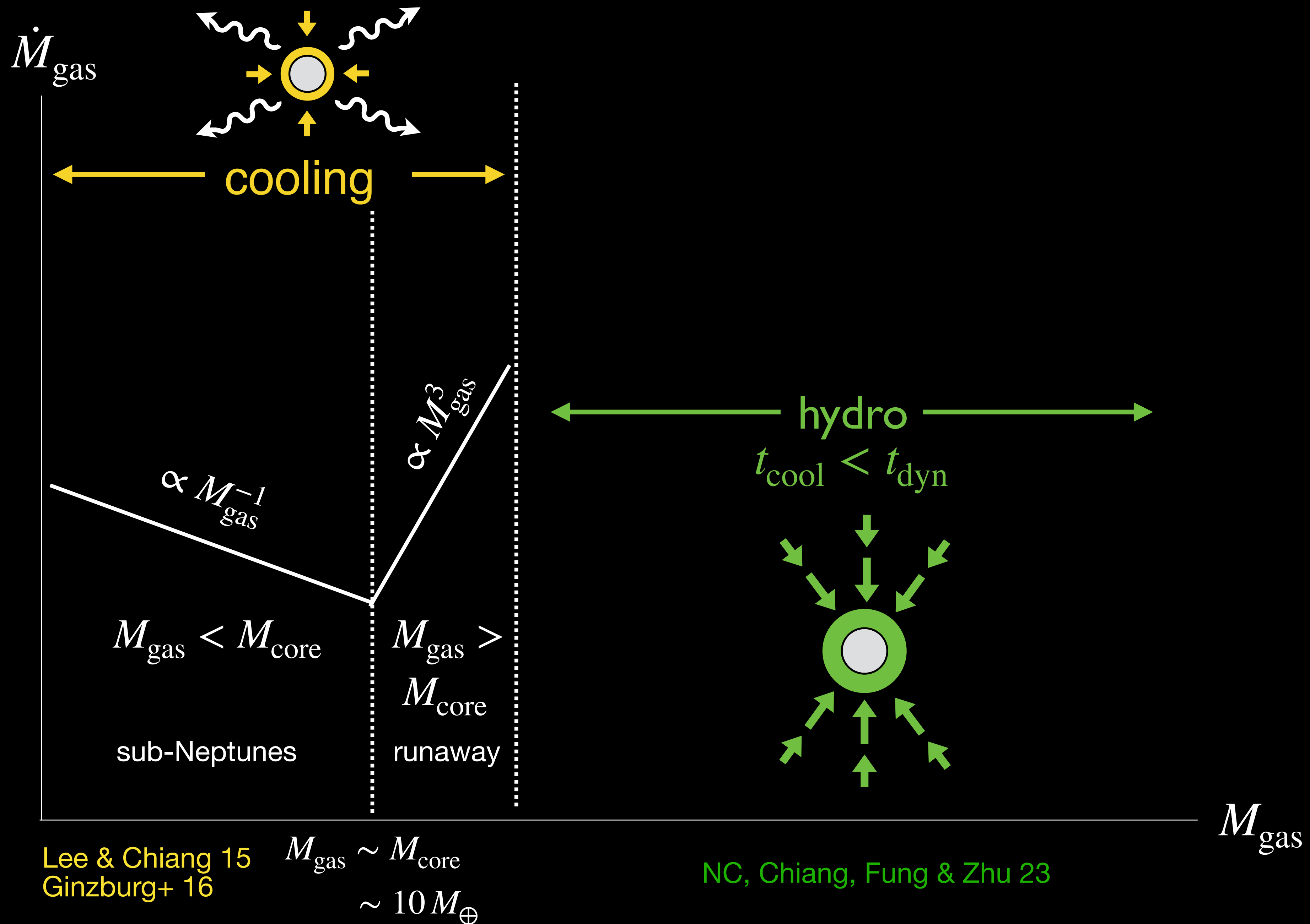
3D isothermal global

also $\dot{M}_p \propto M_p^{2/3}$
Zhu+ 12



3D isothermal global

HOW FAST DO PLANETS ACCRETE GAS?



THE MAXIMUM ACCRETION RATE OF A PROTOPLANET

$$\dot{M}_{\text{Bondi}} \sim r_{\text{Bondi}}^2 \times c_s \times \rho$$

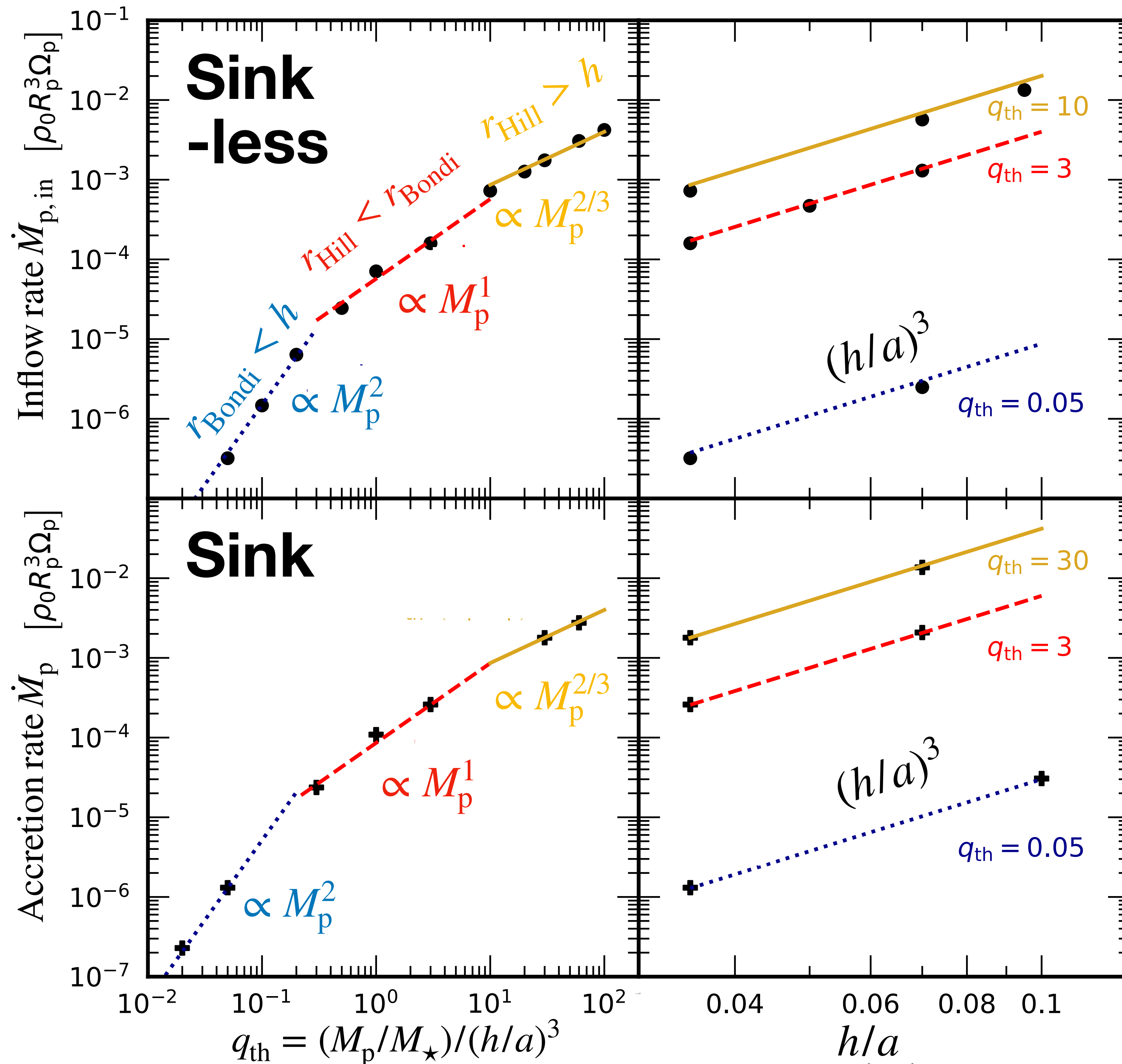
$$\propto M_p^2 (h/a)^{-3}$$

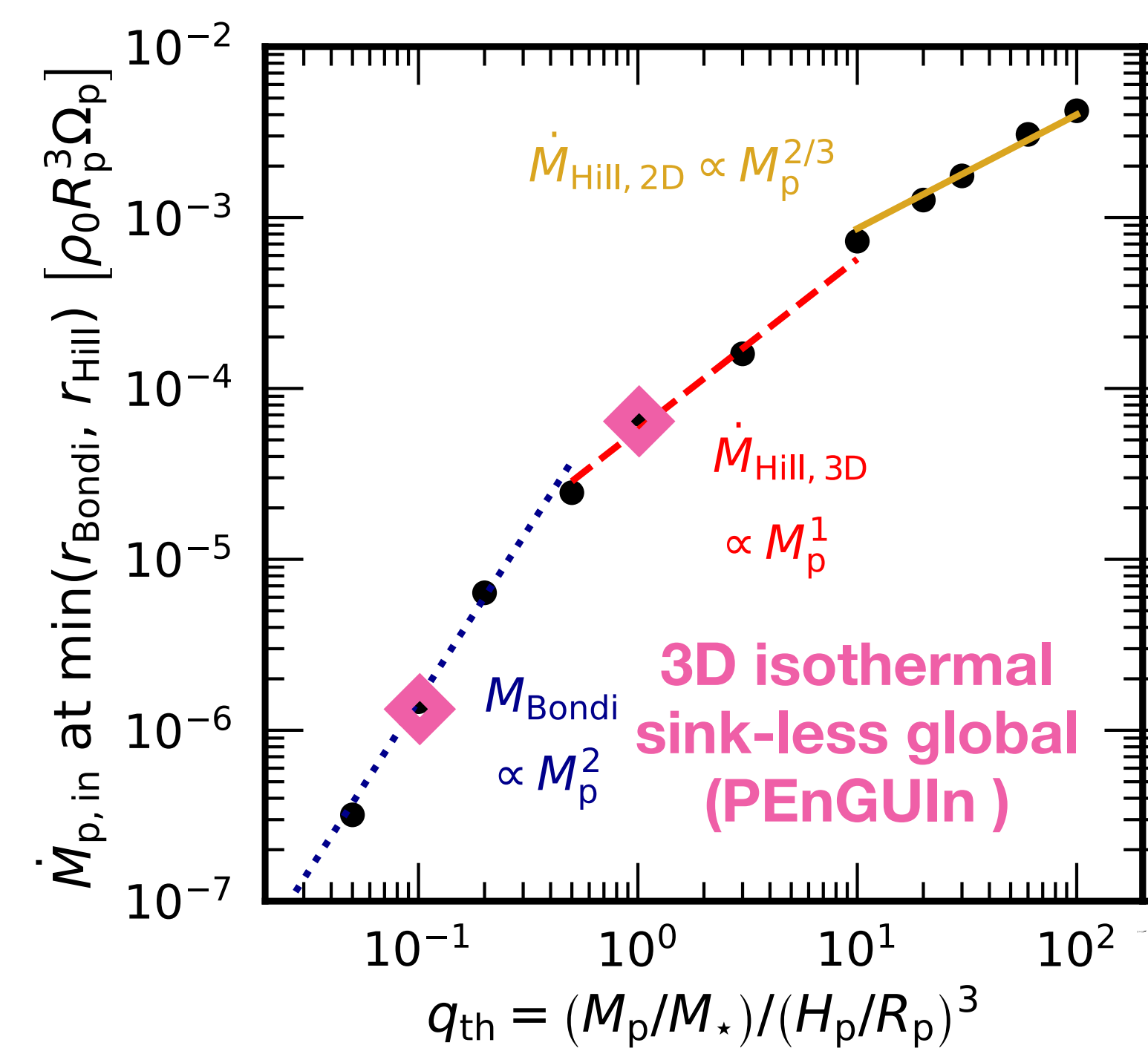
$$\dot{M}_{\text{Hill,3D}} \sim r_{\text{Hill}}^2 \times (3/2)\Omega r_{\text{Hill}} \times \rho$$

$$\propto M_p^1 (h/a)^0$$

$$\dot{M}_{\text{Hill,2D}} \sim r_{\text{Hill}} h \times (3/2)\Omega r_{\text{Hill}} \times \rho$$

$$\propto M_p^{2/3} (h/a)^1$$





Scalings supported by other works

$$\dot{M}_{\text{Bondi}} \propto M_p^2$$

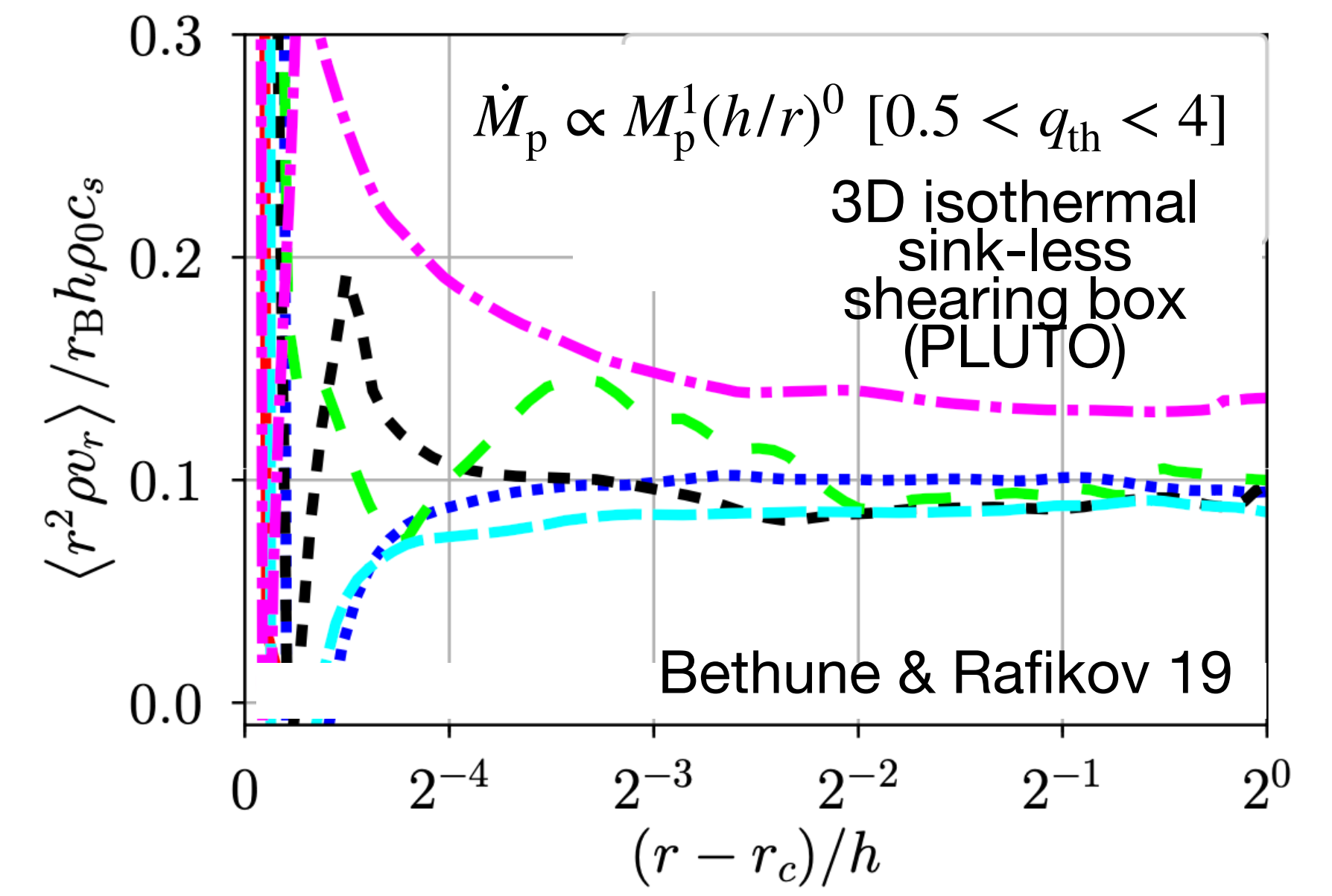
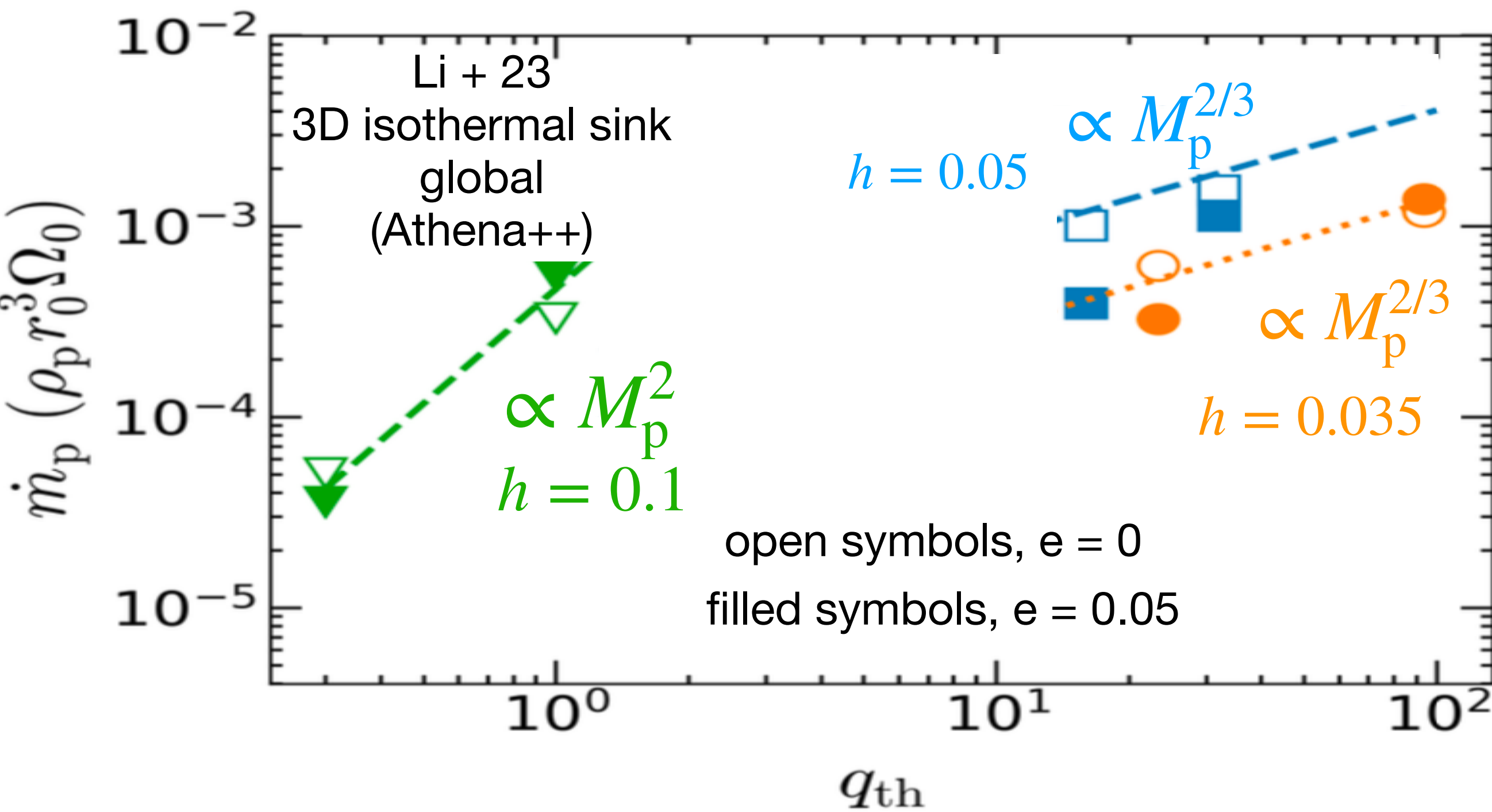
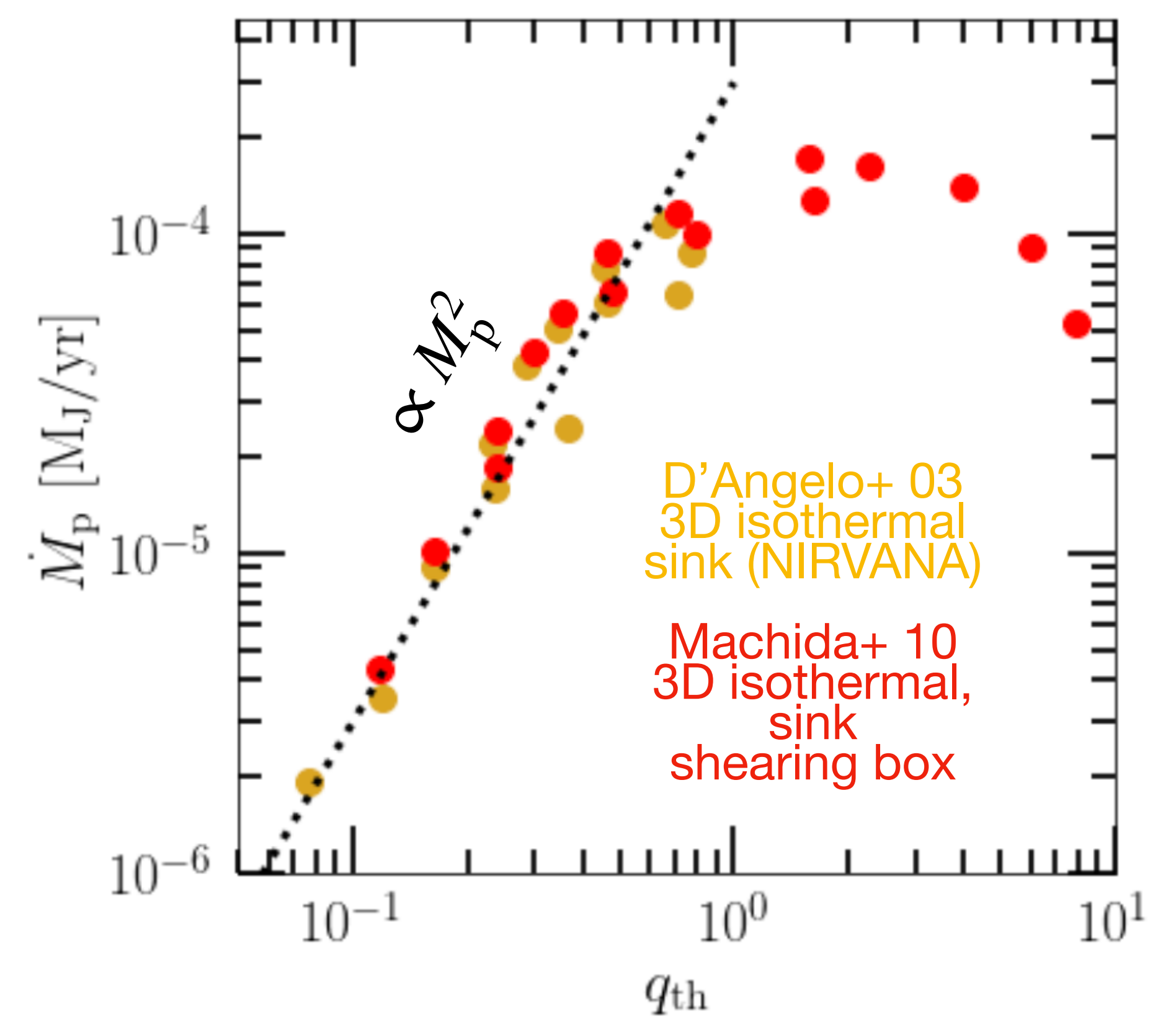
D'Angelo+ 03, Machida+ 10, Li+ 23

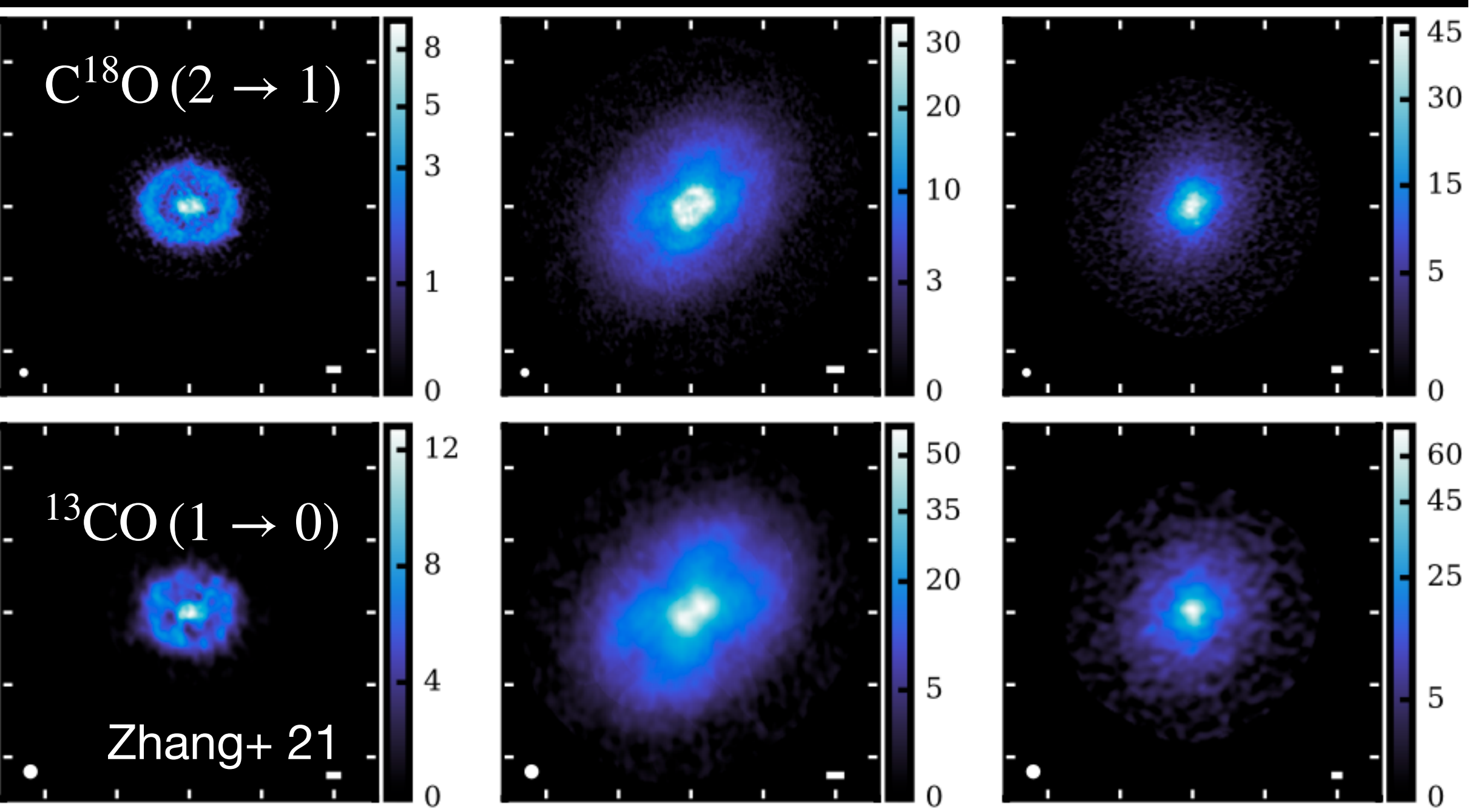
$$\dot{M}_{\text{Hill}} \propto M_p^1$$

Bate+ 03, Bethune & Rafikov+ 19

$$\dot{M}_{\text{Hill, 2D}} \propto M_p^{2/3}$$

Zhu+ 12, Li+ 23





$\min(t_{\text{double}}) \ll t_{\text{age}}$
 $\therefore$ planets accrete below
 theoretical max by $\sim 10\text{-}100\times$

from
 gap-fitting
 (Zhang+ 18,
 Dong & Fung 18)

**Use optically thin
 CO isotopologues
 to estimate gap depth**

$$\Sigma_{\text{H}_2} = X_{\text{CO}} N_{\text{CO}}$$

$$N_{\text{CO}} \sim X_{\text{iso}} \frac{k_B}{\nu S d^2} \frac{Q}{g_J} e^{E_u/k_B T} I_\nu \Delta \nu$$

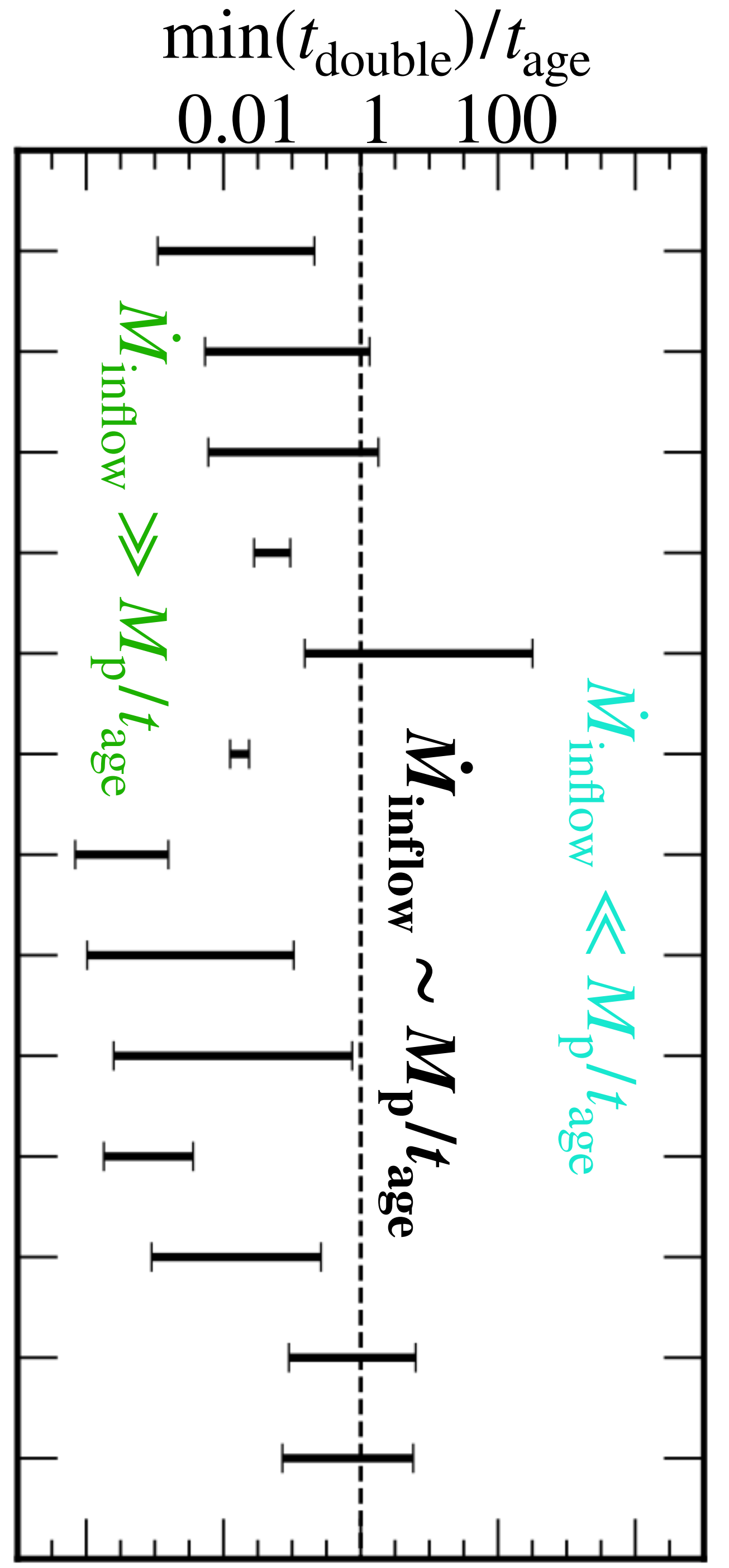
e.g. Mangum & Shirley 15

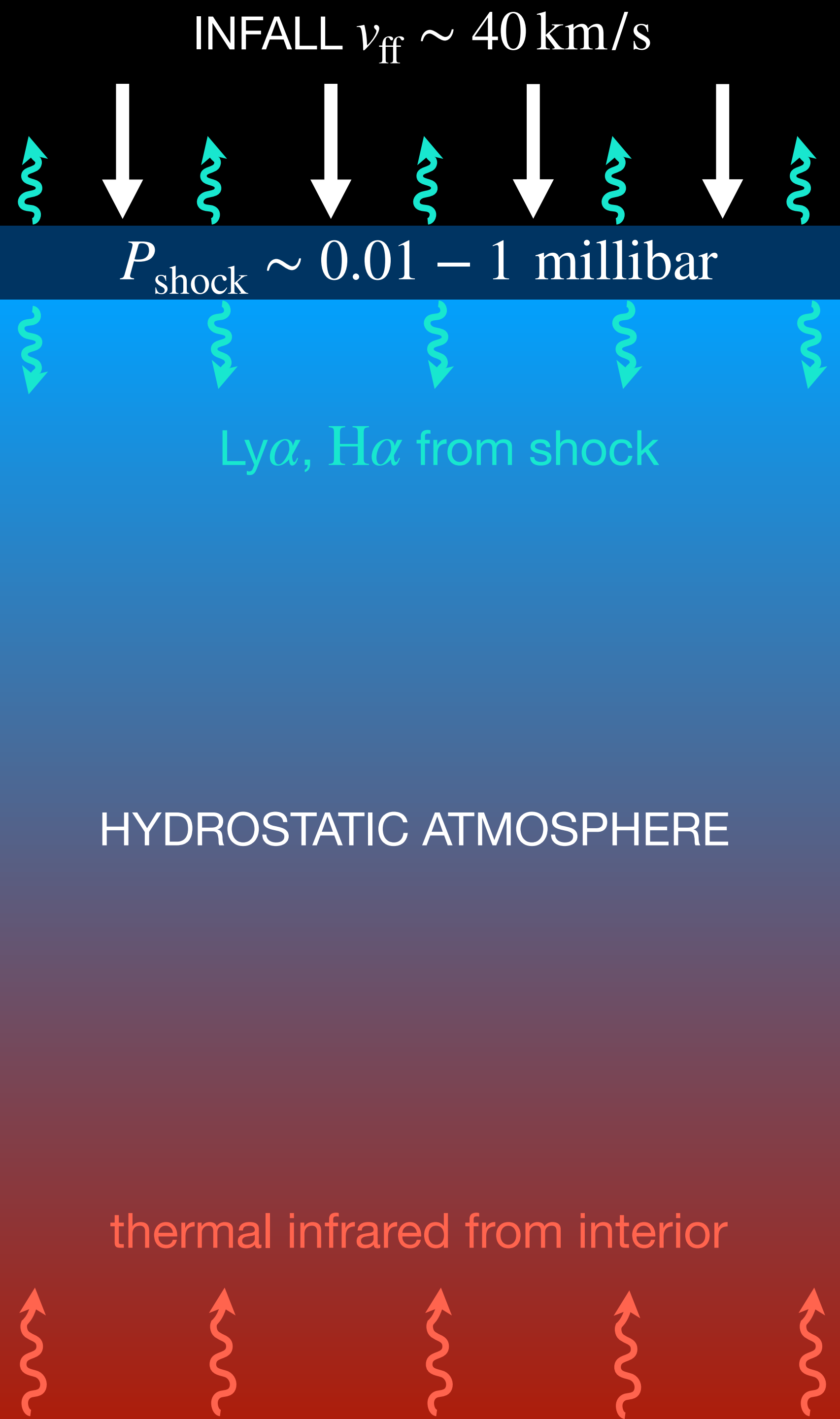
$$\min(t_{\text{double}}) \equiv \frac{M_p}{\dot{M}_{\text{inflow}}}$$

new scalings
 + gap densities from CO

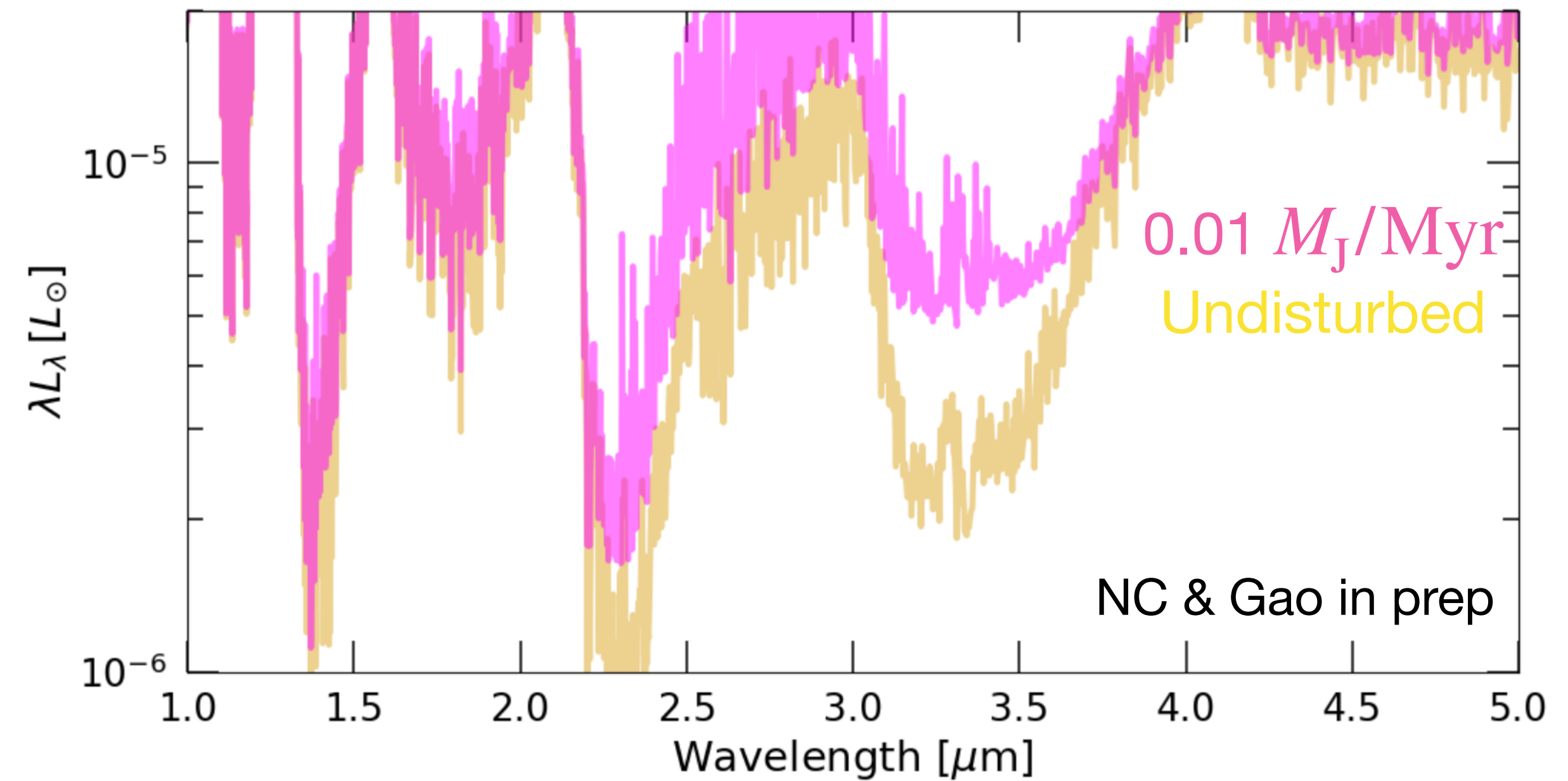
- TW Hya-G1 (21)
- TW Hya-G2 (85)
- IM Lup (117)
- AS 209-G2 (99)
- AS 209-G3 (240)
- HD 169142 (37)
- HD 163296-G2 (48)
- HD 163296-G3 (86)
- HD 163296-G4 (137)
- HD 163296-G234alt (108)
- HD 163296-G5 (260)
- PDS 70b (22)
- PDS 70c (34)

NC & Chiang 22, NC+ 23

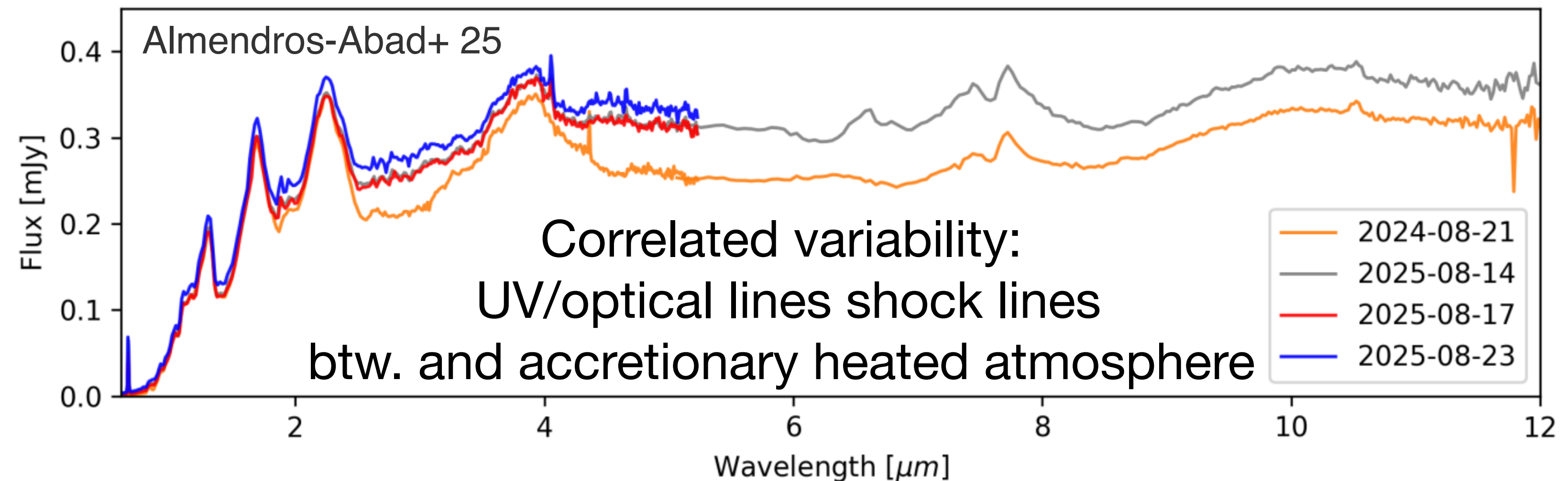




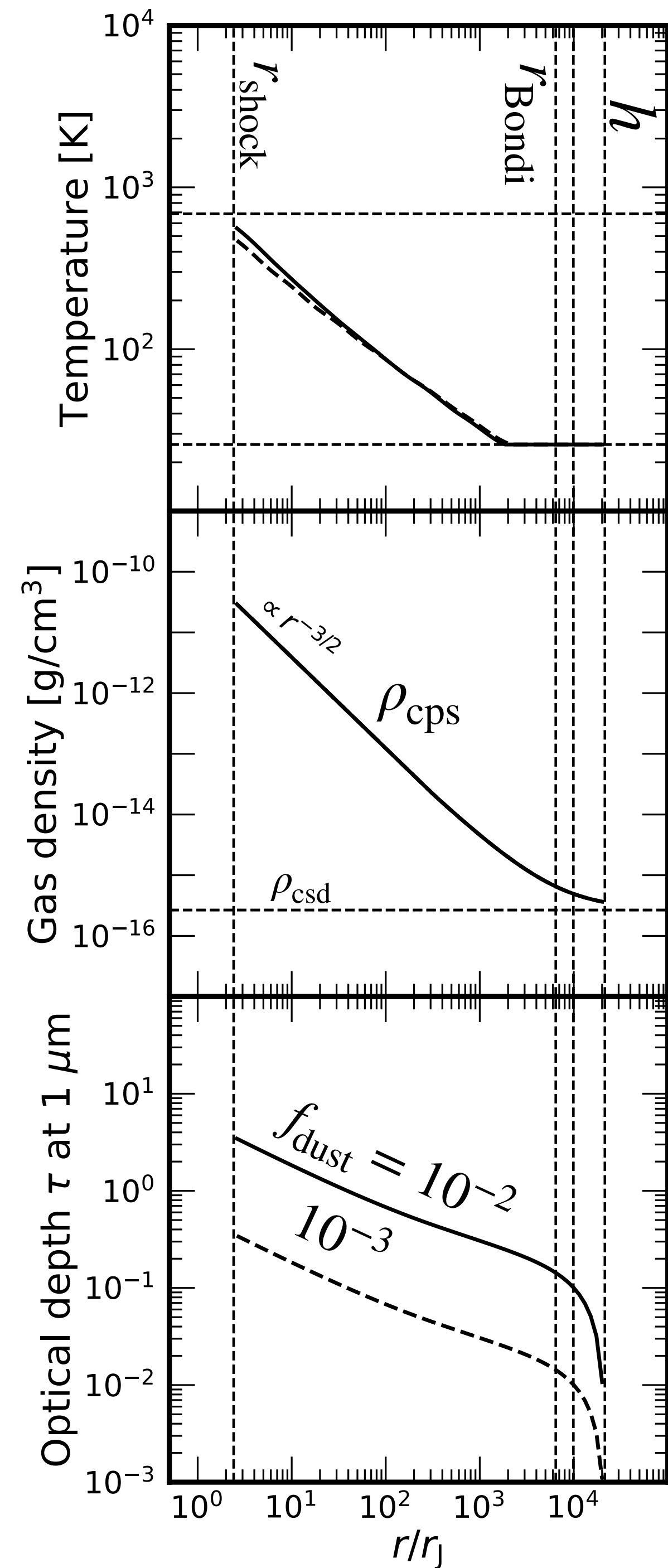
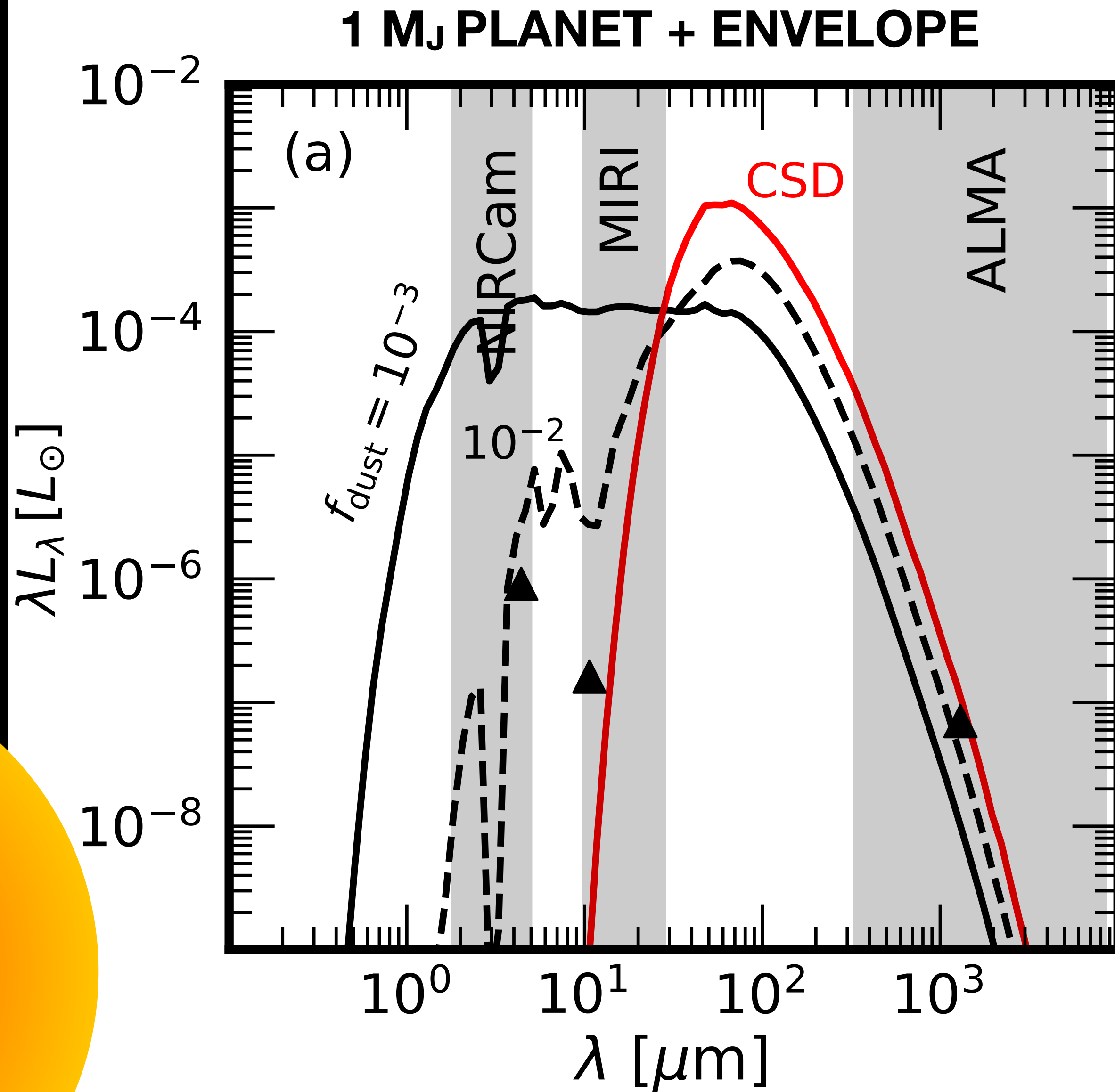
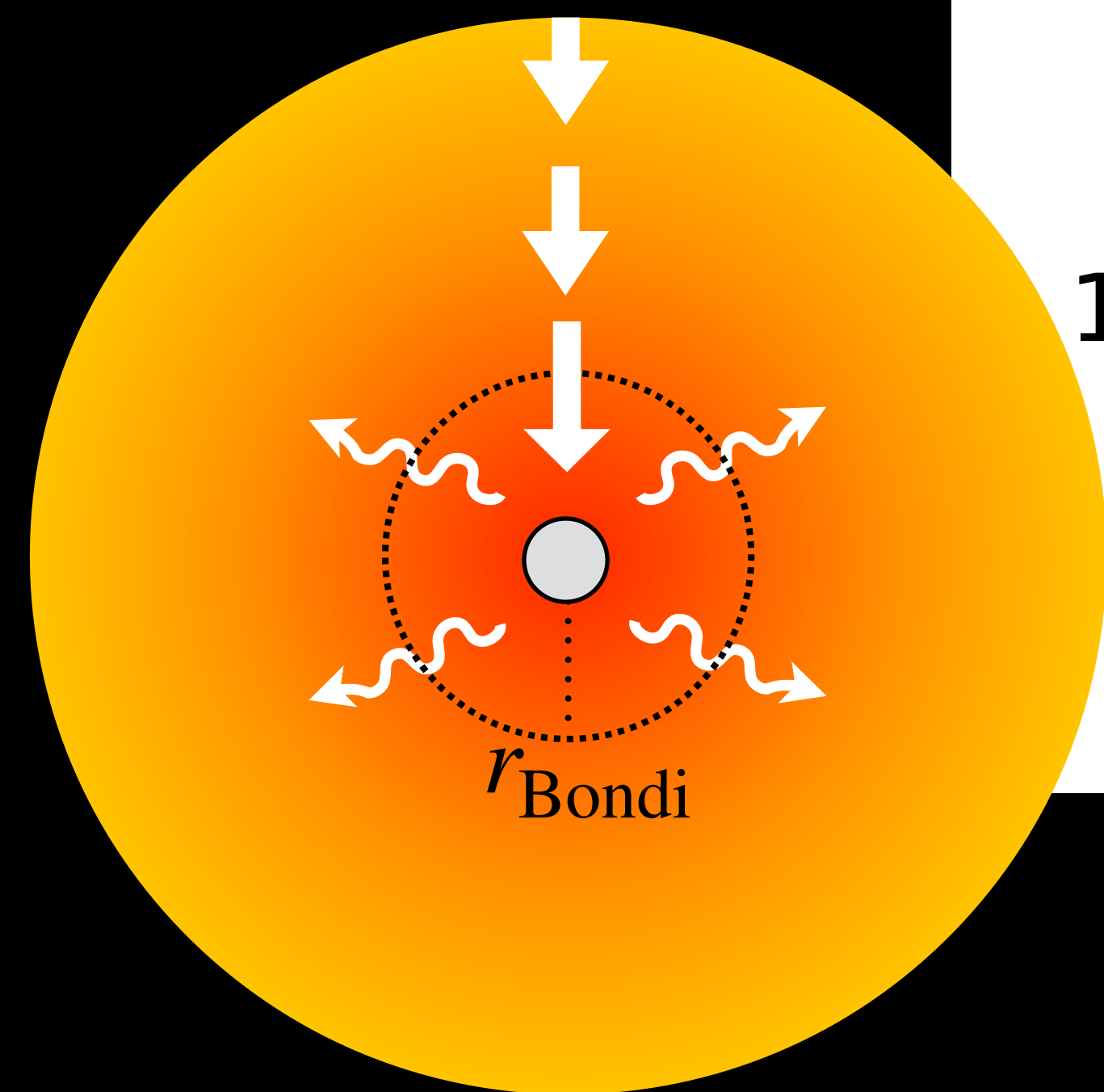
ACCRETIONALLY HEATED ATMOSPHERES

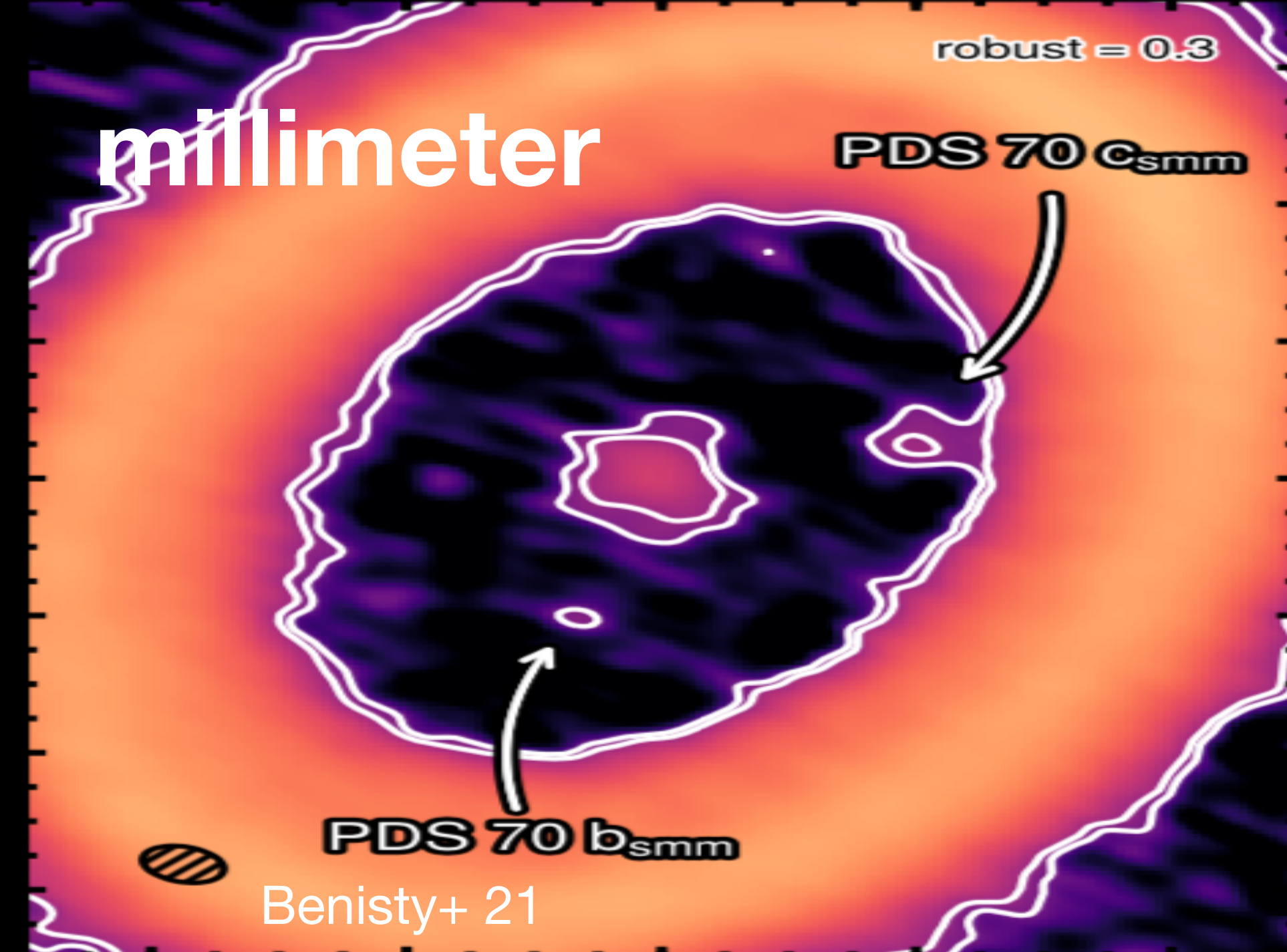
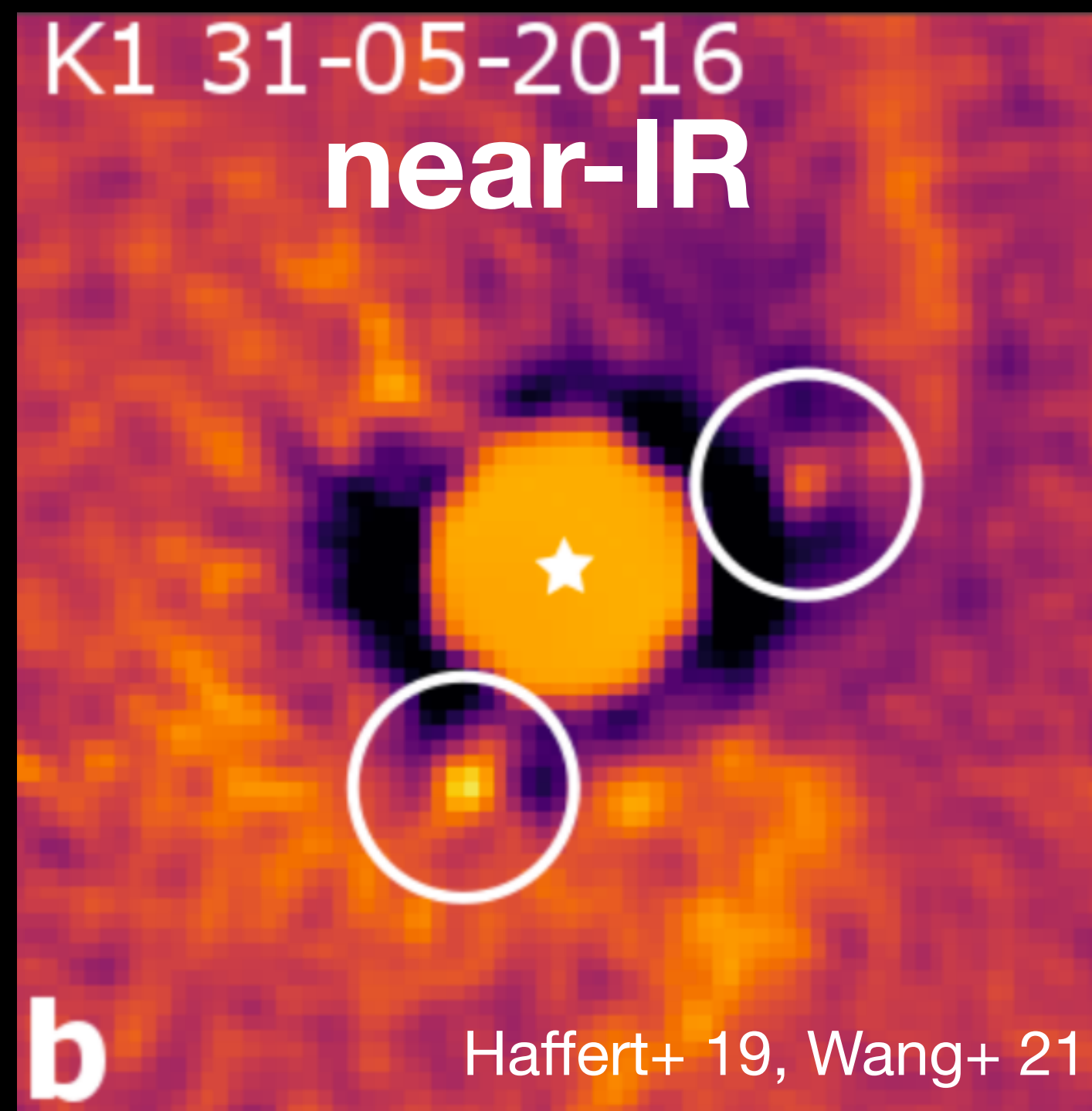
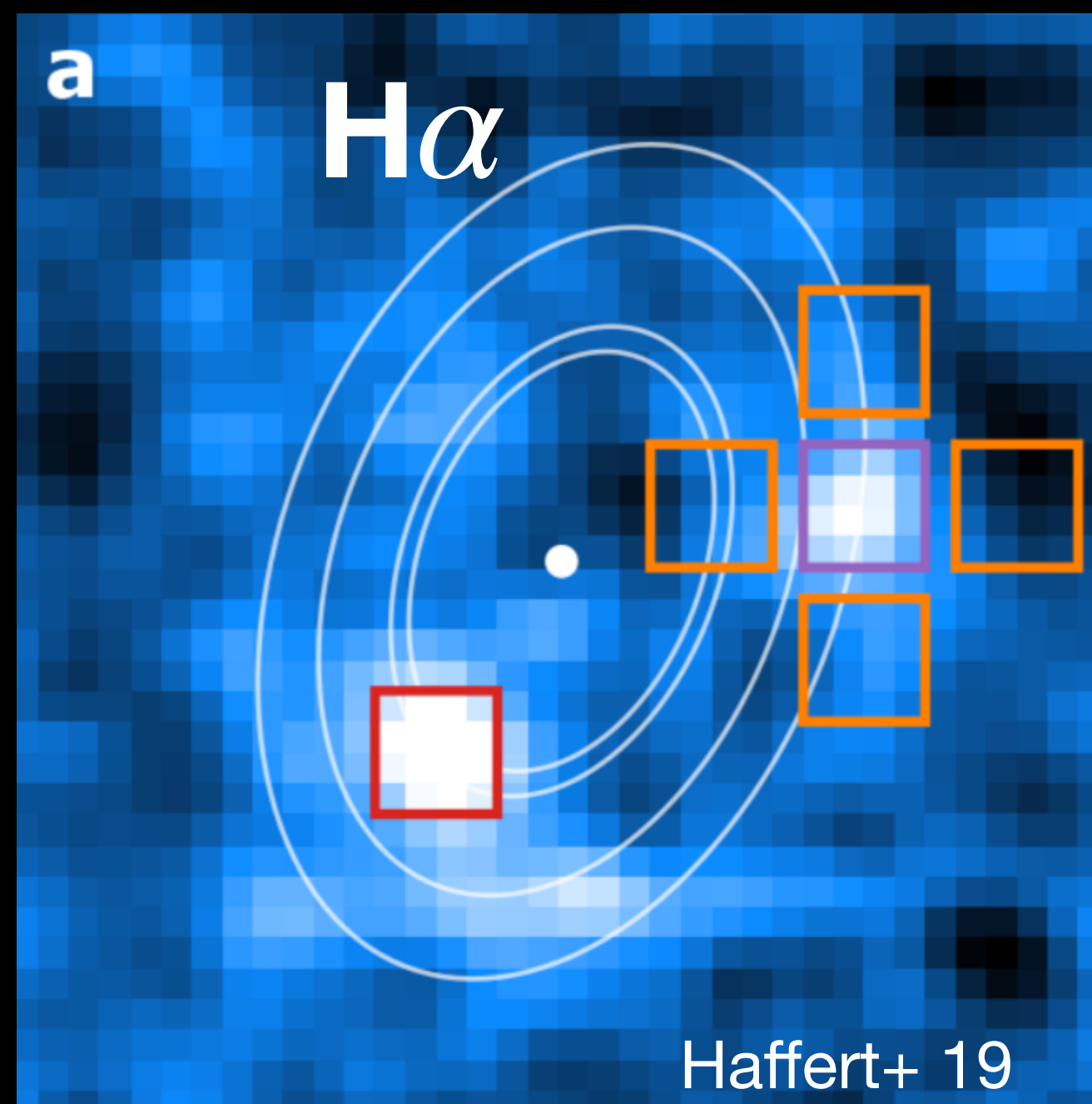


Elevated NIR continuum from an accreting brown dwarf



REPROCESSING BY DUSTY CIRCUM- PLANETARY ENVELOPES





PDS 70 across the spectrum

two giant planets in a disk gap

$$M_p \sim 1 - 10 M_J$$

from cavity size, SED,
and dynamical stability

Bae+ 19, Ginzburg & Chiang 19, Toci+ 20, Wang+ 21

(1) $\langle \dot{M}_p \rangle \sim 3 M_J / 5 \text{ Myr}$
 $\sim 0.6 M_J / \text{Myr}$

(2) $\max \dot{M}_p \sim r_{\text{Hill}} h \times (3/2) \Omega r_{\text{Hill}} \times \rho$
 $\sim 1 M_J / \text{Myr}$

(3) $\dot{M}_p \sim L_{\text{bol}} 4 R_J / (G \times 3 M_J)$
 $\sim 1 M_J / \text{Myr}$ NC & Chiang 22, NC+ 23