

### Main Points

We present a first-principles mechanism demonstrating how a weakly ionized accretion disk sheds angular momentum and produces the electric power for driving bidirectional astrophysical jets.

1. The basic conserved quantity in an axisymmetric system with a magnetic field is the canonical angular momentum  $mrv\theta + q\psi/2\pi$  rather than the ordinary angular momentum  $mrv\theta$ . This comes from fundamental Hamiltonian/Lagrangian mechanics.  $\psi$  is the poloidal magnetic flux.

2. When the Kepler angular velocity and the magnetic field have opposite polarity, collisions between neutrals and charged particles cause:

- ions to move radially inwards,
- electrons to move radially outwards,
- both ions and electrons to gain canonical angular momentum,
- neutrals to lose ordinary angular momentum and move radially inwards (accrete)
- the total canonical angular momentum of the system (electrons, ions, and neutrals) to be conserved.

3. The accumulation of ions at small radius and accumulation of electrons at large radius produces a radially outward electric field.

4. In 3D, this radial electric field would drive an out-of-plane poloidal current that produces the magnetic forces that drive bidirectional astrophysical jets.

### Mathematical Explanation

#### Motion of charged particles colliding with Kepler-motion neutrals

Equation of motion for charged particles

$$\frac{d\mathbf{u}_\sigma}{dt} = -\frac{GM_*}{r^2} \hat{r} + \omega_{c\sigma} \mathbf{u}_\sigma \times \hat{z} - \nu_{\sigma n} (\mathbf{u}_\sigma - \mathbf{u}_n) \approx \omega_{c\sigma} \mathbf{u}_\sigma \times \hat{z} - \nu_{\sigma n} (\mathbf{u}_\sigma - \mathbf{u}_n)$$

The solution is

$$\mathbf{u}_\sigma = \left( \mathbf{u}_L \cos \omega_{c\sigma} t + \mathbf{u}_L \times \hat{z} \sin \omega_{c\sigma} t \right) e^{-\nu_{\sigma n} t} + \frac{\frac{\nu_{\sigma n} \mathbf{u}_n \times \hat{B} + \frac{\nu_{\sigma n}^2}{\omega_{c\sigma}^2} \mathbf{u}_n}{1 + \frac{\nu_{\sigma n}^2}{\omega_{c\sigma}^2}}}{\omega_{c\sigma}} \mathbf{u}_\sigma = \frac{\frac{\nu_{\sigma n} \mathbf{u}_n \times \hat{B} + \frac{\nu_{\sigma n}^2}{\omega_{c\sigma}^2} \mathbf{u}_n}{1 + \frac{\nu_{\sigma n}^2}{\omega_{c\sigma}^2}}}{\omega_{c\sigma}} \mathbf{u}_\sigma = \frac{\xi_\sigma}{1 + \xi_\sigma^2} u_{n\theta} \hat{\theta}$$

$$P_\theta = mrv_\theta + \frac{qBr^2}{2}$$

Ions move inward with increasing CAM

$$u_{ir} < 0 \quad P_{\theta i} \nearrow$$

Electrons move outward with increasing CAM

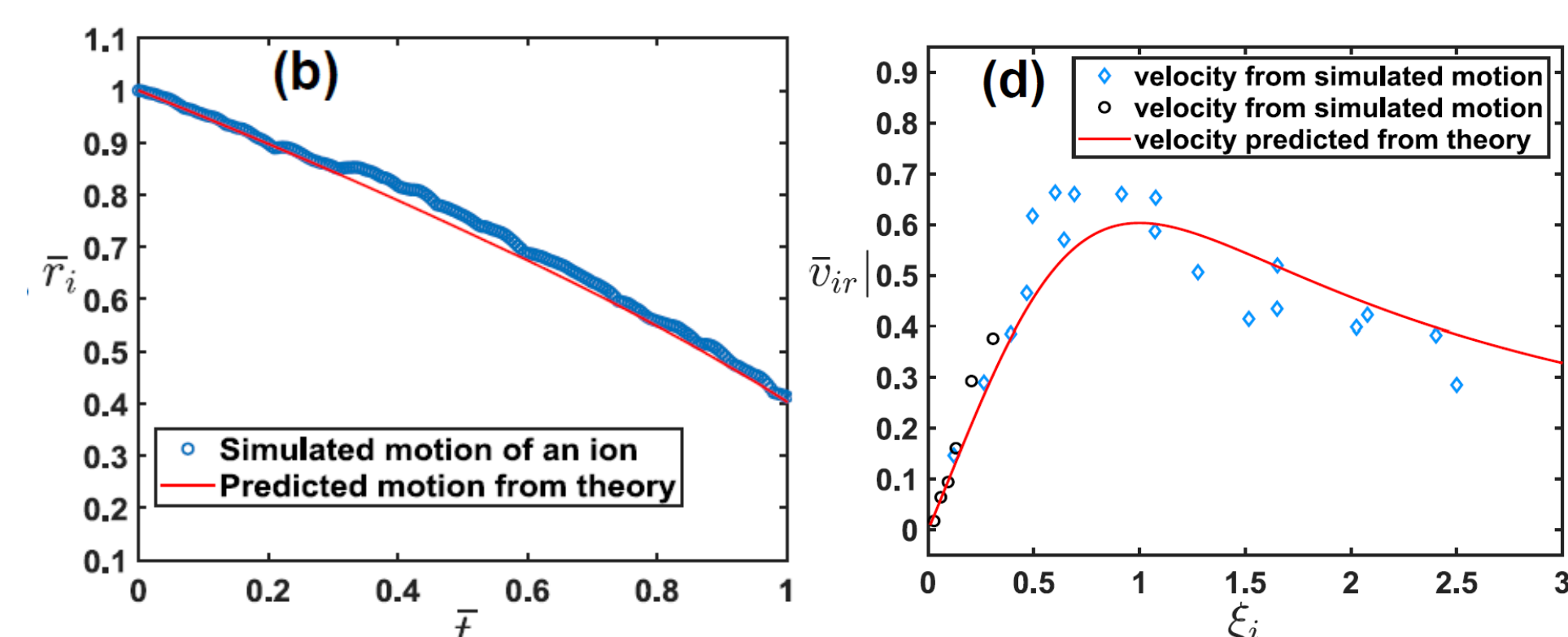
$$u_{er} > 0 \quad P_{\theta e} \nearrow$$

Neutrals accrete with decreasing CAM

$$u_{nr} < 0 \quad P_{\theta n} \searrow$$

Simulation case

Insensitive to the magnetic field polarity!



#### Neutral's radial velocity

At a specific radius  $r$ , the canonical angular momentum density at this radius is

$$P_\theta = n_n m_n r_n u_{n\theta} + n_i m_i r_i u_{i\theta} + n_e m_e r_e u_{e\theta} + \frac{1}{2} n_i q_i B r_i^2 + \frac{1}{2} n_e q_e B r_e^2$$

During a small time  $\Delta t$ , the change of CAM density is

$$\Delta P_\theta = n_n m_n (\Delta r_n u_{n\theta} + r_n \Delta u_{n\theta}) + n_i m_i (\Delta r_i u_{i\theta} + r_i \Delta u_{i\theta}) + n_e m_e (\Delta r_e u_{e\theta} + r_e \Delta u_{e\theta}) + n_i q_i B r_i \Delta r_i + n_e q_e B r_e \Delta r_e = 0$$

During collisions,  $n_n m_n r_n \Delta u_{n\theta} + n_i m_i r_i \Delta u_{i\theta} + n_e m_e r_e \Delta u_{e\theta} = 0$

$$n_i m_i \Delta r_i u_{i\theta} + n_e m_e \Delta r_e u_{e\theta} \ll n_i q_i B r_i \Delta r_i + n_e q_e B r_e \Delta r_e$$

$$\Delta P_\theta \approx n_n m_n \Delta r_n u_{n\theta} + n_i q_i B r_i \Delta r_i + n_e q_e B r_e \Delta r_e = 0$$

$$u_{nr} = -\frac{\omega_{ci} m_i}{\omega_{ce} m_e} (\chi_i u_{ir} - \chi_e u_{er})$$

Mass accretion rate

$$\dot{M} = 2\pi r u_{nr} \Sigma \quad \chi_\sigma = \frac{n_\sigma}{n} \quad \xi_\sigma = \frac{\nu_{\sigma n}}{\omega_{c\sigma}}$$

$$\dot{M} = 2\pi \chi_n n_n r^2 h |qB| \left( \frac{|\xi_i|}{1 + \xi_i^2} + \frac{|\xi_e|}{1 + \xi_e^2} \right)$$

From our model

$$\dot{M} = 3 \times 10^{-8} M_\odot \cdot \text{year}^{-1}$$

From observation

$$\dot{M} = 10^{-9} - 10^{-7} M_\odot \cdot \text{yr}^{-1}$$

### Simulation Results

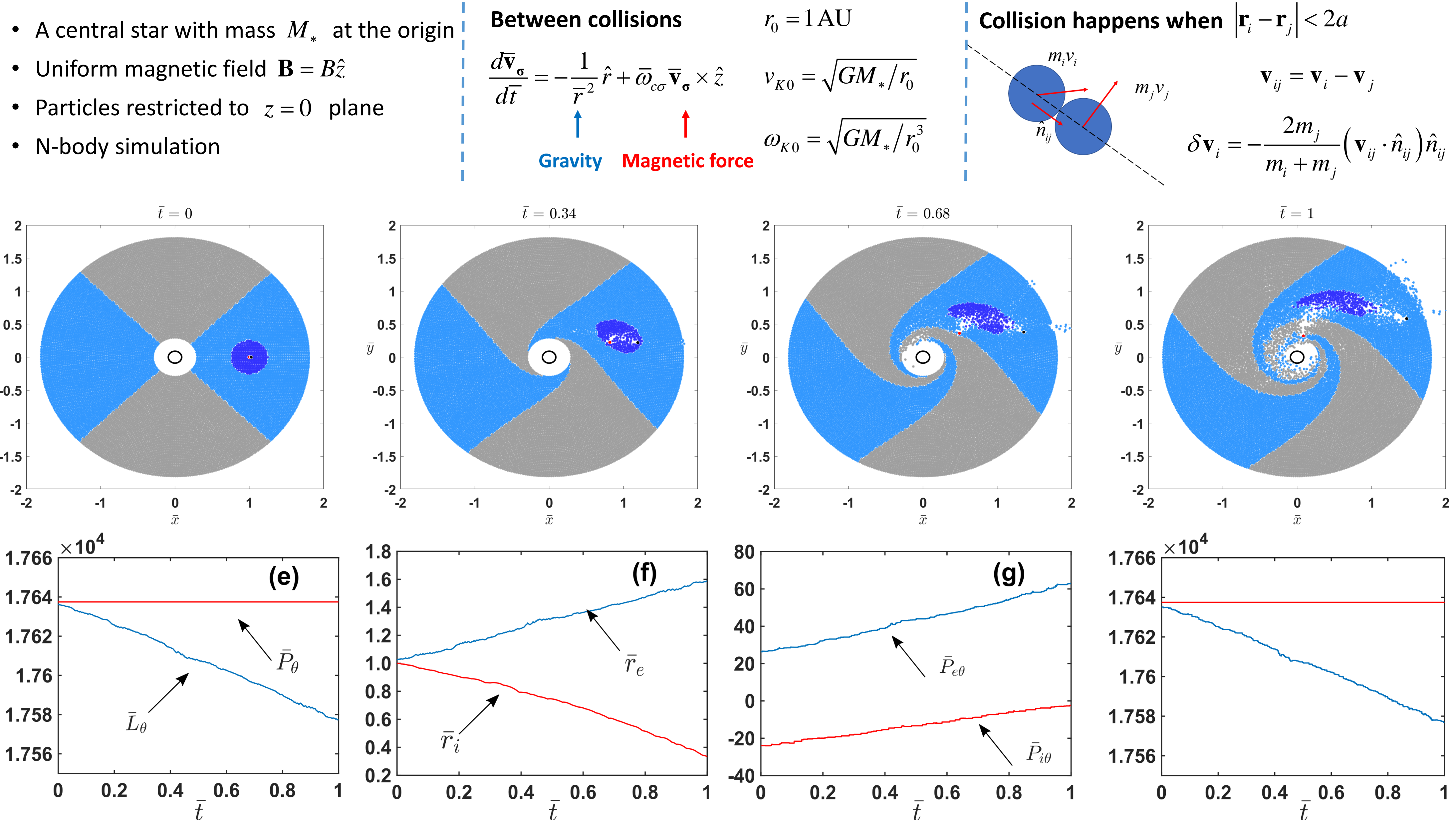


Figure 1. Simulation results when there is an electron-ion pair in the system. (a)-(d) The particle trajectories of the whole system. Neutrals are blue and gray. The ion is red, and the electron is black. The neutrals surrounding the electron-ion pair at  $\bar{t} = 0$  are dark blue. (e) The total canonical angular momentum of the system and the total ordinary angular momentum of the system. (f) The radial positions of the ion and electron. (g) The canonical angular momentum of the ion and electron. (h) The ordinary angular momentum of the neutrals and the total canonical angular momentum of the system

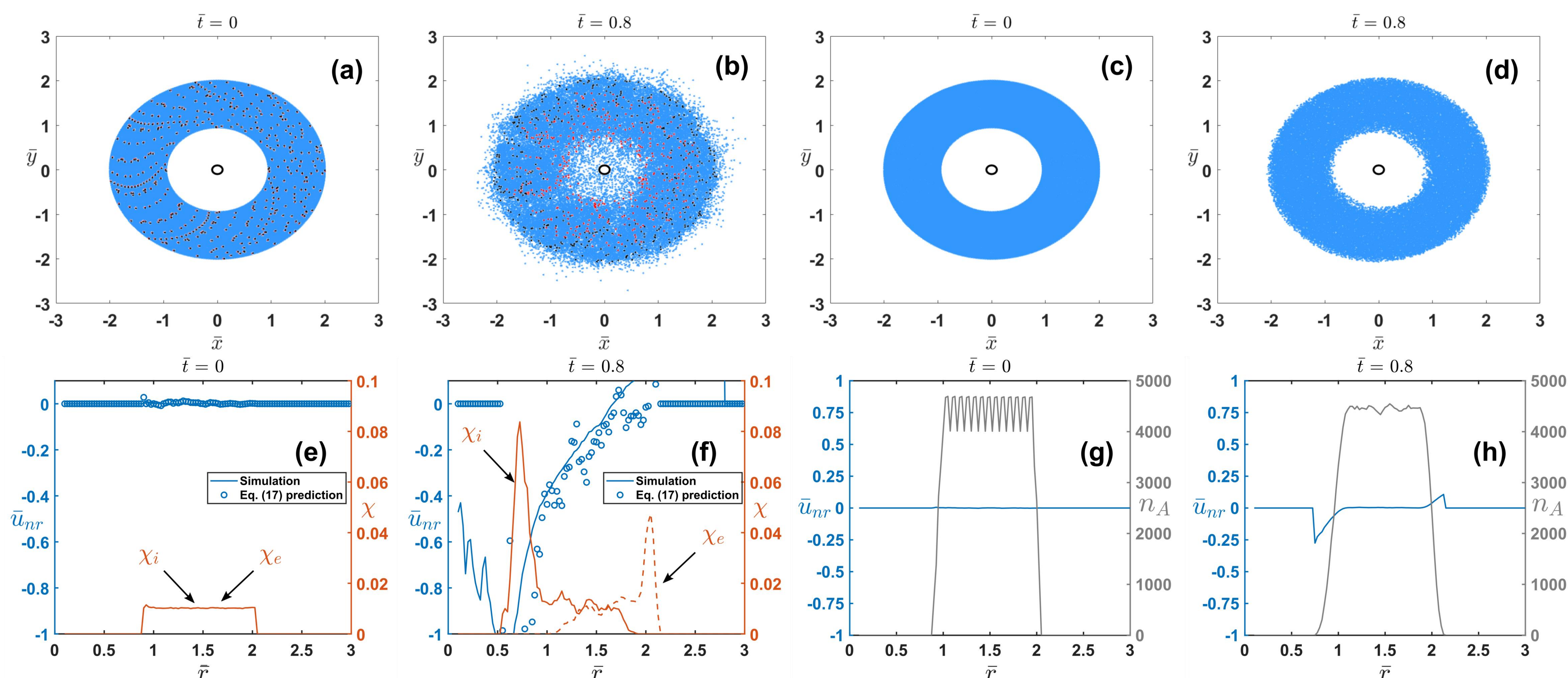
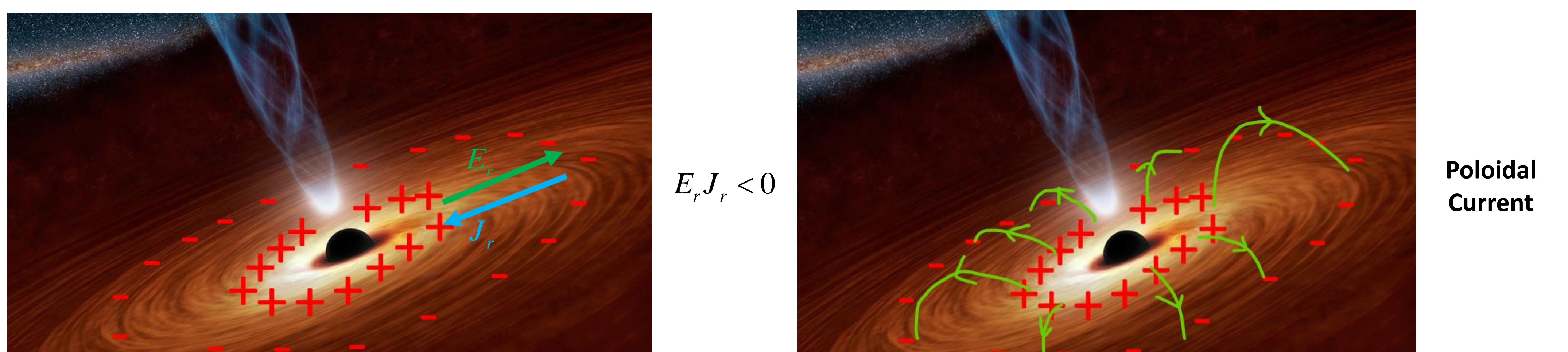


Figure 2. (a), (b) The particle trajectories of a system with ions and electrons. The initial velocity is a Kepler velocity plus a random thermal velocity having 10% of the Kepler velocity magnitude. (c), (d) The particle trajectories of a reference system having neutrals only with the same initial condition for neutrals as in (a). (e), (f) The neutral radial drift velocity profile and the density fraction of ions and electrons of the system in (a), (b). The blue line is the radial drift velocity profile of neutrals obtained from the simulation. The blue circles are the radial velocity of neutrals calculated as a function of the ion radial velocity and electron radial velocity as predicted by Equation (17). The red solid/dashed line shows the ion/electron density fraction vs. radial position. (g), (h) The neutral drift velocity profile and neutral surface density  $n_A$  vs. radial position of the system of (c), (d).

### Astrophysical Jets Generation



#### Caltech Astrophysical Jets Experiment

